

Renormalization of $SU(2)$ Yang–Mills theory with flow equations

Alexander Efremov, Riccardo Guida, Christoph Kopper

► **To cite this version:**

Alexander Efremov, Riccardo Guida, Christoph Kopper. Renormalization of $SU(2)$ Yang–Mills theory with flow equations. *Journal of Mathematical Physics*, American Institute of Physics (AIP), 2017, 58, pp.093503. 10.1063/1.5000041 . cea-01614681

HAL Id: cea-01614681

<https://hal-cea.archives-ouvertes.fr/cea-01614681>

Submitted on 8 Apr 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Renormalization of $SU(2)$ Yang–Mills theory with Flow Equations

Alexander N. Efremov
Université Paris-Saclay, Doctoral School ED564,
CPHT, Ecole Polytechnique, F-91128 Palaiseau, France.
alexander@efremov.fr

Riccardo Guida
Institut de Physique Théorique, Université Paris-Saclay, CEA, CNRS,
F-91191 Gif-sur-Yvette, France
ORCID iD: orcid.org/0000-0002-2854-7214
ResearcherID: N-7759-2013
riccardo.guida@cea.fr

Christoph Kopper
Centre de Physique Théorique de l'Ecole Polytechnique,
CNRS, UMR 7644, F-91128 Palaiseau, France.
Christoph.Kopper@polytechnique.edu

Abstract

We give a proof of perturbative renormalizability of $SU(2)$ Yang–Mills theory in four-dimensional Euclidean space which is based on the Flow Equations of the renormalization group. The main motivation is to present a proof which does not make appear mathematically undefined objects (as for example dimensionally regularized generating functionals), which permits to parametrize the theory in terms of *physical* renormalization conditions, and which allows to control the singularities of the correlation functions of the theory in the infrared domain. Thus a large part of the proof is dedicated to bounds on massless correlation functions.

Preprint IPhT t17/055
RR 018-04-2017

Contents

1	Introduction	2
2	The formalism	5
2.1	Notations	5
2.2	The Classical Action	8
2.3	Generating functionals and Flow equations	9
2.4	Violated Slavnov–Taylor identities	18
2.4.1	The Antighost Equation	23
3	Momentum bounds	23
3.1	Weighted trees	25
3.2	Theorems	28
4	Proof of Theorems 1–4	34
4.1	Chains of vertex functions	35
4.2	Junction of weighted trees	38
4.3	Irrelevant terms	43
4.4	Marginal terms	46
4.4.1	$\Gamma^{c\bar{c}AA}$ and $\Gamma^{c\bar{c}\bar{c}\bar{c}}$	46
4.4.2	$\partial_A \Gamma^{c\bar{c}A}$	47
4.4.3	Renormalization at $\Lambda = 0$	47
4.5	Strictly relevant terms	50
4.5.1	$\partial \Gamma_1^{\phi^3}$, $\Gamma_1^{\phi^4}$ and $\Gamma_{1\gamma}^{cc}$	51
4.5.2	$\Gamma_{\beta}^{\bar{\phi};w}$, $\Gamma_1^{\phi^2;w}$ and $\Gamma_1^{\phi^3}$	52
4.6	Convergence	52
5	Restoration of the STI	55
5.0.1	Smallness relations	55
5.0.2	The functional $F_{1,rel}$	56
5.0.3	Consistency conditions	57
5.0.4	Existence of a constant solution	60
5.1	Γ_{β}^c	60
5.2	Γ_{β}^{cA}	61
5.3	Γ_{β}^{cAA} and $\Gamma_{\beta}^{c\bar{c}\bar{c}}$	62
5.4	$\Gamma_1^{c\bar{c}\bar{c}\bar{c}}$ and Γ_1^{cAAAA}	62
5.5	$\Gamma_1^{c\bar{c}cAA}$	62

5.6	Γ_1^{cA}	62
5.7	Γ_1^{cAA}	63
5.8	$\Gamma_1^{\bar{c}cA}$ and $\Gamma_{1\gamma}^{cA}$	64
5.9	$\Gamma_1^{\bar{c}c}$ and $\Gamma_{1\gamma}^{cc}$	64
5.10	Γ_1^{cAAA}	64
6	Acknowledgements	65
	Appendices	65
A	Properties of Gaussian measures	65
B	Chains of vertex functions	66
C	Tensors	67
D	Basic estimates	74
E	The functional $\Gamma_{n \leq 4}$	88
F	The functionals $\Gamma_{\gamma; n \leq 2}$ and $\Gamma_{\omega; n \leq 2}$	90
G	Violated STI for $\Lambda > 0$	91
H	List of the renormalization points	92
I	List of insertions	93

1 Introduction

Renormalization theory based on the differential Flow Equations [Wil71, WH73, Pol84] has allowed to build a unified approach to the analysis of renormalizability for a wide class of theories without recourse to Feynman graphs. This method was applied to show momentum bounds of massive ϕ^4 theory [KM02], used to prove renormalization of spontaneously broken SU(2) Yang–Mills [KM09] and to establish uniform bounds on Schwinger functions of massless ϕ^4 fields [GK11]. Starting with the milestone works, [YM54, FP67, Sla72, Tay71, LZJ72, tHV72, BRS75, Tyu75, ZJ75] (see [Lai81]

