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***In situ* analysis of weakly bound proteins reveals molecular basis of soft corona formation**

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ABSTRACT: Few experimental techniques allow the analysis of the protein corona *in situ*. As a result, little is known on the effects of nanoparticles on weakly bound proteins that form the soft corona. Despite its biological importance, our understanding of the molecular bases driving its formation is limited. Here we show that hemoglobin can form either a hard or a soft corona on silica nanoparticles depending on the pH conditions. Using cryoTEM and synchrotron-radiation circular dichroism, we show that nanoparticles alter the structure and the stability of weakly bound proteins *in situ*. Molecular dynamics simulation identified the structural elements driving protein-nanoparticle interaction. Based on thermodynamic analysis, we show that nanoparticles stabilize partially unfolded protein conformations by enthalpy-driven molecular interactions. We suggest that nanoparticles alter weakly bound proteins by shifting the equilibrium towards the unfolded states at physiological temperature. We show that the classical approach based on nanoparticle separation from the biological medium fails to detect destabilization of weakly bound proteins, and therefore cannot be used to fully predict the biological effects of nanomaterials *in situ*.

KEYWORDS. Protein corona, silica nanoparticles, cryoTEM, synchrotron-radiation circular dichroism, molecular dynamics, protein unfolding.

In a biological environment, the adsorption of proteins on nanoparticles (NPs) is an early process that leads to major changes of the physicochemical properties and the biological effects of NPs.¹⁻³ The formation of the so-called protein corona drives NP-cell interaction, NP biodistribution, and NP toxicity *in vivo*.⁴⁻⁷ Numerous studies have investigated the composition and dynamics of the protein corona in various biological systems on one hand, and the changes of the structure and activity of adsorbed proteins on the other.⁸⁻¹² The current paradigm is that one protein may or may not have affinity for one NP, depending on (i) the protein structural and biochemical properties, (ii) the NP size, shape, and surface chemistry, (iii) the pH, ionic strength and composition of the medium.¹³⁻¹⁶ In a model system where one protein encounters one NP, this would result in either the adsorption or the non-adsorption of the protein on the surface. In a more complex biological environment where competitive adsorption occurs, protein adsorption will also depend on the affinity scale of the different proteins and on their concentration in the medium. Proteins with high abundance and lower affinity may be gradually replaced by proteins with higher affinity.¹⁷ The formation of the protein corona is a dynamic process that evolves

with time,^{8,18} changes of the biological medium,¹⁹ and through cell processing.²⁰ *In vitro* and *in vivo* studies have recently shown that flow plays a significant role in shaping the protein corona on NPs in a biological environment.²¹⁻²³

Methodologies have been developed to identify and quantify proteins adsorbed on NPs.²⁴ In most cases, the adsorbed proteins are separated from the free proteins in solution by centrifugation. The pellet (which contains NPs coated with proteins) and the supernatant (which contains free proteins) are then analyzed by proteomics to determine the composition of the corona. These information are used to determine the structural and physicochemical parameters driving protein adsorption.^{13,14}

Proteins that are weakly bound to the surface may be lost during the separation steps. They are referred to as the *soft corona* as opposed to the *hard corona* composed of tightly bound proteins that remain adsorbed after NP separation.²⁵ The reshuffling of the corona was described by Pisaní *et al.* who compared its composition following the separation of magnetic silica NPs by centrifugation or magnetization.²⁶ While the separation steps are known to alter the

composition of the protein corona (see the review by Weber *et al.* on the different separation techniques²⁷), most studies of protein adsorption are still performed following the classical approach of incubation and separation of NPs from the biological medium before protein analysis. One reason is that there are few techniques to analyse the protein corona *in situ*. A methodology based on flow cytometry and immunolabelling was developed to quantify the adsorption of a few specific proteins on NPs in human serum.²⁸ The analysis of the hydrodynamic radius (or gyration radius) of NPs coated by proteins using fluorescence correlation spectroscopy,²⁹ small-angle neutron scattering,³⁰ NMR^{31,32} or confocal laser scanning microscopy³³ can also provide information on the formation of the protein corona in solution. Here we introduce another approach based on cryoTEM and synchrotron-radiation circular dichroism (SRCD) to detect changes in protein structure and stability mediated by nanoparticles *in situ* (Figure 1).

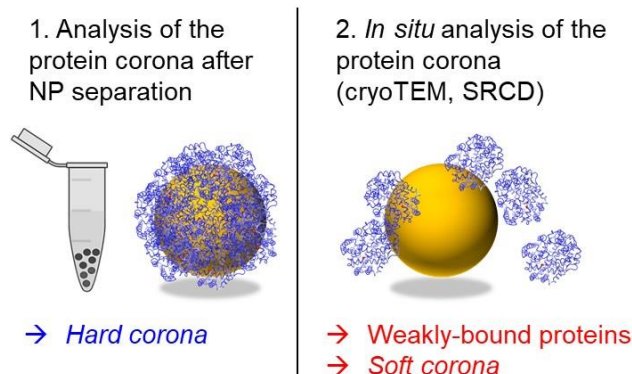


Figure 1. Schematic representation of the analysis of the hard corona following NP separation (1), and the *in situ* analysis of weakly-bound proteins and the soft corona by cryoTEM and SRCD (2).

Changes in protein structure and activity following adsorption on NPs (and more generally on inorganic surfaces) have been well documented in the literature.^{9,12,34,35} Based on these observations, it was hypothesized that NPs may also have a biological impact by affecting protein structure and activity, for example by triggering an immune response following modifications of protein conformation or by altering the activity of an enzyme.¹⁰ It implicitly assumes that biological effects of NPs on proteins would result from protein adsorption on NPs, as evidenced by the formation of a hard corona.

In this study, we show that the classical approach based on NP separation from the biological medium fails to detect protein destabilization mediated by weak protein-NP interactions, and therefore cannot be used to fully predict biological effects of nanomaterials *in situ*.

We investigated the effect of silica nanoparticles (SNPs) on the structure and stability of a model protein, oxyhemoglobin (oxyHb), by SRCD. OxyHb is the functional oxygenated form of hemoglobin. Monodisperse Ludox TM50 silica nanoparticles with a mean diameter of 26.0 nm^{30,36} were used to investigate their effect on the structure of oxyHb. Previously, we reported that oxyHb strongly adsorbs on SNPs at pH 7, but does not form a hard corona on SNPs at pH 9.^{35,37}

We used these two pH conditions to investigate the molecular mechanisms and the structural effects of strong *versus* weak protein-NP interactions respectively during oxyHb thermal unfolding.

RESULTS AND DISCUSSION

First, we checked for the stability of SNPs in the 3 buffers chosen for the SRCD experiments: 100 mM phosphate buffer pH 7, 5 mM ammonium acetate (NH₄Ac) buffer pH 9, and 3 mM N-Cyclohexyl-2-aminoethanesulfonic acid (CHES) buffer pH 9. Two buffer conditions were chosen at pH 9 to take into account potential ionic effects on protein-NP interactions. Experiments were conducted in phosphate buffer only at pH 7 because we did not observe any major buffer effect on the adsorption of oxyHb on SNPs at this pH.³⁵ There is no difference in size or aggregation state of SNPs at pH 7 nor at pH 9 in the different buffer conditions (Figure S1). The surface charge of SNPs is very similar at pH 7 and at pH 9, as indicated by the ζ -potential values: $\zeta = -42.1 \pm 2.2$ mV, $\zeta = -45.1 \pm 0.5$ mV and $\zeta = -47.3 \pm 1.3$ mV in phosphate buffer pH 7, in NH₄Ac buffer pH 9, and in CHES buffer pH 9 respectively. We did not observe any dissolution of SNPs at pH 9 over time (Figure S1).

OxyHb forms a hard corona on SNPs at pH 7

The adsorption isotherm of oxyHb on SNPs at pH 7 is shown in Figure 2a. The adsorption constant K_{ads} was calculated using the Langmuir model assuming reversible adsorption.³⁸ The Langmuir model relates the amount of adsorbed protein to the free protein concentration in solution for a fixed amount of binding sites that is for a fixed NP concentration (eq 1):

$$m_{ads} = \frac{m_{\infty} * K_{ads} * C_p}{1 + K_{ads} * C_p} \quad (\text{eq 1})$$

where C_p is the free protein concentration at equilibrium, m_{ads} the amount of adsorbed protein, m_{∞} the maximum amount of adsorbed protein. Here we measured the amount of adsorbed protein as a function of NP concentration, for a fixed protein concentration.

We adapted the Langmuir model to fit the adsorption isotherms (eq 2):

$$\%_{ads} = \frac{\left(\frac{n \times C_{NP} + C_{protein}}{K_{ads}} - \sqrt{\left(\frac{n \times C_{NP} + C_{protein}}{K_{ads}} \right)^2 - 4 \times n \times C_{NP} \times C_{protein}} \right)}{2 C_{protein}} \quad (\text{eq 2})$$

where $\%_{ads}$ is the percentage of adsorbed protein, C_{NP} the NP concentration (mol/L), and $C_{protein}$ the initial protein concentration. The quantity $n \times C_{NP}$ represents the concentration of binding sites, with n the number of binding sites per NP. The adsorption isotherms measured at pH 7 with two protein concentrations (0.5 mM and 1 mM oxyHb) are accommodated by the model (Figure 2a).

