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To cite this version:
Laurent Colas, P Jacquet, V. Bobkov, M. Brix, L. Meneses, et al.. 2D mappings of ICRF-induced SOL density modifications on JET. 45th EPS Conference on Plasma Physics, Jul 2018, Prague, Czech Republic. cea-02103387

HAL Id: cea-02103387
https://hal-cea.archives-ouvertes.fr/cea-02103387
Submitted on 18 Apr 2019

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2D mappings of ICRF-induced SOL density modifications on JET

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Waves in the Ion Cyclotron Range of Frequencies (ICRF, 20-80MHz) provide core ion heating in tokamaks and are used to limit impurity accumulation in high-performance JET scenarios \cite{1,2}. Before reaching the core, these radiofrequency (RF) waves, excited at the Low-Field Side of the torus, interact with the Scrape-Off-Layer (SOL), causing enhanced wall sputtering, heat loads and local density ($n_e$) changes with a complex 3D geometry. Extending previous studies on JET \cite{3,4,5}, this paper aims at mapping RF-induced SOL $n_e$ patterns in 2D. This puts constraints on SOL RF modelling and provides hints for locating RF-specific W-sources on the JET ITER-Like-Wall (ILW), similar to the ITER vessel.

\section{I. Experimental setup and 2D mapping technique}

At JET, ICRF waves were excited at a frequency of 42MHz by phased toroidal arrays of four poloidal RF current straps (A2 antennas \cite{6}) as well as a 2\times4 ITER-like array (ILA \cite{7}). Four A2 antennas (named A-D) are located toroidally around the torus, while the ILA is placed between A and B (see figure 1). Upper and Lower halves of the ILA can be operated independently of D and of the pair A+B, whose feeding transmission lines are coupled. 500kW power was delivered per antenna, with a toroidal phasing $\Delta \varphi = \pi/2$rad. SOL density ($n_e$) distributions are measured by Lithium Beam Emission Spectroscopy (Li-BES) \cite{8} and X-mode reflectometry \cite{9}, whose Lines Of Sight (LOS) have different magnetic connections to the various antennas (see Figure 1). D[H]-heated L-mode pulses at 2.4T were studied, where

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ICRF antennas were toggled while the plasma current was ramped up (see figure 2). D$_2$ gas fueling from a toroidally-distributed valve was feedback-controlled to maintain a constant core density. In ohmic regime Li-BES and reflectometry measured similar profiles within a 4cm radial shift. During ICRH $n_e$ modifications with respect to ohmic references at similar edge safety factor $q_{95}$, were recorded in the outer SOL, several meters away toroidally from the active antennas. Earlier studies showed that these reproducible changes increased in magnitude with larger RF power and when switching from phasing $\Delta \varphi = \pi$ to $\Delta \varphi = \pi/2$ [4]. They depended on $q_{95}$, on which antenna was active and which diagnostic monitored $n_e$. This is attributed to varying LOS-antenna magnetic connections. Following this idea, supported by other observations (e.g. [5]), and assuming a parallel homogeneity of the RF-induced SOL patterns, the ratio of time-averaged $n_e$ with or without RF waves was plotted versus the
location of observation points along the diagnostic LOS, mapped in front of each RF antenna using field line tracing. In this technique, first proposed in [10], the LOS ensured the radial resolution of 2D maps, while scanning $q_{95}$ provided their vertical resolution.

II. Experimental 2D maps of RF-induced density variations.

Figure 3 maps density ratios during A2 antennas with respect to ohmic reference. Ratios in the range of 0.4 (depletion) to 2.5 (over-density) were observed, mostly in the outer SOL, extending radially 2-3cm in front of antenna limiters, with poloidal asymmetries.

Figure 3: 2D maps of $n_e$ ratios during antenna D and A+B to ohmic, mapped to antennas D and A.

Figure 4 compares $n_e$ modifications when powering lower and upper ILA, along the Li-beam LOS connected magnetically to the upper array. $n_e$ changes in this region even when the un-connected lower ILA is active. When switching to the upper array, the density depletion stays similar in the upper part of the map while an over-density area develops at the bottom.

→ Figure 4: $\log_{10}(n_e(\text{ILA})/n_e(\Omega))$ along Li-beam LOS mapped to ILA, with lower ILA (left panel, JPN 90456) and upper ILA (right panel, JPN 90509) active. Color scale: [-0.4 ;0.4].
III. Discussion and outlook.

Many machines have evidenced ICRF-induced SOL $n_e$ modifications, with a phenomenology similar to JET [10-14]. The emitted oscillating near fields likely cause SOL biasing to large direct current (DC) potentials $V_{DC}$ by RF-sheath rectification, and subsequent $E\times B$ $n_e$ convection in $E_{DC}=-\nabla V_{DC}$. Large $V_{DC}$ exceeding sputtering thresholds [14-16] and SOL cross-field flows were measured during ICRH [14-17]. In addition to other experiments, e.g. [5], reproducing the observed $n_e$ patterns, in relation with antenna electrical settings, can be used to validate RF-sheath modelling tools. Present simulations suggest that powering the lower ILA with $\Delta \varphi=\pi/2$ enhances $V_{DC}$ also in the upper part of the antenna, but less than using the upper array, qualitatively consistent with figure 4 [18]. Within RF-induced convection the maximal $V_{DC}$, most efficient for wall sputtering, is expected at the center of the convective cells. Figure 5 shows that some flux tubes with modified $n_e$ connect to W-components on the ILW, e.g. the tiles at low divertor entrance. Limited W spectroscopy diagnostics at these candidate sputtering locations and the toroidal spread of connection points, might explain why RF-specific W-sources have not yet been evidenced directly on the ILW [3]. This might guide further attempts at detection. It is also proposed to reduce RF-SOL interactions by tuning the power balance between inner and outer straps of A2 antennas [19].

Acknowledgements. This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the European research and training programme under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.

IV. References

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