for more references), a variety of results on the renormalizability of non-abelian gauge theories flourished in the literature, in different contexts and with different level of mathematical rigor. Work on this problem in the context of Flow Equations (FE) includes [RW94, Bec96, BDM95, MD96], and, more recently, [FHH16]. The present work deals with perturbative renormalizability and shares certain aspects with some of these articles: from [Bec96] we borrow the fruitful idea that the local operator describing the violation of Slavnov–Taylor identities (STI) for the one-particle irreducible (1PI) functions [ZJ75, KSZ75] is constrained by the nilpotency of the underlying ”Slavnov differential operator”; as in [BDM95] (and in contrast with [FHH16]) we define the marginal correlators by physical boundary conditions at vanishing IR cutoff and nonexceptional momenta. Our main result is a proof of the renormalizability of Yang–Mills theory that complements the previous work in view of the following features: i) Rigorous control of the IR and UV behavior of the one-particle irreducible functions is established by means of uniform bounds in momentum space. In particular, from the bounds follow the existence of the IR limit of 1PI functions at nonexceptional momenta and the existence of the subsequent UV limit. ii) The vanishing of the STI violation in the UV limit is proven; iii) The relevant part of 1PI functions and STI violation at nonvanishing renormalization momenta, which gives initial conditions for the renormalization group FE, is thoroughly identified.

Our proof refers explicitly to SU(2) Yang–Mills theory. However, it could be extended without important conceptual changes to other semi-simple compact Lie groups.

We proceed as follows. In section 2 we fix the notations, introduce the classical Yang–Mills action, generating functionals and regulators. We define the BRST transformations. Then we derive the FE of the renormalization group, for the connected amputated Schwinger functions and for the one-particle irreducible vertex functions. Finally we study the STI, the Antighost Equation (AGE) and their violation. Remark that in our context gauge invariance is broken through the presence of the cutoffs.

Our proof of renormalizability of Yang–Mills theory is based on momentum bounds for the vertex bounds which permit to take the limits $\Lambda \rightarrow 0$ (IR-cutoff) and $\Lambda_0 \rightarrow \infty$ (UV-cutoff) for nonexceptional external momenta. These bounds are an extension of those in [GK11] and are established inductively with the aid of the FE. They are expressed in terms of tree amplitudes and polynomials of logarithms. For our trees we only have to con-

sider vertices of coordination numbers 1 and 3. In section 3 we present the definitions of the tree structures to be considered, and the aforementioned bounds. These bounds permit to prove the existence of the vertex functions for nonexceptional external momenta on removal of the cutoffs. This statement also holds for vertex functions carrying an operator insertion related to the BRST transformation of the fields. These vertex functions are required to formulate the STI. Then we also have to consider vertex functions with operator insertions which permit to formulate the violation of the STI. Our bounds then permit to show that, for suitable renormalization conditions, these functions describing the violation of the STI vanish in the UV limit.

Section 4 is dedicated to the proof of those bounds. The rhs of the FE is a sum of products of vertex functions in lower loop order joined by free propagators. Our technique of proof then is based on the fact that applying our inductive bounds on these chains, which still have to be closed by a Λ -derived propagator and then integrated over the circulating loop momentum and over Λ , reproduce our inductive bounds at loop order l . The proof treats irrelevant terms first, then marginal and finally strictly relevant ones. Particular attention has to be paid to the renormalization conditions. Section 4 ends with a proof of UV-convergence for $\Lambda_0 \rightarrow \infty$.

In section 5 we prove that the renormalization conditions required to prove the bounds of the previous sections can actually be imposed. These renormalization conditions are such that they leave us free to fix the physical coupling of the theory. At the same time they permit us to make vanish the relevant part of all functions describing the violation of the STI. This is required in the previous proof.

In the Appendices A, B, C, D we present some facts on Gaussian measures, an example of chains of vertex functions, an analysis of linear independence of euclidean tensor structures, a large number of elementary bounds on integrals we encounter in the proofs. We also add a list with bounds on the propagators and their derivatives. In the two subsequent appendices E, F we analyse the generating functionals of the (inserted) vertex functions, as far as they have relevant content. Appendix G is a supplement to section 2.4 and contains the generalization of the STI and AGE to the case of non-vanishing IR cutoff. In the last two appendices H, I we present the list of renormalization points and operator insertions to be considered.

2 The formalism

2.1 Notations

\mathbb{N} is the set of nonnegative integers. $|S|$ is the cardinality of a set S . Furthermore, (a, b, c, \dots) , $\{a, b, c, \dots\}$ denote a sequence and a set, respectively. Unless otherwise stated, sequence stands for finite sequence. For shortness we set $[a : b] := \{i \in \mathbb{Z} : a \leq i \leq b\}$ and $[b] := [1 : b]$. Repeated indices are implicitly summed over, e.g. $A^a t_a := \sum_a A^a t_a$. We choose the following basis of the Lie algebra

$$t_a := \frac{1}{2} \sigma_a, \quad [t_a, t_b] = i \epsilon_{abc} t_c, \quad a, b, c \in \{1, 2, 3\}, \quad (1)$$

where σ_a are the Pauli matrices and ϵ_{abc} is the Levi-Civita symbol, $\epsilon_{123} = 1$. In this article we will deal with tensor fields on \mathbb{R}^4 in Cartesian coordinate systems with metric tensor $\delta_{\mu\nu}$. If A, B are two Cartesian tensors of \mathbb{R}^4 of rank r with components $A_{\mu_1 \dots \mu_r}$ and $B_{\mu_1 \dots \mu_r}$, respectively, then the scalar product (A, B) is the contraction