The adsorption constant $K_{ads} = 2.1 \times 10^5 \text{ M}^{-1}$ and the amount of adsorbed protein $m_{\infty} = 2.2 \text{ mg/m}^2$ reflect the high affinity of oxyHb for SNPs at pH 7. In the conditions used for the SRCD measurements ($C_{oxyHb} = 1 \text{ mM}$, $C_{SNP} = 100 \text{ mg/mL}$), 97% of oxyHb molecules are adsorbed on the SNPs at pH 7 at 22°C after NP centrifugation. These conditions were chosen so that the SRCD spectra of adsorbed oxyHb could be recorded *in situ* without any contribution from the free

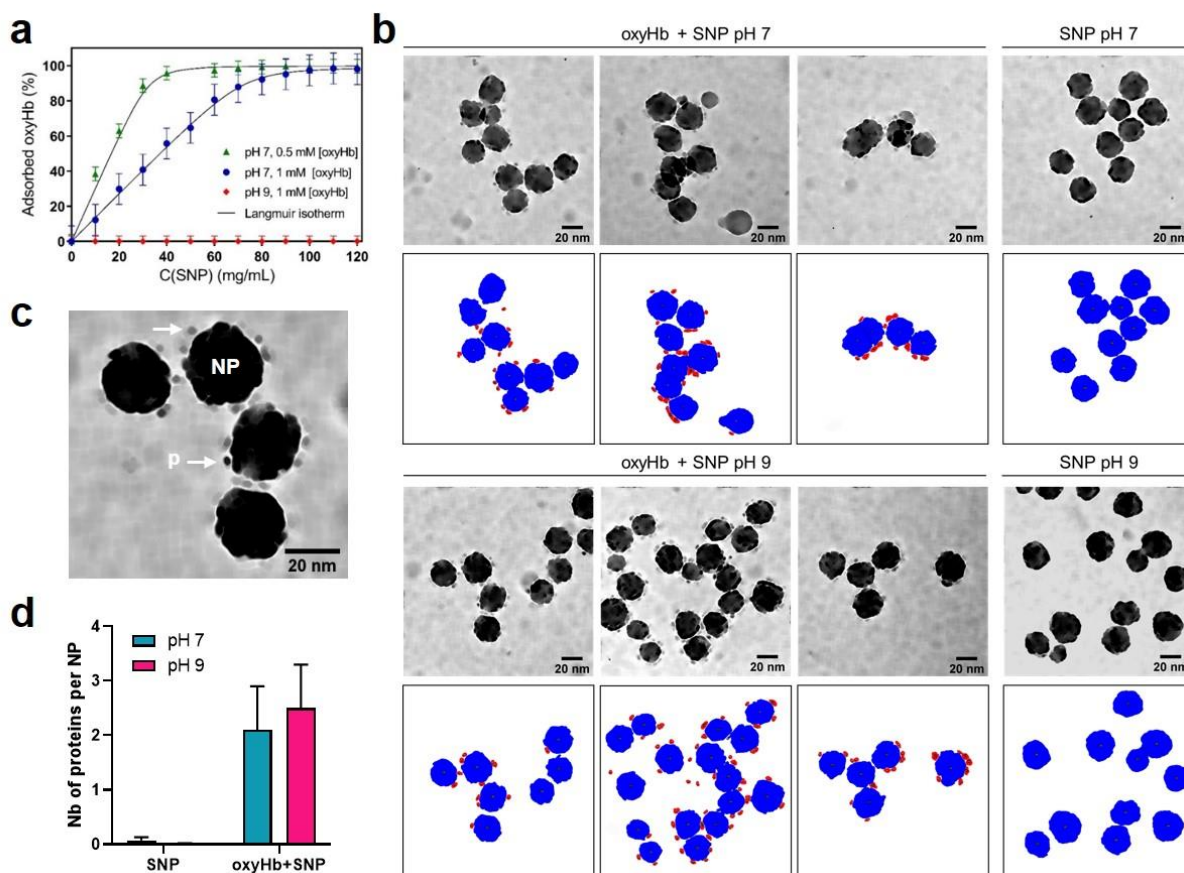


Figure 2. Identification of the hard corona and soft corona formed by oxyHb on SNPs. (a) Adsorption isotherms of oxyHb on SNPs in phosphate buffer at pH 7 and in NH_4Ac buffer at pH 9 at 22°C. $C(\text{oxyHb}) = 0.5 \text{ mM}$ and 1 mM . At pH 9, <3% oxyHb is adsorbed on SNPs after NP centrifugation. Fit by the Langmuir model is shown at pH 7 only. (b) CryoTEM images of SNPs with and without oxyHb at pH 7 and at pH 9. The results of automated particle analysis of SNPs (in blue) and adsorbed oxyHb (in red) are shown below. (c) Zoom on the cryoTEM image of SNPs (NP) with adsorbed oxyHb (p) at pH 9. Some adsorbed proteins are indicated by a white arrow. (d) Number of proteins adsorbed per SNP at pH 7 and at pH 9 identified by automated analysis of the cryoTEM images in 2D projection. A minimum of 100 SNPs were analyzed in each condition.

proteins in solution at pH 7. A hard corona of oxyHb forms on SNPs at pH 7.

OxyHb forms a soft corona on SNPs at pH 9

In contrast, no oxyHb (within the experimental error of 3%) was adsorbed on SNPs at pH 9 in NH_4Ac buffer or in CHES buffer following NP separation. To check whether this observation resulted from a lack of protein-NP interaction at pH 9 or from a loss of weakly-bound proteins during the separation step, we analysed the SNPs *in situ* by cryoTEM (Figure 2b, 2c). CryoTEM imaging was performed with an excess of proteins in solution to maximize corona formation. The presence of proteins at the surface of SNPs is clearly visible both in phosphate buffer at pH 7 and in NH_4Ac buffer at pH 9. The high resolution of the cryoTEM images makes it possible to identify individual nanoparticles and adsorbed proteins based on their size differences (Figure 2b). Automated image analysis was developed with ImageJ software for this purpose (see Methods, Table S1). Untreated images are shown in Figure S2. The same parameters were applied to all the images. Free proteins were excluded from the analysis. A minimum number of 100 nanoparticles was automatically analyzed in each condition.

The average diameter of SNPs was $25.6 \pm 0.6 \text{ nm}$, in agreement with published data.^{30,36} The average diameter of adsorbed proteins was calculated assuming a spherical shape. A diameter of $4.7 \pm 0.2 \text{ nm}$ was measured at pH 7 and at pH 9, close to the value reported for native porcine hemoglobin ($d = 4.6 \text{ nm}$). Interestingly, a SANS study showed that little modification of oxyHb shape occurs following adsorption onto SNPs.³⁰ The cryoTEM images show that oxyHb adsorbs on SNPs both at pH 7 and at pH 9 (Figure 2b, 2c, 2d). This observation is confirmed by the automated image analysis, which detected a minimum of 2 adsorbed proteins on SNPs both at pH 7 and at pH 9 for more than 100 nanoparticles analyzed (Figure S2). The number of proteins per NP (Figure 2d) reflects the number of adsorbed proteins detected automatically in a 2D projection. This indicator does not quantify the number of proteins in the corona, but provides a numerical value to evaluate the presence or absence of adsorbed proteins from the observation of a larger number of objects by cryoTEM. It confirms that oxyHb adsorbs onto SNPs both at pH 7 and at pH 9 *in situ*. The adsorption of oxyHb on SNPs at pH 9 was not seen in the adsorption isotherm (Figure 2a) suggesting that adsorbed oxyHb is lost following SNP centrifugation. These results show that oxyHb forms a soft corona of weakly

bound proteins on SNPs at pH 9, whereas it can form a hard corona that is retained after SNP centrifugation at pH 7. Therefore oxyHb can form either a hard or a soft corona on SNPs depending on pH conditions. Here, the formation of the soft corona was observed *in situ* by cryoTEM analysis.

In situ structural analysis of adsorbed proteins by SRCD

Following this observation, we investigated the effect of SNPs on oxyHb secondary structure in the hard corona or in the soft corona *in situ* at pH 7 and at pH 9 respectively by SRCD (Figure 3). The high photon fluxes provided by SRCD in the UV region can be used to access additional electronic transitions of the polypeptide backbone from 170 nm to 260 nm and to accurately detect changes in protein secondary structures in colloidal suspensions.^{39,40} The SRCD spectra of free oxyHb are identical at pH 7 and at pH 9 in the 3 buffers (Figure S3) showing that the average secondary structure of oxyHb does not vary between pH 7 and pH 9. The SRCD spectrum of oxyHb is characterized by two minima at 222 nm and 209 nm, and a maximum at 194 nm with a shoulder around 175 nm. These features are typical of proteins with a predominant α -helical secondary structure. The peaks originate from the $n \rightarrow \pi^*$ electronic backbone transition (222 nm) and the exciton splitting of the $\pi \rightarrow \pi^*$ transition (194 nm and 209 nm) which is specific to helical structures.⁴¹

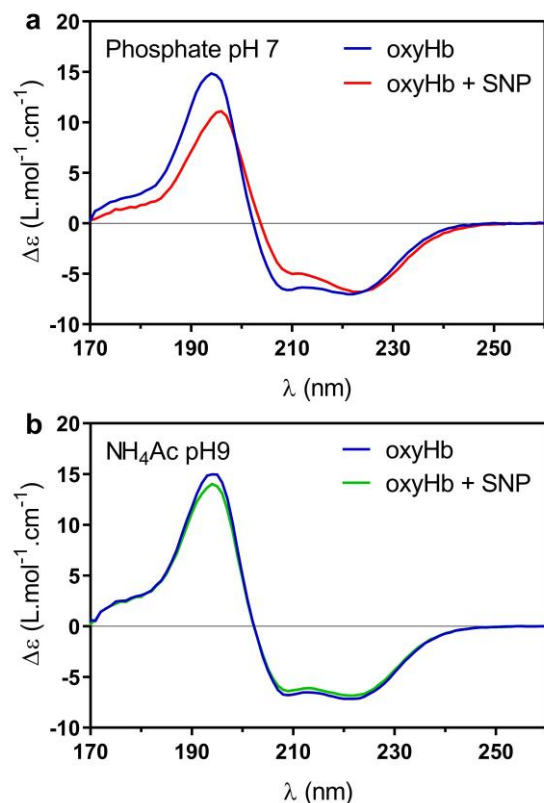


Figure 3. SRCD spectra of oxyHb with and without SNPs at 22°C (a) in phosphate buffer pH 7 and (b) in NH_4Ac buffer pH 9. $C(\text{oxyHb}) = 1 \text{ mM}$, $C(\text{SNP}) = 100 \text{ mg/mL}$.

