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Impact of non-axisymmetric magnetic field perturbations on flows

R.Varenes¹, X.Garbet¹, L.Vermare², Y.Sarazin¹, V.Grandgirard¹, G.Dif-Pradalier¹, M.Peret²

¹CEA, IRFM, France ²LPP, CNRS, Ecole polytechnique, France

Introduction

- Plasma rotation plays a key role in plasma confinement [1][2];
- Control of plasma rotation in reactor-sized tokamak is challenging [3];
- Intrinsic bulk plasma rotation is driven by turbulence and **Neoclassical Toroidal Viscosity (NTV)**;
- Magnetic field ripple is responsible for the NTV.

Objective: Understand the competition/synergy between turbulence and NTV with gyrokinetic simulations.

Theoretical model

- A reduced model based on the mean toroidal velocity V_T reads :

$$\frac{\partial V_T}{\partial t} = \text{Neoclassical Toroidal Viscosity} + \text{Turbulent torque}$$

Ripple constrains V_T through neoclassical friction v_ϕ

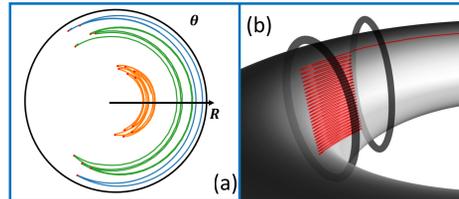


FIG.1 – Kinetic effects induced by ripple : drift of banana bounce points (a) and toroidal trapping between coils (b). Neoclassical friction v_ϕ comes from the collisions between trapped populations.

Turbulence constrains V_T through turbulent viscosity χ_{turb}

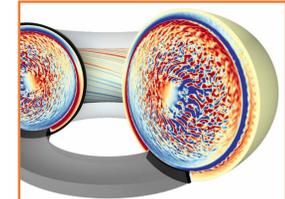


FIG.2 – Snapshot of the turbulent structures seen through a colormap on the electric potential. From rest, V_T grows due to wave-particle interactions.

Goal: Obtain coefficients of **turbulent momentum transport from simulations** and use **neoclassical friction v_ϕ predictions** to find the ripple amplitude threshold for which **neoclassical effects overcome turbulence**.

Turbulent momentum transport

- Turbulence impacts V_T through the toroidal Reynold's stress Π_ϕ
- Axisymmetric theory [4][5] :
$$\Pi_\phi = \underbrace{-\chi_{turb} \frac{\partial V_T}{\partial r}}_{\text{Viscosity}} + \underbrace{\mathcal{V} V_T}_{\text{Pinch}} + \underbrace{\Pi_{res}}_{\text{Residual}} \quad (\text{Eq.1})$$
- Obtention of χ_{turb} , \mathcal{V} and Π_{res} using gyrokinetic simulations :
 - Simulations of turbulent plasma without ripple with different initial toroidal velocity → FIG.3 shows the dominance of the viscosity : $-\chi_{turb} \frac{\partial V_T}{\partial r}$;
 - Eq.1 defines a plane in the $(V_T, \partial_r V_T, \Pi_\phi)$ space;
 - Mean square plane fit using simulations output gives an estimation of χ_{turb} within a radial range (see FIG.4 for an example).

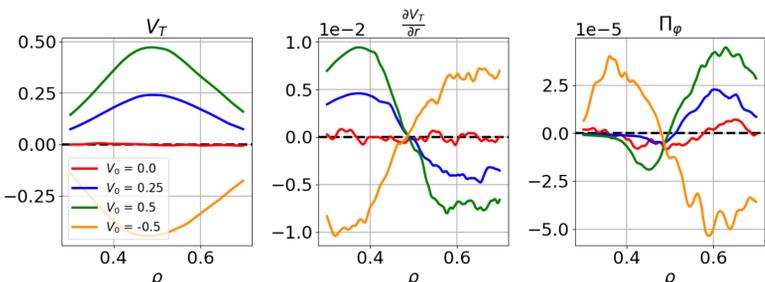


FIG.3 - Radial profiles of V_T , $\partial_r V_T$ and Π_ϕ taken at turbulent saturation for simulations initialized at different toroidal velocity $V_T(t=0) = V_0 e^{4(\rho-0.5)^2}$

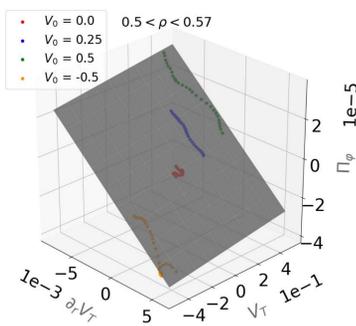


FIG.4 - Plane fit of Π_ϕ using point cloud from simulations giving an estimation of the turbulent viscosity within a radial range.