$$(A, B) := A_{\mu_1 \dots \mu_r}^* B_{\mu_1 \dots \mu_r}. \quad (2)$$

Given a Cartesian tensor T , we use the norm

$$|T| := (T, T)^{\frac{1}{2}}. \quad (3)$$

For instance, for $p \in \mathbb{R}^4$, $|p|^2 = \sum_\mu p_\mu^2$. Let T, A , and B be Cartesian tensors such that $T_{\vec{\mu}\vec{\nu}} = A_{\vec{\mu}\vec{\rho}} B_{\vec{\rho}\vec{\nu}}$ where $\vec{\mu}, \vec{\rho}, \vec{\nu}$ are multi-indices, for example $\vec{\mu} := (\mu_1, \dots, \mu_n)$. Then using the Cauchy-Schwarz inequality

$$|T|^2 = \sum_{\vec{\mu}, \vec{\nu}} \left| \sum_{\vec{\rho}} A_{\vec{\mu}\vec{\rho}} B_{\vec{\rho}\vec{\nu}} \right|^2 \leq \sum_{\vec{\mu}, \vec{\nu}} \sum_{\vec{\rho}, \vec{\sigma}} |A_{\vec{\mu}\vec{\rho}}|^2 |B_{\vec{\sigma}\vec{\nu}}|^2 = |A|^2 |B|^2. \quad (4)$$

The integral over \mathbb{R}^4 of the product of two functions is denoted by

$$\langle f_1, f_2 \rangle := \int d^4x f_1(x) f_2(x), \quad (5)$$

and the Fourier transform of a function is defined by

$$f(p) := \int d^4x e^{-ipx} f(x). \quad (6)$$

The convolution of two functions f, g is denoted as below

$$(f * g)(x) = \int d^4 y f(y)g(x - y). \quad (7)$$

For functions $\phi_i(p_i)$ and $F(p_1, \dots, p_k)$ with $p_i \in \mathbb{R}^4$ the symbol $\langle F | \phi_1 \dots \phi_k; p \rangle$ denotes the following integral in momentum space

$$\langle F | \phi_1 \dots \phi_k; p \rangle := \int (2\pi)^4 \delta\left(\sum_{j=1}^k p_j + p\right) F(p_1, \dots, p_k) \phi_1(p_1) \dots \phi_k(p_k) \prod_{i=1}^k \frac{d^4 p_i}{(2\pi)^4}. \quad (8)$$

We also use the shorthands

$$\langle \phi_1 \dots \phi_k; p \rangle := \langle 1 | \phi_1 \dots \phi_k; p \rangle, \quad \langle \phi_1 \dots \phi_k \rangle := \langle 1 | \phi_1 \dots \phi_k; 0 \rangle. \quad (9)$$

A cumulative notation for the elementary fields and corresponding sources is

$$\Phi := (A_\mu^a, B^a, c^a, \bar{c}^a), \quad K := (j_\mu^a, b^a, \bar{\eta}^a, \eta^a), \quad (10)$$

where $c, \bar{c}, \eta, \bar{\eta}$ are generators of an infinite-dimensional anticommuting algebra. Furthermore, we use the following shorthand

$$K \cdot \Phi := \langle j_\mu^a, A_\mu^a \rangle + \langle b^a, B^a \rangle + \langle \bar{\eta}^a, c^a \rangle + \langle \bar{c}^a, \eta^a \rangle. \quad (11)$$

We will have to consider one-particle irreducible functions, also known as vertex functions, whose generating functional is denoted by Γ . These functions are translation-invariant in position space. Their *reduced* Fourier transforms $\Gamma^{\vec{\phi}}$ are defined as follows

$$\Gamma^{\vec{\phi}}(p_1, \dots, p_{N-1}) := \int \left(\prod_{i=1}^{N-1} d^4 x_i e^{-ip_i x_i} \right) \Gamma^{\vec{\phi}}(0, x_1, \dots, x_{N-1}), \quad (12)$$

where $\vec{\phi} := (\phi_0, \dots, \phi_{N-1})$ is a sequence of field labels, $\phi_i \in \{A, B, c, \bar{c}\}$, and

$$\Gamma^{\vec{\phi}}(x_0, x_1, \dots, x_{N-1}) := \left(\frac{\delta}{\delta \phi_0(x_0)} \cdots \frac{\delta}{\delta \phi_{N-1}(x_{N-1})} \right) \Gamma \Big|_{\vec{\phi}=0}. \quad (13)$$

Note that, with the exception of appendix A, derivatives wrt Grassmannian variables are always intended to be left derivatives, see e.g. [Ber66]. The complete Fourier transformed N-point vertex function then satisfies

$$(2\pi)^{4(N-1)} \left(\frac{\delta}{\delta \phi_0(-p_0)} \cdots \frac{\delta}{\delta \phi_{N-1}(-p_{N-1})} \right) \Gamma \Big|_{\vec{\phi}=0} = \delta\left(\sum_{i=0}^{N-1} p_i\right) \Gamma^{\vec{\phi}}, \quad (14)$$

where $\Gamma^{\vec{\phi}}$ stands for $\Gamma^{\vec{\phi}}(p_1, \dots, p_{N-1})$. The reduced and complete Fourier transforms with $n_\chi \geq 1$ composite operator insertions of sources $\vec{\chi} = (\chi_0, \dots, \chi_{n_\chi-1})$ and $\vec{\phi} := (\phi_{n_\chi}, \dots, \phi_{n-1})$ are correspondingly related by

$$(2\pi)^{4(n-1)} \left(\prod_{i=n_\chi}^{n-1} \frac{\delta}{\delta\phi_i(-p_i)} \prod_{i=0}^{n_\chi-1} \frac{\delta}{\delta\chi_i(-p_i)} \right) \Gamma \Big|_{\substack{\vec{\phi}=0 \\ \vec{\chi}=0}} = \delta \left(\sum_{i=0}^{n-1} p_i \right) \Gamma_{\vec{\chi}}^{\vec{\phi}}, \quad (15)$$

where the order of the derivatives $\delta/\delta\phi_i$ is the same as before, the derivatives $\delta/\delta\chi_i$ are ordered with left-to-right increasing indices, and $\Gamma_{\vec{\chi}}^{\vec{\phi}}$ stands for $\Gamma_{\vec{\chi}}^{\vec{\phi}}(p_1, \dots, p_{n-1})$. Note that $n-1$ is the total number of arguments, e.g. for $\Gamma_{\vec{\chi}}^{\vec{\phi}}(p_1, \dots, p_{n-1})$ we have $n = n_\chi + N$.