The secondary structure of oxyHb was determined by deconvolution of the SRCD spectra with the BeStSel algorithm, showing good consistency between the experimental and the deconvoluted spectra (Figure S4).^{42,43} The structural analysis of oxyHb in solution was validated by comparing our results to the secondary structure calculated with DSSP for the crystal structure of porcine oxyHb 1QPW.pdb (Table S2). The DSSP program calculates the secondary structure of a protein from the atomic coordinates of its crystal structure.^{44,45} Good agreement was observed between DSSP calculation and SRCD spectra deconvolution.

The differences between the SRCD spectra of oxyHb with and without SNPs at pH 7 (Figure 3a) reflect the structural changes induced by the formation of the hard corona. The spectrum of adsorbed oxyHb is characterized by a decrease of the intensity at 209 nm, a shift of the maximum from 194 nm to 196 nm, whereas the peak at 222 nm is unaffected. The structural analysis reveals a relatively limited decrease of the helix content from 68% to 59% following adsorption. The $\Delta\epsilon(222\text{nm})/\Delta\epsilon(209\text{nm})$ ratio (corresponding to the two minima of the SRCD spectra of oxyHb) is equal to 1.04 for free oxyHb and increases to 1.38 for adsorbed oxyHb at pH 7. This ratio is indicative of the intra- and inter-molecular interactions between helices in oligomeric proteins.⁴⁶ MD simulations carried out by Jiang *et al.* on tetrameric hemoglobin suggest that the electronic transition at 208 nm is stronger at helix termini due to the parallel orientation of the polarisation to the helix axis, whereas the electronic transition at 222 nm is homogeneously distributed within the protein structure, with a polarisation perpendicular to the helix axis.⁴⁷ The decrease of the ellipticity of adsorbed oxyHb at 209 nm and the unchanged ellipticity at 222 nm could suggest preferential alteration or reorientation of helix termini following adsorption on SNPs. Differential structural modifications between the α and β subunits following adsorption could, for example, play a role.

At pH 9, we did not observe any differences between the SRCD spectra of oxyHb with and without SNPs at 22°C in NH_4Ac buffer (Figure 3b) or in CHES buffer (Figure S4). Therefore there is no effect of SNPs on the secondary structure of weakly bound oxyHb forming the soft corona at 22°C. This result also confirms that there is no artefact due to light scattering by SNPs on the SRCD spectrum of proteins in the far UV range.

Then, we investigated the effect of SNPs on the structure and stability of oxyHb in the hard and soft corona during thermal unfolding. We measured the SRCD spectra of oxyHb with and without SNPs at pH 7 and at pH 9 from 22°C to 97°C (Figure 4, Figure S5). The melting curves are displayed as $\Delta\epsilon(194\text{nm}) = f(T)$ to compare the unfolding pathways of oxyHb with and without SNPs (Figure 5a). The thermal unfolding of oxyHb is characterized by a moderate decrease of the ellipticity between 20°C and 60°C followed by a sharp decrease between 60°C and 85°C. The melting curves were best fitted by two-sigmoidal curves (Table 1).

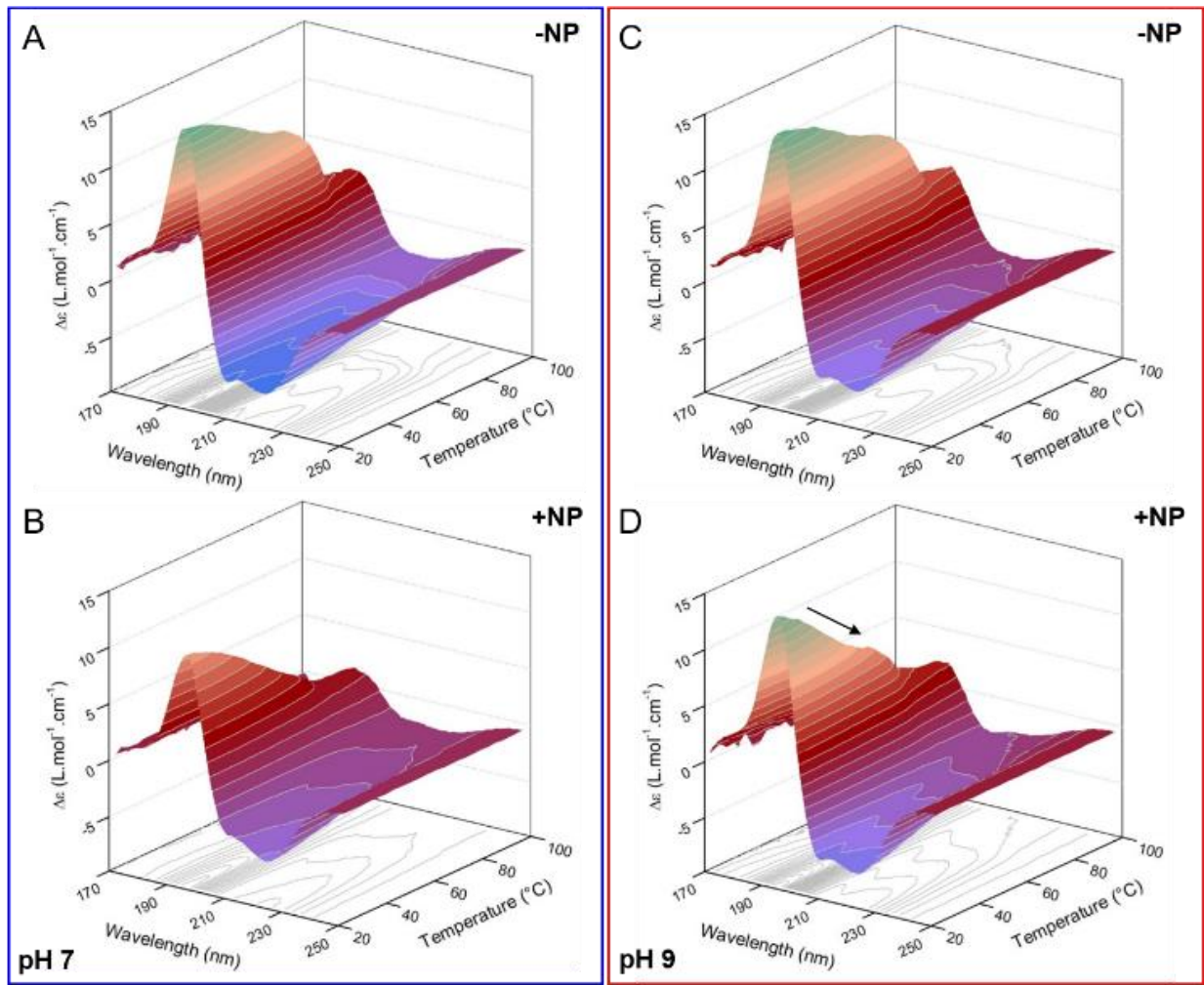


Figure 4. Thermal unfolding of oxyHb with and without SNPs (A,B) in phosphate buffer pH 7, (C,D) in NH_4Ac buffer pH 9. The SRCD spectra were recorded from 22°C to 97°C. The arrow in D highlights the difference between the SRCD spectra of oxyHb with and without SNPs around physiological temperature at pH 9.

The first phase centered on $T_{m1} = 47.0 \pm 0.6$ °C corresponds to the initiation of oxyHb unfolding at moderate temperature. It is followed by irreversible unfolding during the second phase centered on $T_{m2} = 77.4 \pm 0.5$ °C. Similar results were obtained by fitting the melting curves $\Delta\epsilon(222\text{nm}) = f(T)$ showing good consistency in the analysis of the different polypeptide backbone electronic transitions of the SRCD spectrum (Figure S6, Table S3). The melting curves of oxyHb are very similar in phosphate buffer at pH 7, in NH_4Ac buffer at pH 9 and in CHES buffer at pH 9. The close values of the melting temperatures of oxyHb at pH 7 and at pH 9 suggest that the unfolding pathway of oxyHb is similar in both pH conditions (Table 1).

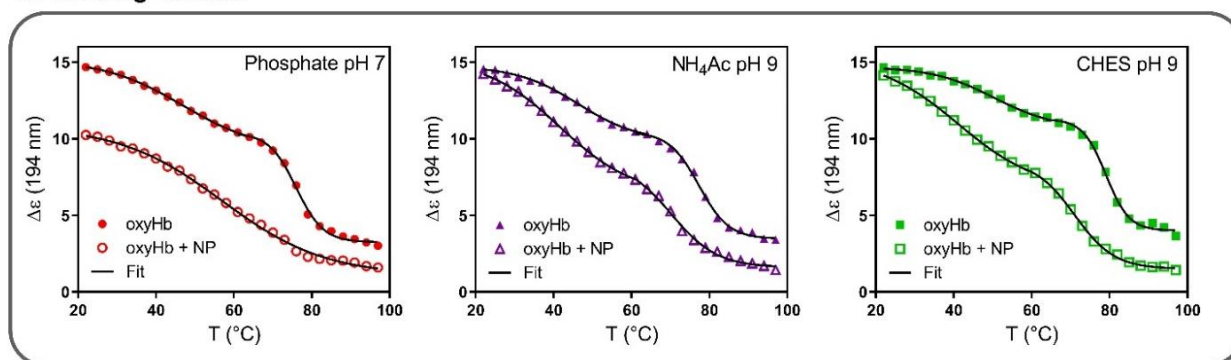
This two-step process was also observed by Yan *et al.* using 2D FTIR spectroscopy to follow the unfolding and the aggregation of hemoglobin with temperature.⁴⁸ The authors describe Hb thermal unfolding as a two-phase process with an initial structural perturbation stage (30°C–50°C) characterized by the partial unfolding of the α -helices associated with higher solvent exposure; and a thermal aggregation stage (50°C–80°C) dominated by the unfolding of buried structures and the formation of aggregates.