Neoclassical friction

- Ripple is responsible for magnetic braking m in the toroidal direction : the NTV;
- Neoclassical theory [6][7] with ripple gives :

$$m = -v_\phi \left(V_T - k_T \frac{\nabla T}{e B_p} \right)$$

\downarrow Friction
 \downarrow Thermal drive

- GYSELA code shows good agreement with neoclassical predictions;
- v_ϕ increases with the ripple amplitude.

Competition ripple/turbulence

- Complete model reads :
$$\frac{\partial V_T}{\partial t} \approx m - \nabla \cdot \Pi_\phi \Rightarrow \tau_V \sim \max \left(\frac{\chi_{turb}}{L_V^2}, v_\phi \right)^{-1}$$

τ_V relaxation time

L_V^2 gradient's length
- Boundary physics not in the model → radially gaussian ripple (cf FIG.5)

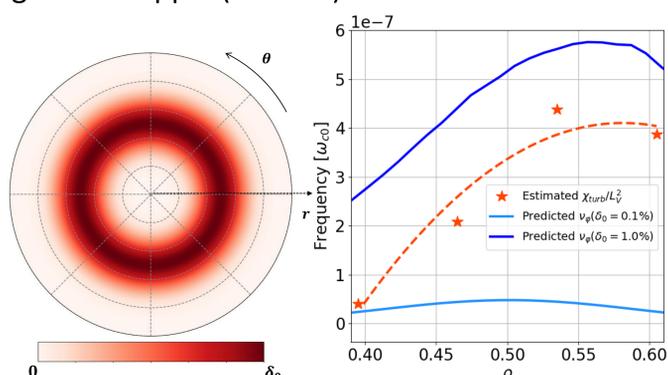


FIG.5 – Polar map of the ripple amplitude used for simulations.

FIG.6 – Radial profiles of estimated turbulent relaxation frequency and predicted neoclassical frictions.

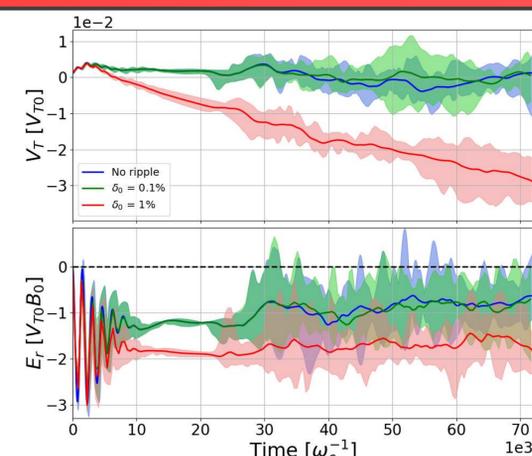


FIG.7 – Time evolution of V_T and E_r in the range $0.45 < \rho < 0.55$ for different ripple amplitudes.

- Model verification using 2 simulations :
 - With **dominant neoclassical friction**
 - With **dominant turbulent viscosity**
- Main result FIG.7
 - $\delta_0 = 0.1\%$ case dominated by turbulence
 - $\delta_0 = 1.0\%$ case driven by neoclassical friction → increase of E_r to fulfill force balance

Conclusion

- Ripple is responsible for a neoclassical friction that constrains the toroidal velocity;
- Turbulence is a source of intrinsic rotation;
- Evolution of mean toroidal flow ruled by competing turbulent stress and ripple drag forces;
- The radial electric field grows in response to the modification of the toroidal velocity;
- Simulations suggest that neoclassic drag overcomes turbulent stress for typical realistic ripple amplitudes in WEST;
- Future work taking into account boundary physics with realistic WEST ripple amplitude is necessary to understand the shape of $\frac{dE_r}{dr}$ and the transition toward high-confinement modes.

References

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