It will be useful to keep a bijective relation between momenta and field labels (including possible insertion labels), $p_i \leftrightarrow \phi_i$. Hence, we assume that p_0 is the negative subsum of all other momenta,

$$\mathbb{P}_n := \{ \vec{p} \in \mathbb{R}^{4n} : \vec{p} = (p_0, \dots, p_{n-1}), p_0 = - \sum_{i=1}^{n-1} p_i \}, \quad |\vec{p}|^2 := \sum_{i=0}^{n-1} p_i^2, \quad (16)$$

and, referring to (14), (15), we use the notation

$$\Gamma^{\vec{\phi}}(\vec{p}) = \Gamma^{\vec{\phi}}(p_1, \dots, p_{N-1}), \quad \vec{p} = (p_0, \dots, p_{N-1}) \in \mathbb{P}_N, \quad (17)$$

$$\Gamma_{\vec{\chi}}^{\vec{\phi}}(\vec{p}) = \Gamma_{\vec{\chi}}^{\vec{\phi}}(p_1, \dots, p_{n-1}), \quad \vec{p} = (p_0, \dots, p_{n-1}) \in \mathbb{P}_n. \quad (18)$$

A momentum configuration $\vec{p} \in \mathbb{P}_n$ is said **exceptional** iff there exists a nonempty $S \subset [1 : n-1]$ such that $\sum_{i \in S} p_i = 0$. If this is not the case \vec{p} is said **nonexceptional**.

We rely on the multi-index formalism for derivatives with respect to momenta. Taking in account that there are no derivatives wrt p_0 , we set:

$$\mathbf{w} := (w_{0,1}, w_{0,2}, \dots, w_{n-1,4}), \quad w_{i,\mu} \in \mathbb{N}, \quad w_{0,\mu} := 0, \quad (19)$$

$$w := (w_0, \dots, w_{n-1}), \quad w_i := \sum_{\mu=1}^4 w_{i,\mu}, \quad (20)$$

$$\mathbb{W}_n := \{ w \in \{0\} \times \mathbb{N}^{n-1} : \|w\| \leq w_{max} \}, \quad \|w\| := \sum_{i=0}^{n-1} w_i, \quad (21)$$

$$\partial^{\mathbf{w}} \Gamma := \prod_{i=0}^{n-1} \prod_{\mu=1}^4 \left(\frac{\partial}{\partial p_{i,\mu}} \right)^{w_{i,\mu}} \Gamma, \quad \partial^w \Gamma := \left(\prod_{i=0}^{n-1} \partial_{p_i}^{w_i} \right) \Gamma, \quad (22)$$

where ∂_p^k is the tensor with components $\partial_{p_{\mu_1}} \dots \partial_{p_{\mu_k}}$, and w_{max} is an arbitrary integer ≥ 4 fixed throughout the paper. The following shorthands will be helpful:

$$\dot{\Gamma}^{\Lambda\Lambda_0} := \partial_\Lambda \Gamma^{\Lambda\Lambda_0}, \quad \partial_{\phi_i} \Gamma^{\vec{\phi}} := \frac{\partial \Gamma^{\vec{\phi}}}{\partial p_i}, \quad \Gamma^{;w} := \partial^w \Gamma, \quad (23)$$

$$\tilde{\delta}_{\phi(p)} := (2\pi)^4 \frac{\delta}{\delta \phi(-p)}, \quad \check{\phi}(p) := \phi(-p), \quad \log_+ x := \log \max(1, x). \quad (24)$$

2.2 The Classical Action

We start with a four-dimensional Euclidean classical field theory invariant under local $SU(2)$ transformations and with a coupling $g \neq 0$. Given $G = e^{ig\alpha}$ with $\alpha = \alpha^a t_a$ an element of the Lie algebra, a gauge transformation reads:

$$A_\mu \mapsto \frac{i}{g} G \partial_\mu G^{-1} + G A_\mu G^{-1}, \quad \phi \mapsto G \phi G^{-1}, \quad (25)$$

and in infinitesimal form

$$A_\mu \mapsto A_\mu + \partial_\mu \alpha - ig[A_\mu, \alpha], \quad \phi \mapsto \phi + ig[\alpha, \phi], \quad (26)$$

where we also show the transformation for a field ϕ in the adjoint representation of the Lie algebra. The quantization of gauge theories using auxiliary ghost fields is presented, for example, by Slavnov and Faddeev in their book [SF91]. Following [PF67, tH71], with an appropriate linear gauge-fixing functional the semiclassical Lagrangian density takes the form

$$\mathcal{L}_0^{tot} = \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \frac{1}{2\xi} (\partial_\mu A_\mu^a)^2 - \partial_\mu \bar{c}^a (D_\mu c)^a, \quad (27)$$

$$D_\mu c = \partial_\mu c - ig[A_\mu, c], \quad (28)$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu - ig[A_\mu, A_\nu], \quad (29)$$