Table 1. Melting temperatures T_m of oxyHb with and without SNPs obtained by fitting the melting curves $\Delta\epsilon(194 \text{ nm}) = f(T)$ with one (*) or two sigmoids. The experimental error related to the temperature measurement (0.1 °C) was added to the curve fitting error.

Buffer	oxyHb \pm SNP	T_{m1} (°C)	T_{m2} (°C)
Phosphate pH 7	oxyHb	47.0 ± 0.6	76.0 ± 0.4
	oxyHb + SNP	-	$57.0 \pm 0.7^*$
NH_4Ac pH 9	oxyHb	46.3 ± 0.6	77.4 ± 0.5
	oxyHb + SNP	41.7 ± 0.4	70.2 ± 0.7
CHES pH 9	oxyHb	50.4 ± 1.2	79.1 ± 0.7
	oxyHb + SNP	41.2 ± 0.5	70.6 ± 0.6

The structural analysis of the SRCD spectra supports this sequence of events (Figure 5b). We observed a first unfolding phase (20°C–60°C) characterized by a moderate decrease of the helical content and an increase of the disordered content of oxyHb (included in the ‘others’ section), and a second phase (60°C–80°C) wherein a larger portion

A. Melting curves



B. Secondary structure deconvolution

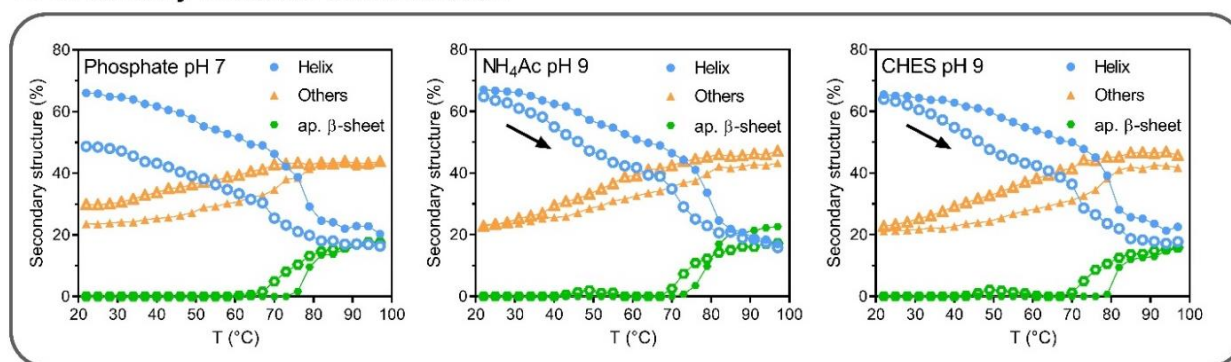


Figure 5. (A) Melting curves of oxyHb with and without SNPs expressed as $\Delta\epsilon(194 \text{ nm})$ in phosphate buffer pH 7, in NH_4Ac buffer pH 9 and in CHES buffer pH 9. Curves are fitted with two-sigmoids, except for oxyHb with SNPs at pH 7, which is best fitted with one sigmoid. (B) Evolution of the helix, the antiparallel β -sheet (ap. β -sheet) content and ‘others’ (which includes the disordered fraction) of oxyHb with SNPs (open symbols) and without SNPs (closed symbols). Turns which show only minor changes and represent $<10\%$ of the secondary structure (Table S1) are not included in the figure for clarity. The arrows highlight the partial loss of helices in oxyHb with SNPs at physiological temperature at pH 9.

of the helices are unfolded and antiparallel β -sheets are formed. Since oxyHb does not contain any β -sheet in its native structure, we assume that the antiparallel β -sheets formed at high temperature correspond to intermolecular β -sheets associated with protein aggregation.⁴⁸ The observation of the aggregation of free oxyHb by UV-vis spectroscopy at 60°C supports this hypothesis (data not shown).

Interestingly, a single-step process best describes the unfolding pathway of adsorbed oxyHb in the hard corona at pH 7. The melting curve $\Delta\epsilon(194 \text{ nm}) = f(T)$ is fitted by one sigmoidal curve centered on $T_{m2} = 57.0 \pm 0.7^\circ\text{C}$ (Figure 5A). The large decrease of the melting temperature shows that the adsorption of oxyHb on SNPs greatly impairs the thermal stability of the protein. This effect can be rationalized by the partial structural loss associated with the initial adsorption step which further reduces the enthalpic cost of protein unfolding with temperature.^{12,49,50}

SNPs alter the stability of weakly-bound proteins

At pH 9, the evolution of the SRCD spectra of weakly-bound oxyHb forming the soft corona revealed a completely different picture. Similarly to the thermal unfolding of free oxyHb, a two-step unfolding pathway is maintained. However, we observed a large decrease of $\Delta\epsilon(194 \text{ nm})$ with

SNPs which gradually builds up with temperature. Identical effects were observed for oxyHb unfolding with SNPs in NH_4Ac and in CHES buffers. The structural analysis shows that SNPs induce a partial loss of the helical structure of oxyHb and an increase of the disordered structure at moderate temperature (Figure 5B). At 40°C , the helical content decreases to 55% in both buffer conditions (Table 2). Moreover, the decrease of T_{m1} and T_{m2} suggests that SNPs undermine the thermal stability of oxyHb at pH 9 both at moderate and high temperature.

Table 2. Percentage of oxyHb adsorbed on SNPs (top line) and difference in the helical content of oxyHb with and without SNPs (*below, in italics*) at 22°C , 40°C , 50°C in phosphate buffer pH 7, in NH_4Ac buffer pH 9, and in CHES buffer pH 9. *At 50°C , oxyHb starts aggregating in solution.

T ($^\circ\text{C}$)	Phosphate pH 7	NH_4Ac pH 9	CHES pH 9
22	91.2 \pm 0.2 % - 17.3 %	< 3 % - 2.3 %	< 3 % - 1.6 %
40	87.7 \pm 1.1 % - 18.4 %	< 3 % - 7.5 %	< 3 % - 8.0 %
50 *	85.7 \pm 0.7 % - 18.5 %	< 12.2 \pm 3.9 % - 10.0 %	< 7.0 \pm 3.7 % - 12.3 %

We focused our analysis on the moderate temperature range (20°C–49°C) to investigate the destabilization effect of SNPs on oxyHb within the hard or soft corona before the initiation of aggregation. The SRCD spectra of oxyHb with and without SNPs at 22°C, 37°C and 49°C are shown in Figure S7. We can see that the secondary structure of oxyHb is already altered by SNPs at pH 9 at physiological temperature with an additional loss of 8% of the helix content in the presence of SNPs. These results show that SNPs mediate the destabilization of weakly bound oxyHb within the soft corona at physiological temperature.

We observed that SNPs alter oxyHb secondary structure for $T > 30^\circ\text{C}$. To check whether this effect could be due to the formation of a hard corona at higher temperature, we measured the amount of oxyHb adsorbed on SNPs from 22°C to 50°C (Table 2). Less than 3% of oxyHb adsorbs on SNPs at 40°C, at pH 9 in NH_4Ac buffer and in CHES buffer, while an additional loss of 8% of the helical content is observed within oxyHb structure. This result shows that SNPs can destabilize oxyHb at moderate temperature even in the soft corona. At 50°C, only 10% oxyHb remains adsorbed on SNPs after NP separation. Adsorption could be favoured by partial unfolding of oxyHb at this temperature. However protein unfolding also leads to the onset of aggregation of the free oxyHb in solution. As the amount of adsorbed oxyHb is measured by the depletion method after centrifugation of SNPs, the aggregation of oxyHb could artificially increase this value. This hypothesis is supported by the concomitant decrease of the amount of adsorbed oxyHb on SNPs at pH 7.

OxyHb is neutral around pH 7 ($\text{pI} = 7.2$ for porcine oxyHb⁵¹) and globally negatively charged at pH 9, whereas the SNPs are negatively charged at pH 7 and at pH 9 with similar surface charge in the 3 buffer conditions. Ionic strength can modulate protein adsorption on SNPs,^{50,52} but we did not observe differences between NH_4Ac and CHES buffer at pH 9. The formation of a hard or soft corona of oxyHb on SNPs at pH 7 and at pH 9 respectively could primarily be explained by the difference in protein charge. Besides, the secondary structure of oxyHb is very similar in phosphate buffer pH 7, in NH_4Ac buffer pH 9 and in CHES buffer pH 9 at 22°C. Therefore we hypothesize that the change of oxyHb affinity for SNPs from pH 7 (high affinity) to pH 9 (low affinity) is mostly due to the deprotonation of basic residues in this pH range and to the electrostatic repulsion between negatively charged oxyHb and the anionic siloxide groups at the surface of SNPs.^{13,14,52}

Molecular basis of soft corona formation

To gain insights into the structural changes of oxyHb as a function of pH and temperature, we performed molecular dynamic (MD) simulations of the oxyHb structure at pH 7 and at pH 9 at 295 K (22°C), 322 K (49°C), 353 K (80°C) and 400 K (127°C) (Figure 6). As charge plays a key role in protein-NP interactions in our system, first we calculated the protonation state of oxyHb at pH 7 and at pH 9. Each oxyHb α and β chain contains 3 Arg, 11 Lys, 11 His, and 5 Arg, 11 Lys, 8 His respectively. His side chain has a pK_a around 6. Three protonation states of His were considered: HID,

HIE, and HIP, which correspond to the protonation of only ND1 (HID), only ND2 (HIE), or both ND1 and ND2 (HIP) (Figure S8). At pH 9, His is neutral and can only adopt one of the two tautomer forms, HID or HIE. At pH 7, His can also be in the protonated HIP form and carry a global charge of +1. We used two independent programs, Proton33,54 and Propka^{55,56}, to evaluate the protonation state of His in oxyHb structure at pH 7 and at pH 9 (see Methods).