where $\xi > 0$ is the Feynman parameter and $\{F_{\mu\nu}, A_\mu, c\}$ are elements of the $su(2)$ algebra, e.g. $A_\mu = t^a A_\mu^a$. We will study the quantum theory in a framework which breaks local gauge invariance due to the presence of momentum space regulators. However, the Lagrangian respects global $SU(2)$ symmetry, Euclidean isometries $ISO(4)$ and has ghost number zero. We

admit all counterterms compatible with these symmetries:

$$\begin{aligned}
\mathcal{L}_{ct} = & r^0 \bar{c} \bar{c} \bar{c} \bar{c} c^b c^b \bar{c}^a c^a + r_1^0 \bar{c} c A A c^b c^b A_\mu^a A_\mu^a + r_2^0 \bar{c} c A A \bar{c}^a c^b A_\mu^a A_\mu^b \\
& + r_1^0 A A A A A_\mu^b A_\nu^b A_\mu^a A_\nu^a + r_2^0 A A A A A_\nu^b A_\nu^b A_\mu^a A_\mu^a + 2 \epsilon_{abc} r^0 A A A (\partial_\mu A_\nu^a) A_\mu^b A_\nu^c \\
& - r_1^0 \epsilon_{abd} \bar{c}^a (\partial_\mu \bar{c}^a) A_\mu^b c^d - r_2^0 \epsilon_{abd} \bar{c}^a A_\mu^b \partial_\mu c^d + \Sigma^0 \bar{c} c \bar{c}^a \partial^2 c^a \\
& - \frac{1}{2} \Sigma_T^0 A_\mu^a (\partial^2 \delta_{\mu\nu} - \partial_\mu \partial_\nu) A_\nu^a + \frac{1}{2\xi} \Sigma_L^0 (\partial_\mu A_\mu^a)^2 \\
& + \frac{1}{2} \delta m_{AA}^2 A_\mu^a A_\mu^a - \delta m_{\bar{c}c}^2 \bar{c}^a c^a.
\end{aligned} \tag{30}$$

There are eleven marginal counterterms and two strictly relevant counterterms. To denote the marginal counterterms we use the symbols r^0 and Σ^0 . For the strictly relevant counterterms we use δm_{AA}^2 and $\delta m_{\bar{c}c}^2$.

2.3 Generating functionals and Flow equations

In the next two sections we will introduce the essential structural tools required for our proof of renormalizability of nonabelian Yang–Mills theory. These tools are on the one hand the differential Flow Equations of the renormalization group, and on the other hand the (violated) Slavnov–Taylor identities. They are both obtained from the functional integral representation of the theory.

Functional integrals are known to exist, beyond perturbation theory, if suitably regularized [GJ87]. To restrict their support to sufficiently regular functions in position space one has to introduce an ultraviolet (UV) cutoff. To avoid infrared problems one has to introduce a corresponding infrared cutoff, the most practical one being a finite volume cutoff, introduced by putting the system on a torus with periodic boundary conditions. For an analysis of the support properties of (Gaussian) functional integrals see [Ree73].

We do not make explicit the finite volume cutoff in this paper, which would still blow up the manuscript to some degree. For some more details in this respect see [Mül03]. On the other hand the UV cutoff is explicit.

Once the finite volume cutoff has been introduced, the FEs (with UV regulator) can be shown to hold up to $\Lambda = 0$, Λ being the flow parameter, even beyond perturbation theory (as far as the bosonic degrees of freedom are concerned). One can then construct perturbative solutions of the FEs, which are well-defined and regular in the open domain of nonexceptional external momenta, and which satisfy the Slavnov–Taylor identities, in the

limits where the flow parameter is sent to zero and the UV cutoff is sent to infinity. If keeping the finite volume until the end, one then has to show that these solutions have uniform limits (in compact subsets of the open domain of nonexceptional momenta) once the volume is sent to infinity. Short-circuiting this reasoning, as we do here, is tantamount to postulating from the beginning that the FEs and the (violated) Slavnov-Taylor identities hold, for nonexceptional momenta, also in the limit $\Lambda \rightarrow 0$.

By definition, the characteristic function of a finite measure $d\nu$ is its Fourier transform, see e.g. [DF91, GJ87],

$$\chi(K) := \int d\nu(\Phi) e^{\frac{i}{\hbar} K \cdot \Phi}. \quad (31)$$

We consider a Gaussian measure $d\nu_{\Lambda\Lambda_0}$ with characteristic function

$$\chi^{\Lambda\Lambda_0}(j, b, \bar{\eta}, \eta) = e^{\frac{1}{\hbar} \langle \bar{\eta}, S^{\Lambda\Lambda_0} \eta \rangle - \frac{1}{2\hbar} \langle j, C^{\Lambda\Lambda_0} j \rangle - \frac{1}{2\hbar\xi} \langle b, b \rangle}, \quad (32)$$

and we denote as usual

$$(C^{\Lambda\Lambda_0} j)_\mu(x) = \int d^4y C_{\mu\nu}^{\Lambda\Lambda_0}(x, y) j_\nu(y), \quad (33)$$

$$C_{\mu\nu}^{\Lambda\Lambda_0}(x, y) = \int \frac{d^4p}{(2\pi)^4} e^{ip(x-y)} C_{\mu\nu}^{\Lambda\Lambda_0}(p), \quad (34)$$

where in momentum space the regularized propagators are defined by the expressions

$$C_{\mu\nu}(p) := \frac{1}{p^2} (\delta_{\mu\nu} + (\xi - 1) \frac{p_\mu p_\nu}{p^2}), \quad C_{\mu\nu}^{\Lambda\Lambda_0}(p) := C_{\mu\nu}(p) \sigma_{\Lambda\Lambda_0}(p^2), \quad (35)$$

$$S(p) := \frac{1}{p^2}, \quad S^{\Lambda\Lambda_0}(p) := S(p) \sigma_{\Lambda\Lambda_0}(p^2), \quad (36)$$

$$\sigma_{\Lambda\Lambda_0}(s) := \sigma_{\Lambda_0}(s) - \sigma_\Lambda(s), \quad \sigma_\lambda(s) := \exp\left(-\frac{s^2}{\lambda^4}\right). \quad (37)$$

For shortness we will also write $C_{\Lambda\Lambda_0}^{-1}$ instead of $(C^{\Lambda\Lambda_0})^{-1}$. The parameters Λ, Λ_0 , such that $0 < \Lambda \leq \Lambda_0$, are respectively IR and UV cutoffs.

Definition 1 *Let $d\mu_{\Lambda\Lambda_0}$ be the measure defined by*

$$d\mu_{\Lambda\Lambda_0}(A, B, c, \bar{c}) := d\nu_{\Lambda\Lambda_0}(A, B - i\frac{1}{\xi} \partial A, c, \bar{c}). \quad (38)$$

For $\Phi = (A, B, c, \bar{c})$ and $K = (j, b, \bar{\eta}, \eta)$ and an infinitesimal variation $\delta\Phi = (\delta A, 0, \delta c, \delta \bar{c})$, using the properties of Gaussian measures from appendix A, we have

$$d\mu_{\Lambda\Lambda_0}(\Phi + \delta\Phi) = d\mu_{\Lambda\Lambda_0}(\Phi) \left(1 + \frac{1}{\hbar} \langle \bar{c}, S_{\Lambda\Lambda_0}^{-1} \delta c \rangle + \frac{1}{\hbar} \langle \delta \bar{c}, S_{\Lambda\Lambda_0}^{-1} c \rangle - \frac{1}{\hbar} \langle A, C_{\Lambda\Lambda_0}^{-1} \delta A \rangle - \frac{1}{\hbar} \langle i\partial(B - i\frac{1}{\xi}\partial A), \delta A \rangle \right). \quad (39)$$

Definition 2 The free partition function $Z_0^{\Lambda\Lambda_0}$ is defined by

$$Z_0^{\Lambda\Lambda_0}(K) := \int d\mu_{\Lambda\Lambda_0}(\Phi) e^{\frac{1}{\hbar} K \cdot \Phi}. \quad (40)$$

It follows that

$$Z_0^{\Lambda\Lambda_0}(j, b, \bar{\eta}, \eta) = \chi^{\Lambda\Lambda_0}(-i j - \frac{1}{\xi} \partial b, -i b, -i \bar{\eta}, -i \eta), \quad (41)$$

$$Z_0^{\Lambda\Lambda_0}(K) = e^{\frac{1}{2\hbar} \langle K, \mathbf{C}^{\Lambda\Lambda_0} K \rangle}, \quad (42)$$

where $\mathbf{C}^{\Lambda\Lambda_0}$ is a 7x7 matrix,

$$\mathbf{C}^{\Lambda\Lambda_0} := \begin{pmatrix} C_{\mu\nu}^{\Lambda\Lambda_0}, & S^{\Lambda\Lambda_0} p_\mu, & 0, & 0 \\ -S^{\Lambda\Lambda_0} p_\nu, & \frac{1}{\xi}(1 - \sigma_{\Lambda\Lambda_0}), & 0, & 0 \\ 0, & 0, & 0, & -S^{\Lambda\Lambda_0} \\ 0, & 0, & S^{\Lambda\Lambda_0}, & 0 \end{pmatrix}. \quad (43)$$

For $\Lambda < \Lambda_0$, $\mathbf{C}^{\Lambda\Lambda_0}$ is invertible:

$$\mathbf{C}_{\Lambda\Lambda_0}^{-1} = \begin{pmatrix} C_{\Lambda\Lambda_0; \mu\nu}^{-1} - \frac{1}{\xi} p_\mu p_\nu, & -p_\mu, & 0, & 0 \\ p_\nu, & \xi, & 0, & 0 \\ 0, & 0, & 0, & S_{\Lambda\Lambda_0}^{-1} \\ 0, & 0, & -S_{\Lambda\Lambda_0}^{-1}, & 0 \end{pmatrix}. \quad (44)$$

We write the bosonic part of $\mathbf{C}_{\Lambda\Lambda_0}^{-1}$ as $\mathbf{P}^T \mathbf{Q}^{\Lambda\Lambda_0} \mathbf{P}$, where \mathbf{P} is a diagonal matrix with $\mathbf{P}_{\mu\mu} = |p|$ for each index $\mu \in \{1, \dots, 4\}$ and $\mathbf{P}_{55} = 1$. The eigenvalues q_α of $\mathbf{Q}^{\Lambda\Lambda_0}$ are:

$$q_{1,2,3} = \sigma_{\Lambda\Lambda_0}^{-1}, \quad q_{4,5} = \frac{\xi_{\Lambda\Lambda_0} \pm (\xi_{\Lambda\Lambda_0}^2 - 4\sigma_{\Lambda\Lambda_0}^{-1})^{\frac{1}{2}}}{2}, \quad \xi_{\Lambda\Lambda_0} := \xi + \frac{1}{\xi}(\sigma_{\Lambda\Lambda_0}^{-1} - 1).$$

The fact that the real part of these eigenvalues is positive is known to be a prerequisite for the definition of a Gaussian measure for the bosonic part of the theory.