Among the 38 His residues, only 3 were predicted to be protonated: His50, His112 in α subunit, and His120 in β subunit. Accordingly we assigned the HIP form (charge +1) to these 3 His at pH 7. For the other His residues, we assigned the tautomer form (HIE or HID) predicted by Propka 3.1 in the presence of ligands (Table S4). This program checks the local environment of each His to predict the tautomer form which is most likely.

We followed the same approach to define the protonation state of His residues in oxyHb structure at pH 9. In this case, only HID and HIE protonation states were considered (Figure S9). After assigning the calculated His protonation state to oxyHb structure at pH 7 and pH 9, we simulated the evolution of the protein structure in both pH conditions at 4 different temperatures by MD simulation (Figure 6). The fluctuations which appear on the RMSF (Root Mean Square Fluctuations) graphs are reported on the 3D structure of oxyHb. The most mobile regions (in red) are visible in the protein structure. When temperature increases, the regions showing the largest fluctuations remain the same. The existing mobile substructures become larger, without any appearance of new mobile element.

Note that the simulation at 400 K is used here to emphasize the fluctuations in the protein structure and to check whether new mobile elements appear at high temperature. MD simulations minimize protein unfolding, that is structural modifications are usually underestimated at the target temperature, unless much longer simulation times are used.⁵⁷ This explains why the simulated structure appears relatively preserved at a target temperature above the real unfolding temperature of oxyHb.

The most mobile regions are similar at pH 7 and pH 9. Interestingly, the intensities of the fluctuations are more homogeneous within the oxyHb structure at pH 7. For the 2 systems, 2 regions are particularly altered with temperature: the helix-loop-helix motif (residues 43 to 56 in α subunit), and the helix motif (residues 218 to 236 in β subunit). Even if the amino acid sequences of α and β subunits differ, their folding is the same.

The structures of these 2 motifs are shown at 295 K, 322 K, and 400 K in Figure 6. We can see that the hydrogen bond between Asn47 and Ser52 that closes the loop is broken at high temperature in the helix-loop-helix structure. Polar amino acids become exposed, in particular Ser49, Ser51, and His50, which is protonated at pH 7. These residues can then interact with a ligand through molecular interactions. This mechanism is also reflected by the decrease of the percentage of H bonds within this motif when temperature increases. The secondary structure of the second helix motif

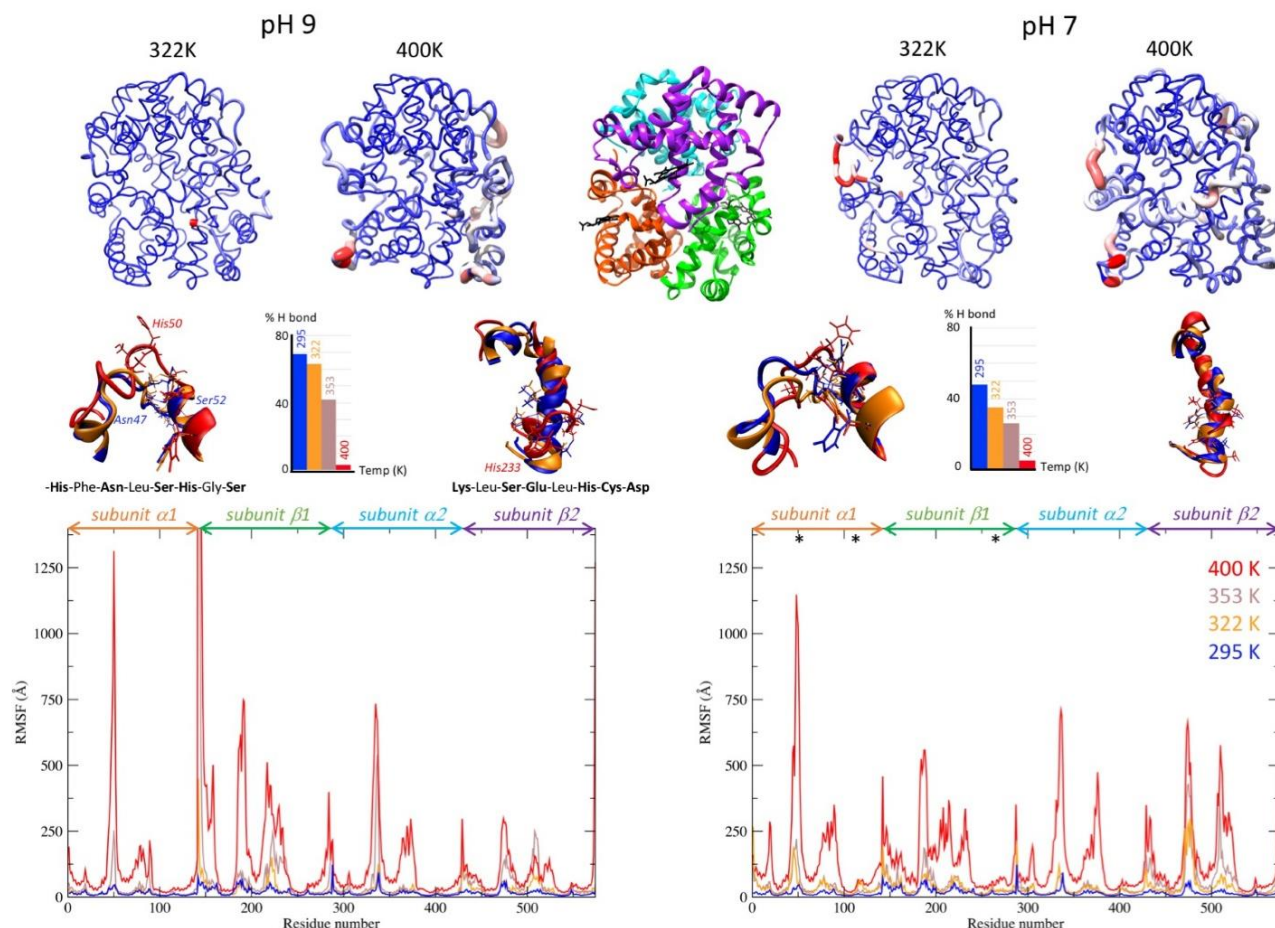


Figure 6. MD simulations of oxyHb structure at pH 7 and pH 9 as a function of temperature. Structures obtained at the end of the simulation at 322 K and 400 K are represented at the top (at pH 9 on the left, at pH 7 on the right) and coloured according to the RMSF of each residue (small fluctuations in blue, large fluctuations in red). The same colour gradient was applied to all the structures. Each subunit of the tetrameric structure is represented in cartoon mode in a different colour at the top centre of the figure with the 4 hemes represented in black licorice mode. The limits between each subunit are reported on the RMSF graphs below. The RMSF of the backbone atoms is shown along oxyHb sequence at 295 K (blue), 322 K (orange), 353 K (brown), and 400 K (red). * indicates HIP histidine with a +1 charge at pH 7. The inset structures represent the helix-turn-helix motif centered on His50 and the helix motif from Leu222 and Asp235, at 295 K, 322 K, and 400 K. The corresponding sequences of the turn and helix motifs that break with temperature are displayed. Polar or charged residues area highlighted in bold. The histogram represents the percentage of hydrogen bonds between Asn47 and Ser52 closing the turn at each temperature.

is totally ruptured at 353 K. This also leads to the exposure of polar and charged residues, notably His233.

These results suggest that the exposure of polar and charged amino acids following the disruption of these 2 motifs with temperature, and in particular Asn47, Ser52, His 50, and His233, may drive the interactions of oxyHb with SNPs. This hypothesis is supported by the increase in the surface accessibility of oxyHb with temperature at pH 7 and at pH 9 (Figure 7a, Figure S10), which could favour interactions with the silica surface. The increase in the solvent-accessible surface area (SASA) of oxyHb is explained here by two factors: -the increase in protein flexibility as expressed by its B-factor; -the loss of some structural elements exposing some residues to the solvent, such as the 2 motifs described here. The sequence 43-56, in particular, shows a marked increase in the residue SASA with temperature at pH 9 (Figure S11).

To determine if this mobile region can interact with the silica surface, we calculated the interaction energy between the corresponding peptide and the silica surface.

The motif structures at 295 and at 400 K (hereafter called P295 and P400) were extracted from the MD simulation of the oxyHb structure (Figure 6). The peptide sequence is Phe43-Pro44-His45-Phe46-Asn47-Leu48-Ser49-His50-Gly51-Ser52-Asp53-Gln54-Val55-Lys56. The terminal -CO and -NH groups of the sequence were replaced by H atoms. The interaction of P295 and P400 with the silica surface was analyzed at the quantum chemistry level using the Dispersion-Corrected Self Consistent Charge Density Functional Tight Binding (DC-SCC-DFTB), which is an approximated DFT method adapted to the system size (445 atoms).⁵⁸ The silica surface was represented by a silica slab of 17 x 17 x 5 Å along the 001 plane composed of 42 SiO₄ units with 68 OH surface groups. The N and C atoms of the two first and last residues were frozen. The peptide was positioned 2 Å away from the silica slab at the beginning of the calculation, with the most mobile region corresponding to residues 50-56 oriented towards the surface. The minimized structures of P295 and P400 in interaction with the silica surface are shown in Figure 7.