A useful relation follows from (32), (41):

$$\int d\mu_{\Lambda\Lambda_0}(\Phi) B e^{\frac{1}{\hbar}K\cdot\Phi} = \frac{1}{\xi} \int d\mu_{\Lambda\Lambda_0}(\Phi) (b + i\partial A) e^{\frac{1}{\hbar}K\cdot\Phi}. \quad (45)$$

Definition 3 *The partition function $Z^{\Lambda\Lambda_0}$ of $SU(2)$ Yang–Mills field theory is given by*

$$Z^{\Lambda\Lambda_0}(K) := \int d\mu_{\Lambda\Lambda_0}(\Phi) e^{-\frac{1}{\hbar}L^{\Lambda_0\Lambda_0}} e^{\frac{1}{\hbar}K\cdot\Phi}, \quad L^{\Lambda_0\Lambda_0} := \int d^4x \mathcal{L}^{\Lambda_0\Lambda_0}. \quad (46)$$

The interaction Lagrangian density $\mathcal{L}^{\Lambda_0\Lambda_0} := \mathcal{L}_0^{\Lambda_0\Lambda_0} + \mathcal{L}_{ct}^{\Lambda_0\Lambda_0}$ is given by (30) and

$$\mathcal{L}_0^{\Lambda_0\Lambda_0} := g\epsilon_{abc}(\partial_\mu A_\nu^a)A_\mu^b A_\nu^c + \frac{g^2}{4}\epsilon_{cab}\epsilon_{cda}A_\mu^a A_\nu^b A_\mu^d A_\nu^s - g\epsilon_{abc}(\partial_\mu \bar{c}^a)A_\mu^b c^c. \quad (47)$$

Since we restrict to perturbation theory, all generating functionals are formal series in terms of \hbar and of their source/field arguments. We also emphasize that $\mathcal{L}_0^{\Lambda_0\Lambda_0}$ does not depend on the B field.

From the expression for $\mathbf{C}_{0\infty}^{-1}$ and equation (47) one deduces the full tree level Lagrangian density

$$\mathcal{L}_0^{tot} = \frac{1}{4}F_{\mu\nu}F_{\mu\nu} + \frac{\xi}{2}B^2 - iB\partial_\mu A_\mu - \partial_\mu \bar{c}D_\mu c. \quad (48)$$

The density \mathcal{L}_0^{tot} is invariant under the infinitesimal BRST transformation [BRS75], [Tyu75]

$$\delta^{BRS} A = \epsilon Dc, \quad \delta^{BRS} c = \epsilon \frac{1}{2}ig\{c, c\}, \quad \delta^{BRS} \bar{c} = \epsilon iB, \quad \delta^{BRS} B = 0, \quad (49)$$

where ϵ is a Grassmann parameter, and $\{c, c\}^d = i\epsilon_{abd}c^a c^b$. Defining the classical operator s by $\delta^{BRS}\Phi = \epsilon s\Phi$ one shows that s is nilpotent.

Definition 4 *The generating functional of the Connected Schwinger (CS) functions is*

$$e^{\frac{1}{\hbar}W^{\Lambda\Lambda_0}(K)} := Z^{\Lambda\Lambda_0}(K). \quad (50)$$

The derivation of the FE is usually given considering the generating functional L^{Λ_0} of the Connected Amputated Schwinger (CAS) functions.

Definition 5 *The generating functional L^{Λ_0} of CAS functions is*

$$e^{-\frac{1}{\hbar}L^{\Lambda_0}(\Phi)} := \int d\mu_{\Lambda_0}(\Phi') e^{-\frac{1}{\hbar}L^{\Lambda_0\Lambda_0}(\Phi'+\Phi)}. \quad (51)$$

From definition 2 we have for any polynomial $P(\Phi)$

$$\frac{d}{d\Lambda} \int d\mu_{\Lambda_0}(\Phi) P(\Phi) = \hbar \int d\mu_{\Lambda_0}(\Phi) \left\langle \frac{\delta}{\delta\Phi}, \hat{\mathbf{1}}\dot{\mathbf{C}}^{\Lambda_0} \frac{\delta}{\delta\Phi} \right\rangle P(\Phi), \quad (52)$$

where

$$(\hat{\mathbf{1}})_{\phi\phi'} = \begin{cases} -\delta_{\phi\phi'} & \text{if } \phi, \phi' \in \{c, \bar{c}\}, \\ \delta_{\phi\phi'} & \text{otherwise.} \end{cases} \quad (53)$$

Using equation (52) one obtains the FE, see e.g. [GJ87], [KKS92],

$$\dot{L}^{\Lambda_0}(\Phi) = \frac{\hbar}{2} \left\langle \frac{\delta}{\delta\Phi}, \hat{\mathbf{1}}\dot{\mathbf{C}}^{\Lambda_0} \frac{\delta}{\delta\Phi} \right\rangle L^{\Lambda_0} - \frac{1}{2} \left\langle \frac{\delta L^{\Lambda_0}}{\delta\Phi}, \hat{\mathbf{1}}\dot{\mathbf{C}}^{\Lambda_0} \frac{\delta L^{\Lambda_0}}{\delta\Phi} \right\rangle. \quad (54)$$

From appendix A on Gaussian measures it follows that

$$d\mu_{\Lambda_0}(\Phi - \mathbf{C}^{\Lambda_0} \delta\Phi) = d\mu_{\Lambda_0}(\Phi) e^{-\frac{1}{2\hbar} \langle \delta\Phi, \mathbf{C}^{\Lambda_0} \delta\Phi \rangle} e^{\frac{1}{\hbar} \langle \delta\Phi, \Phi \rangle}. \quad (55)$$