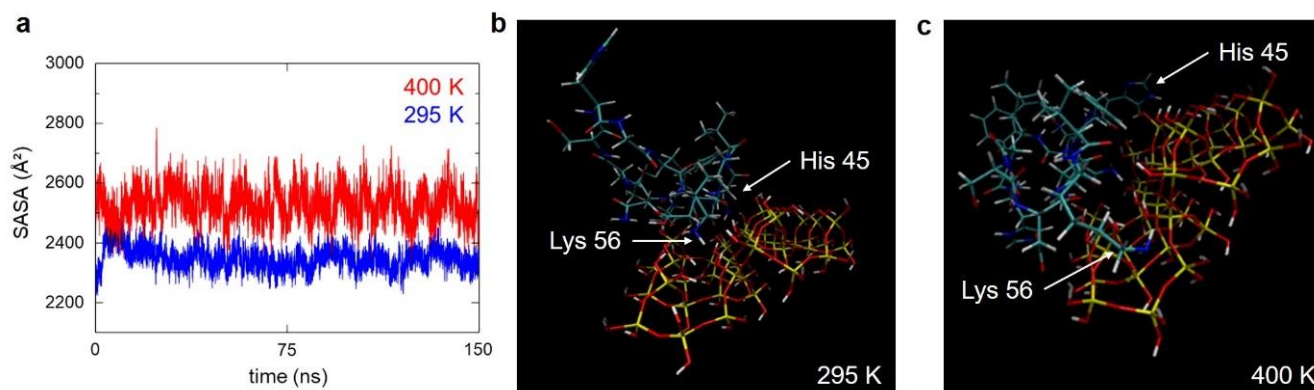


Figure 7. Surface accessibility of oxyHb and interaction of the mobile region with the silica surface at pH 9 as a function of temperature. (a) Calculated solvent-accessible surface area (SASA) of oxyHb at pH 9 at 295 K (in blue) and at 400 K (in red). The calculation time is 150 ns. (b,c) DC-SCC-DFTB calculations of the interaction of peptide 43-56 with the silica surface starting from the structure at 295 K (b) and from the structure at 400 K (c). The peptide sequence is Phe43-Pro44-His45-Phe46-Asn47-Leu48-Ser49-His50-Gly51-Ser52-Asp53-Gln54-Val55-Lys56. The peptide structure was extracted from MD simulation of oxyHb structure at pH 9 at 295 K and 400 K respectively. Residues Lys56 and His45 are shown by white arrows in the minimized structure. The silica surface is represented by a silica slab of 17 x 17 x 5 Å along the 001 plane. The system is composed of 445 atoms with Si (yellow), O (red), H (white), C (grey), N (blue).

The interaction energy between P294 and P400 with the silica surface was calculated by subtracting the energy of the peptide and the silica slab alone minimized at the same level of theory. The interaction energy with the silica surface was -3.3 eV for P295 and -4.2 eV for P400. This result shows that the interaction with the silica surface is much more favourable for the denatured peptide.

The positioning of P295 and P400 on the silica surface (Figures 7b and 7c) suggests that this favourable interaction comes from a larger contact area of the denatured peptide with the surface, compared with the native peptide. Interestingly, the interaction with the silanol groups does not derive solely from basic amino acids (here Lys56, His45), but also from the carbonyl group of the peptide bond.

To test if the increase in the contact area between the peptide and silica was responsible for the favourable interaction, we calculated the interaction energy of P295 and P400 with a silica cluster composed of 20 SiO₄ units that provides a more limited surface area (average diameter 7 Å (Figure S12)). The same parameters and initial position of the peptides were applied. In this case, similar interaction energies were calculated for P295 and P400 (-0.58 eV for P295 and -0.59 eV for P400). These data confirm that the higher interaction energy of P400 with the silica surface is not due to stronger interactions of individual amino acids, but to a higher number of interactions or “contact points”. We conclude that the increase in surface accessibility and the partially unfolded state of this motif favours its interaction with the silica surface.

We observed that partial unfolding of oxyHb with temperature favours the interactions of weakly bound proteins with SNPs at pH 9, as shown by the destabilization effect observed by SRCD. A lower number of “contact points” (e.g. following His exposure with temperature) could account for the weak interaction between oxyHb and SNPs at pH 9 leading to the formation of a soft corona, as opposed to pH 7 where strong adsorption occurs at 22°C between native oxyHb and SNPs.

More generally, the folded structure of a protein results from the equilibrium between stabilizing and destabilizing forces, such as the formation of non-covalent interactions and the entropic contributions of both the protein and the solvent. However, the interplay of molecular forces that determine protein stability in solution may differ when a protein interacts with a surface.⁵⁹ We observed nanoparticle-mediated destabilization of weakly-bound proteins at moderate temperature during the first phase of oxyHb unfolding, that is when an equilibrium is formed between folded and partially unfolded conformations.⁶⁰ The intrinsic affinity of a protein for a surface is not a fixed parameter and changes in protein structure can initiate protein adsorption. For example, GFP (green fluorescent protein) does not exhibit any affinity for SNPs in its native form, and adsorbs spontaneously following protein unfolding.⁶¹ The conformational changes of Complement protein C3 following its proteolytic activation promote the binding of the C3b fragment to the surface of pathogens or nanomaterials which triggers the Complement cascade.^{62,63} It is also known that denatured proteins have a higher tendency to stick to surfaces. This effect is rationalized by the exposure of hydrophobic regions, the partial loss of structure and the higher flexibility of unfolded proteins that favour protein-surface interactions.⁵⁹

Interestingly, recent studies reported that even slight changes of the incubation temperature of NPs with plasma or serum alter the composition of the protein corona in the 20°C-50°C temperature range.^{64,65} Studies also reported a decreased affinity of some proteins for NPs at moderate temperature.⁶⁶ These observations suggest that structural events predating irreversible protein unfolding can alter the adsorption profile of a protein in a temperature-dependent way within a complex biological medium composed of hundreds of different proteins such as human serum. However, the wider impact of moderate temperature on the protein-NP interactions and the associated structural changes remain unclear.

SNPs stabilize partially unfolded conformations

We hypothesize that partially unfolded conformations of oxyHb exhibit a higher affinity for SNPs, which shifts the unfolded-folded equilibrium towards the unfolded state (Figure 8). The affinity of a protein for one NP is expressed by K_{ads} determined by fitting the adsorption isotherm of the protein with the Langmuir model (or other adsorption models). K_{ads} is calculated from the amount of adsorbed protein measured after NP separation from the biological medium. Therefore, K_{ads} only reflects the affinity of tightly bound high-affinity proteins forming the hard corona. This parameter does not take into account proteins weakly or transiently bound to NPs, which are lost during the separation step.

This concept is easily understood when one considers an assemblage of various proteins interacting with NPs. Here we suggest that a similar mechanism can drive protein-NP interaction wherein different conformational states in equilibrium exhibit different affinities for NPs. Our experimental results clearly show that SNPs destabilize oxyHb at pH 9. Because this interaction does not lead to the formation of a stable layer of adsorbed proteins on SNPs (hard corona), the classical approach based on NP separation from the biological medium fails to consider this mechanism. This is due in part to the low concentration of partially unfolded states which cannot be detected using classical methodologies.

The stability of each conformation is defined by its free energy $\Delta G(T)$. We determined the free energy of unfolding ΔG_u of oxyHb with and without SNPs to investigate the mechanisms leading to the destabilization of weakly bound proteins. The decrease of oxyHb stability can result either from a decrease of the enthalpy of unfolding (ΔH_u) or from an increase of the entropy of unfolding (ΔS_u) (eq. 3):

$$\Delta G_u(T) = \Delta H_u - T\Delta S_u \quad (\text{eq. 3})$$

ΔH_u and ΔS_u are usually calculated using a two-state model $F \leftrightarrow U$ by applying the Van't Hoff law to the CD data on the full temperature range (assuming ΔC_p is constant).⁶⁷ However, equilibrium thermodynamics analysis can be applied only if protein unfolding is reversible, which is no longer the case when protein aggregation occurs simultaneously.⁶⁸ We saw that oxyHb thermal unfolding is characterized by two thermal transitions T_{m1} and T_{m2} associated to a first partial unfolding phase without aggregation (20°C-50°C), and a second unfolding phase where aggregation and unfolding happens simultaneously (50°C-80°C). We also showed that oxyHb does not form a hard corona on SNPs during the first unfolding phase at pH 9. Therefore, we focused our analysis on the first unfolding phase to determine the thermodynamic parameters of oxyHb partial unfolding with and without SNPs before the initiation of aggregation, and close to physiological temperature.

We defined $\Delta G_u'$, $\Delta H_u'$ and $\Delta S_u'$ as the equilibrium thermodynamic parameters associated with the partial unfolding of oxyHb during the first unfolding phase. The corresponding midpoint temperature of unfolding is T_{m1} so $\Delta G_u'(T_{m1})$

= 0. The fraction of partially unfolded oxyHb $f(T)$ was calculated from the helix content of oxyHb A(T) (eq 4).

$$f(T) = \frac{A_n - A(T)}{A_n - A_u} \quad (\text{eq. 4})$$

where A(T) is the helix content of oxyHb as a function of temperature, A_n the helix content of native oxyHb at 22°C and A_u the helix content of partially unfolded oxyHb at the end of the first unfolding phase. The evolution of $f(T)$ is different with and without SNPs (Figure S13) and seems independent on the pH and buffer conditions. The constant of partial unfolding K_u' is defined as (eq 5):

$$K_u' = \frac{f(T)}{1-f(T)} \quad (\text{eq. 5})$$

The thermodynamic parameters of oxyHb partial unfolding during the first unfolding phase were obtained by fitting $K_u'(T)$ with the Van't Hoff law (eq 6):

$$\ln K_u' = -\frac{\Delta H_u'}{RT} + \frac{\Delta S_u'}{R} \quad (\text{eq. 6})$$

The Van't Hoff plots are shown in Figure S14. The thermodynamic parameters of oxyHb partial unfolding are summarized in Table S6. Then, the variation of $\Delta G_u'(T)$ was calculated for each condition (Figure S15). SNPs induce a decrease of $\Delta G_u'(T)$ of oxyHb at pH 9 on the full temperature interval of the first unfolding phase, that is a decrease of protein stability. Based on the decrease of the melting temperature and the preferential interaction of unfolded proteins with SNPs (Figure 7c),⁶¹ we hypothesize that this reduction in stability (decrease of $\Delta G_u'(T)$) is associated with an enhanced stability of the partially unfolded conformations forming the soft corona (formally a decrease of the Gibbs energy of the unfolded state).⁶⁹ The differences in enthalpy and entropy of oxyHb partial unfolding with and without SNPs were expressed as $\Delta\Delta H_u'$ and $\Delta\Delta S_u'$ where $\Delta\Delta X_u' = \Delta X_u'(\text{oxyHb+SNP}) - \Delta X_u'(\text{oxyHb})$ for $X=H,S$ (Table S7). In NH_4Ac buffer, both $\Delta\Delta H_u'$ and $\Delta\Delta S_u'$ are negative ($\Delta\Delta H_u' = -4.0 \text{ kJ.mol}^{-1}$, $\Delta\Delta S_u' = -6.3 \text{ J.mol}^{-1}.\text{K}^{-1}$). These results show that the destabilization effect of SNPs on weakly-bound proteins forming the soft corona is enthalpy-driven.

This analysis strongly supports our hypothesis that the destabilization effect of SNPs is due to molecular interactions that stabilize partially unfolded conformations in the soft corona and shift the folded \leftrightarrow unfolded equilibrium towards the unfolded states (Figure 8). This effect primarily comes from the range of affinity of different protein conformations for the NP surface when weak protein-NP interactions take place. As temperature increases, molecular interactions between partially unfolded conformations and NPs potentiate oxyHb thermal unfolding.

The free energy of protein unfolding can be affected by a high concentration of macromolecules or solutes, such as polymers, sugars, or urea, which can stabilize or destabilize proteins by different mechanisms by excluded volume effect, preferential hydration, or preferential binding.^{70,71} Based on the analysis of ubiquitin unfolding, Senske *et al.* classified cosolute effects in terms of their enthalpic contribution to protein stability.⁷⁰ According to this classification, a negative $\Delta\Delta H_u$ and a negative ΔT_m as observed for the unfolding of oxyHb with SNPs at pH 9 is associated

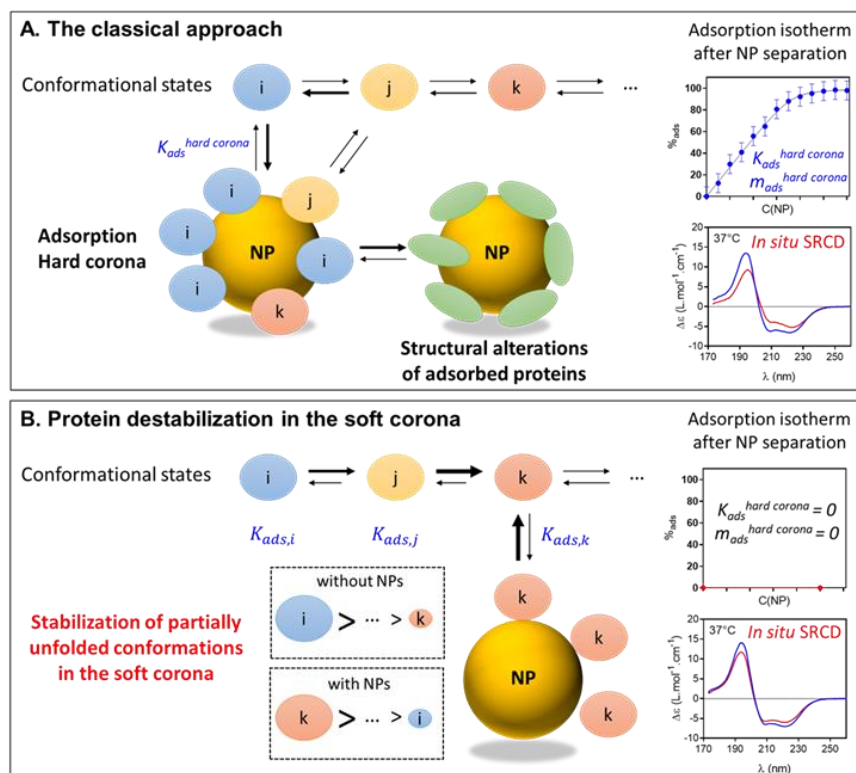


Figure 8. Nanoparticle-mediated protein destabilization. (A) In the classical approach, NPs are separated from the biological medium before protein analysis. The adsorption constant and the structural analysis of adsorbed proteins solely reflects the high affinity and the structural alterations of proteins forming the hard corona. (B) Destabilization of weakly-bound proteins in the soft corona. Nanoparticles stabilize partially unfolded conformations in the soft corona and shift the folded \leftrightarrow unfolded equilibrium towards the unfolded states. Protein destabilization remains unseen in the classical approach, but is visible with *in situ* analytical techniques such as SRCD. Representative adsorption isotherms and SRCD spectra of oxyHb forming a hard corona (top panel) or a soft corona (bottom panel) are shown. The reference spectra of oxyHb without SNPs are shown in blue for each condition.

with preferential binding of the solute on the accessible protein surface.⁷⁰

The extent of the destabilizing effect of SNPs on oxyHb ($\Delta\Delta H_u' = -4.0\text{ kJ}\cdot\text{mol}^{-1}$) is of the same order of magnitude (a few $\text{kJ}\cdot\text{mol}^{-1}$) as the data reported by Kumar *et al.* and Christiansen *et al.* for enthalpy-driven and entropy-driven macromolecular crowding effects respectively.^{71,72}

Whether the destabilization effects of SNPs on oxyHb and the formation of the soft corona could be described as enthalpy-driven macromolecular crowding effect remains to be elucidated. However, we can advance that nanoparticles are as diverse as other solutes in terms of their effects on protein stabilization or destabilization. Indeed, if we measured here an enthalpic destabilization of weakly bound proteins, Ortega *et al.* observed an entropic stabilization of protein L covalently bound to a gold surface.⁵⁹

This study was carried out on a model protein to allow the *in situ* analysis of the structural modifications associated with soft corona formation. The structural analysis of multiple proteins in a biological environment remains challenging. Studies suggested that the composition of the hard and soft corona differ from each other in the biological medium.^{24,73} This observation shows that both strong and weak interactions will drive the formation of the biomolecular corona *in vivo* and highlights the potential effect of weak protein-surface interactions on the biological effects of nanomaterials *in situ*.⁷⁴

Finally, we discuss the effects of affinity changes on the evolution of the protein corona in biological environments. The driving force of the destabilization process observed here is the interaction of partially unfolded proteins with SNPs. Previously we observed that the exacerbated affinity of unfolded proteins for SNPs leads to their adsorption,

even when SNPs are already covered by other proteins.⁶¹ The release of proteins from the hard corona shows that unfolded proteins displaced proteins that initially formed a stable corona. Furthermore, unfolded proteins accessed surface spaces that were completely inaccessible to folded ones. Therefore, unfolded proteins can directly interact with NP surface and reshuffle the hard corona, highlighting the role of protein-surface interactions in the dynamic steps of the protein corona formation in complex environments.

Protein-protein interactions can also take place, such as recognition of adsorbed unfolded proteins by chaperonins,⁶¹ showing that both protein-protein and protein-surface interactions will contribute to shaping a constantly evolving corona.

Using Surface-Plasmon Resonance, Kari *et al.* analyzed the soft and hard protein corona formed on liposomes anchored on a gold surface by flushing human plasma first, followed by buffer.⁷³ Both a hard and a soft corona formed on all liposomes. The analysis of their composition for different liposome formulation revealed that numerous proteins were identified in the hard corona on some liposomes, whilst being found in the soft corona on others. This observation suggests that differences in protein-surface interactions, corresponding to differences in protein affinity, may explain their presence either in the hard or soft corona. These authors conclude that transient protein-surface interactions contribute to the formation of the soft corona in human plasma.

Simon *et al.* showed that the evolution the protein corona in human serum, following changes in protein affinity as a function of temperature, critically affects NP cellular up-

take.⁶⁶ Therefore, a change in protein affinity for NP induced by structural modifications, which we recorded here in a model protein by SRCD and MD simulation, may contribute to the constant reshuffling of the protein corona in a biological environment, where weak protein-surface interactions play a role in NP-cell interactions.

CONCLUSION

Using cryoTEM and SRCD, we show that NPs alter the structure and the stability of weakly bound proteins *in situ*. MD simulation identified the molecular bases driving the formation of the soft corona. Based on thermodynamic analysis, we show that NPs stabilize partially unfolded protein conformations in the soft corona by enthalpy-driven molecular interactions. We suggest that NPs alter weakly bound proteins by shifting the equilibrium towards the unfolded states at physiological temperature. We show that the classical approach based on NP separation from the biological medium fails to detect this effect and therefore cannot be used to fully predict the biological effects of nanomaterials *in situ*.

Further work is required to extend our knowledge on nanoparticle-mediated destabilization effects. In this respect, *in situ* analytical techniques such as SRCD combined with MD simulations represent a powerful analytical approach to investigate destabilization effects of nanoparticles on proteins in biological environments. Advanced methodologies by NMR^{31,32} and SANS^{30,75} have recently been developed to probe the formation of the protein corona *in situ*. They have given us access to a new observation window of NPs in biological environments. Similar techniques could be further extended to investigate the destabilization of proteins by NPs *in situ*.

MATERIALS AND METHODS

Purification of porcine oxyhemoglobin. Hemoglobin in its oxygenated form (oxyHb) was purified from fresh porcine blood (*Sus scrofa domestica*) as described previously.³⁵ Briefly, oxyHb was extracted by erythrocyte membrane precipitation in 280 mM phosphate buffer at pH 7, dialyzed against pure water at 4°C using a membrane cut-off of 14 kDa (Spectra/Pore), stripped on a mixed-bed ion-exchange resin (AG 501-X8, Bio-Rad), and centrifuged at 20,000 g at 4°C for 10 min. The concentration of oxyHb expressed as heme molar concentration ($M = 16,125$ Da) was measured on a Shimadzu UV-2450 spectrophotometer using $\epsilon_{576\text{nm}} = 15,150 \text{ M}^{-1}\cdot\text{cm}^{-1}$. Protein concentration and potential iron oxidation were checked before each experiment by recording the UV-vis spectrum in the 350 nm – 700 nm range.

Silica nanoparticles. SNPs (Ludox TM-50, Sigma) were dialyzed against pure water at 4°C using a membrane cut-off of 3.5 kDa (Spectra/Pore). Then the solution was diluted in pure water, sonicated and filtered through 0.45 μm filters (Minitart). The SNP concentration was measured by drying 1 mL of the solution at 90°C overnight in open Eppendorf tubes and by weighing the final dry mass of 3 samples.

Synchrotron-radiation circular dichroism. Three buffers were used for the SRCD experiments: 100 mM phosphate buffer pH 7 (Sigma 71649, Fisher Scientific S3720), 5 mM ammonium acetate (NH_4Ac) buffer pH 9 (Sigma A7330) and 3 mM N-Cyclohexyl-2-aminoethanesulfonic acid (CHES) buffer

pH 9 (Sigma C8210). The SNPs tend to buffer the solution towards basic pH, especially at high concentration. A minimum concentration of 50 mM phosphate buffer was required to maintain the solution at pH 7 under our experimental conditions. Because of the buffering effect of SNPs, a smaller concentration of CHES or NH_4Ac buffer was required at pH 9. All buffers are stable in the given temperature range. The SRCD experiments were conducted on the DISCO beamline at SOLEIL Synchrotron⁷⁶ (Saint-Aubin, France). 1 mM oxyHb was mixed with 100 mg/mL SNPs in 0.5 mL Eppendorf tubes and gently mixed on a wheel for 1 hour at 22°C. The final pH of the solution was checked using a pH microelectrode (Biotrode, Metrohm). All the samples were vortexed at low speed for a few seconds before analysis. 4 μL of the solution were deposited in a CaF_2 round cell with a 12 μm pathlength⁷⁷ (Hellma). The pathlength was measured by light interference of an empty cell on a spectrophotometer in the visible domain. The spectra were recorded from 170 nm to 260 nm. Temperature experiments were conducted from 22°C to 97°C with a 3°C step and an equilibration time of 5 min. 3 scans were recorded at each temperature step. The experimental spectra were averaged, baseline subtracted and smoothed using CDtoolX software⁷⁸. The intensity of the SRCD signal was calibrated at 192 nm and at 290 nm with a standard solution of camphor-sulfonic acid (CSA) analyzed in a 100 μm pathlength cell at 22°C at a concentration of 6.19 g.L⁻¹. The experimental spectra were converted from θ (millidegrees) to $\Delta\epsilon$ (L.mol⁻¹.cm⁻¹) using a Mean Residue Weight (MRW) of 113.15 for porcine oxyHb. The melting curves represented as $\Delta\epsilon(194 \text{ nm})$ and $\Delta\epsilon(222 \text{ nm})$ as a function of temperature were fitted by one or two successive sigmoids using Igor software. The same results were obtained by applying a two-sigmoidal curve equation to the melting curve following work by Rodnin *et al.*⁷⁹

$$\Delta\epsilon = \frac{P_1}{1 + e^{-\frac{\Delta H_1}{RT}(1 - \frac{T}{T_{m1}})}} + \frac{P_2}{1 + e^{-\frac{\Delta H_2}{RT}(1 - \frac{T}{T_{m2}})}}$$

where P_1 , P_2 are the percentages associated with each transition.

Cryogenic transmission electron microscopy. 200 μM oxyHb was mixed with 9 g/L SNPs in 100 mM phosphate buffer pH 7 or in 5 mM NH_4Ac buffer pH 9 for 1h at room temperature. In these conditions, 56 ± 4 % oxyHb is adsorbed on SNPs at pH 7, whereas no adsorbed protein was detected at pH 9 (<3%). A 4 μL drop was deposited on a glow-discharged holey carbon grid (Quantifoil R2/2). The grid was blotted with a filter paper for 2s and plunged into liquid ethane cooled down by liquid nitrogen using a Vitrobot Mark IV (Thermo Fisher Scientific). The nanoparticles were observed in areas without carbon film to reduce any potential effect of the grids on the samples. Frozen samples were transferred to a Gatan 626 cryo-holder and observed on a JEOL 2010F cryo transmission electron microscope at 200 kV. A minimum of 2 grids per sample were observed and analyzed. Samples were imaged at x60000 magnification using minimal dose system. Images were collected on a Gatan Ultrascan 4K CCD camera with a 2.2 μm nominal defocus. Automated image analysis was developed using ImageJ software for treatment and analysis. CryoTEM images were treated to reduce noise using the Non-local Means Denoising plugin and the Minimum Filter. Single nanoparticles and proteins in the corona were detected with the Particle Analysis plugin and the Watershed function. The same parameters were automatically applied to all the images

(Table S1). A minimum of 100 nanoparticles were analyzed in each condition. A total number of 310 and 244 SNPs were analyzed with and without oxyHb respectively. The average SNP and protein diameters were calculated for a sphere from the area measured by automatic image analysis and expressed as mean \pm std. dev.

Secondary structure determination. The secondary structure of oxyHb was determined by deconvolution of the SRCD spectra with and without SNPs as a function of temperature with BeStSel algorithm.^{42,43} The structural analysis was validated by comparing the results with the crystal structure of porcine oxyHb. The secondary structure of oxyHb was calculated with DSSP tool from the crystal structure of porcine oxyHb 1QPW.pdb.^{44,45}

Molecular dynamics simulations. Initial coordinates of oxyHb were obtained from the crystal structure available on the Protein Data Bank (1QPW.pdb). Crystallographic solvent atoms were removed from the structure. All atom simulations were performed with the AMBER 16 package.⁸⁰ The ff99SB-ildn force field was employed for the protein and those of the heme were obtained from Giammona *et al.*⁸¹ The starting structures were immersed in a cubic box with boundaries extending from the protein periphery for at least 10 Å in all directions. The box was filled with TIP3P water molecules, and an appropriate number of counter ions was added to neutralize the total charge of the system. The systems were minimized using steepest descent algorithm and used to initiate molecular dynamics. For each system, 4 molecular dynamic simulations were carried out at different temperatures: 295 K, 322 K, 353K, and 400 K. The SHAKE algorithm was used to constrain the motion of hydrogen-containing bonds within a 2 fs time step to integrate the equations of motions. The cut-off distance for van der Waals interactions was set to 10 Å and the Ewald particle mesh was used for long-range electrostatic interactions. Production runs of 150 ns were performed on the equilibrated structure with the NPT ensemble using a 2 fs time step. The trajectories were clusterized with TTclust program⁸² and analysed with CPPTRAJ module, VMD 1.9.3⁸³ and Chimera 1.10.2⁸⁴ programs. The solvent-accessible surface area was calculated using the "surf" command in the CPPTRAJ module.

Density Functional Theory calculation. The interaction of the peptides with the silica surface was analyzed at the quantum chemistry level using Dispersion-Corrected Self Consistent Charge Density Functional Tight Binding (DC-SCC-DFTB) with DFTB+ software⁵⁸ using Lennard Jones forces from the universal force field. Parameters from the "matsci-0-3" set were used for the dispersion correction from the universal force field.⁸⁵ The terminal -CO and -NH groups of the sequence were replaced by H atoms. The positions of the N and C atoms of the peptide bond were frozen for the initial minimization of the average peptide structure extracted from the MD simulation of the oxyHb structure. The silica surface was obtained with CHARMM GUI.⁸⁶ During the interaction of the peptide with the silica surface, the N and C atoms of the two first residues only, were frozen. The peptides were positioned 2 Å away from the silica surface at the beginning of the calculation, with the most mobile region corresponding to residues 50-56 oriented towards the surface. The energy of the system

(445 atoms) was minimized by DC-SCC-DFTB using a conjugated gradient approach. The same approach was used for the calculation of the interaction with the silica slab and the silica cluster. The uncertainty on the calculated interaction energy is 0.2 eV.

Dynamic light scattering. The size of SNPs was measured in each buffer at 22°C by Dynamic Light Scattering (DLS). The measurement was performed on a Zetasizer Nano-ZS (Malvern Instruments) with a particle concentration of 0.01 mg/mL. Each measurement was repeated at least 3 times. The correlation coefficient measured by DLS was used to monitor the dissolution of the SNPs at pH 9 over time. The SNP suspensions were analyzed after 3 days of incubation at 22°C in NH₄Ac buffer and in CHES buffer at pH 9. The zeta potential (ζ) of SNPs was measured at 22°C on the same instrument with a particle concentration of 1 mg/mL. Each measurement was repeated at least 3 times. The ζ -potential was calculated by fitting the electrophoretic mobility with the Smoluchowski model.

Adsorption isotherms. The same conditions of incubation of oxyHb with SNPs were used for the SRCD experiments and for the measurement of the adsorption isotherms. After incubation, the samples were centrifuged at 30,000 g for 20 minutes. The concentration of oxyHb in the supernatant was measured by spectrophotometry on a Shimadzu UV-2450 spectrophotometer. The adsorption of oxyHb on SNPs was measured at 22°C, 30°C, 40°C, and 50°C in each buffer condition. The experiments were performed in triplicate. The amount of adsorbed oxyHb is expressed as average \pm standard deviation.

ASSOCIATED CONTENT

Supporting Information. Characterization of SNPs; analysis of cryoTEM images; secondary structure determination; thermal unfolding of oxyHb; protonation state of histidine in oxyHb structure; MD simulation of oxyHb structure; calculation of the surface accessibility of oxyHb; DC-SCC-DFTB calculation of peptide-silica interactions; thermodynamic analysis of oxyHb unfolding. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Author Contributions

SD, SP and JPR designed the study. SD, DSG and FW performed the SRCD experiments. LM and JD carried out cryo-TEM imaging. GCC contributed to sample preparation and measured the adsorption isotherms. YB performed the MD simulations. JPR performed DFTB calculations. SD developed thermodynamic data analysis with input from JPR. SD wrote the manuscript. All authors discussed the results and contributed to the final manuscript.

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