This gives the relation between the generating functionals W^{Λ_0} and L^{Λ_0}

$$W^{\Lambda_0}(K) = \frac{1}{2} \langle K, \mathbf{C}^{\Lambda_0} K \rangle - L^{\Lambda_0}(\hat{\mathbf{1}}_{\bar{c}} \mathbf{C}^{\Lambda_0} K), \quad (56)$$

where

$$(\hat{\mathbf{1}}_{\bar{c}})_{\phi\phi'} = \begin{cases} -1 & \text{if } \phi = \phi' = \bar{c}, \\ \delta_{\phi\phi'} & \text{otherwise,} \end{cases} \quad (\hat{\mathbf{1}}_c)_{\phi\phi'} = \begin{cases} -1 & \text{if } \phi = \phi' = c, \\ \delta_{\phi\phi'} & \text{otherwise.} \end{cases} \quad (57)$$

Definition 6 (Legendre transform) *For $0 < \Lambda < \Lambda_0$, let $K^{\Lambda_0}(\underline{\Phi})$ be a solution of the system of equations*

$$\underline{A} - \frac{\delta W^{\Lambda_0}}{\delta j} \Big|_{K^{\Lambda_0}(\underline{\Phi})} = 0, \quad \underline{B} - \frac{\delta W^{\Lambda_0}}{\delta b} \Big|_{K^{\Lambda_0}(\underline{\Phi})} = 0, \quad (58)$$

$$\underline{c} - \frac{\delta W^{\Lambda_0}}{\delta \bar{\eta}} \Big|_{K^{\Lambda_0}(\underline{\Phi})} = 0, \quad \underline{\bar{c}} + \frac{\delta W^{\Lambda_0}}{\delta \eta} \Big|_{K^{\Lambda_0}(\underline{\Phi})} = 0. \quad (59)$$

The effective action is

$$\Gamma^{\Lambda_0}(\underline{\Phi}) := K^{\Lambda_0}(\underline{\Phi}) \cdot \underline{\Phi} - W^{\Lambda_0}(K^{\Lambda_0}(\underline{\Phi})). \quad (60)$$

A solution of the system of equations (58), (59) always exists as a formal series in \hbar and in the fields. Same statement for the system (61).

Definition 7 For $0 < \Lambda \leq \Lambda_0$, let $\Phi^{\Lambda\Lambda_0}(\underline{\Phi})$ be a solution of the equation

$$\underline{\Phi} = \left(\Phi - \mathbf{C}^{\Lambda\Lambda_0} \frac{\delta L^{\Lambda\Lambda_0}}{\delta \Phi} \right) \Big|_{\Phi^{\Lambda\Lambda_0}(\underline{\Phi})}. \quad (61)$$

The reduced effective action is

$$\Gamma^{\Lambda\Lambda_0}(\underline{\Phi}) := \left(L^{\Lambda\Lambda_0}(\Phi) - \frac{1}{2} \left\langle \frac{\delta L^{\Lambda\Lambda_0}}{\delta \Phi}, \hat{\mathbf{1}} \mathbf{C}^{\Lambda\Lambda_0} \frac{\delta L^{\Lambda\Lambda_0}}{\delta \Phi} \right\rangle \right) \Big|_{\Phi^{\Lambda\Lambda_0}(\underline{\Phi})}. \quad (62)$$

From definitions 6 and 7 it follows that for $0 < \Lambda < \Lambda_0$

$$\Gamma^{\Lambda\Lambda_0}(\underline{\Phi}) = \Gamma^{\Lambda\Lambda_0}(\underline{\Phi}) - \frac{1}{2} \langle \underline{\Phi}, \mathbf{C}_{\Lambda\Lambda_0}^{-1} \underline{\Phi} \rangle, \quad (63)$$

$$\Phi^{\Lambda\Lambda_0}(\underline{\Phi}) = \hat{\mathbf{1}}_{\bar{c}} \mathbf{C}^{\Lambda\Lambda_0} K^{\Lambda\Lambda_0}(\underline{\Phi}), \quad (64)$$

$$\Gamma^{\Lambda\Lambda_0}(\underline{\Phi}) = \left(L^{\Lambda\Lambda_0}(\Phi) - \frac{1}{2} \langle (\underline{\Phi} - \Phi), \mathbf{C}_{\Lambda\Lambda_0}^{-1} (\underline{\Phi} - \Phi) \rangle \right) \Big|_{\Phi^{\Lambda\Lambda_0}(\underline{\Phi})}. \quad (65)$$

Using equation (61) we see that, before the replacement, the rhs of (65), as a function of Φ for fixed $\underline{\Phi}$, has an extremum at $\Phi = \Phi^{\Lambda\Lambda_0}(\underline{\Phi})$. Applying ∂_Λ to equation (65), substituting \dot{L} with the rhs of (54) and using the property of extremum we obtain

$$\dot{\Gamma}^{\Lambda\Lambda_0}(\underline{\Phi}) = \frac{\hbar}{2} \left\langle \frac{\delta}{\delta \Phi}, \hat{\mathbf{1}} \dot{\mathbf{C}}^{\Lambda\Lambda_0} \frac{\delta}{\delta \Phi} \right\rangle L^{\Lambda\Lambda_0} \Big|_{\Phi^{\Lambda\Lambda_0}(\underline{\Phi})}. \quad (66)$$

Defining

$$W_{q,-k}^{\Lambda\Lambda_0} := (2\pi)^4 \frac{\delta^2 W^{\Lambda\Lambda_0}}{\delta K(-q) \delta K(k)}, \quad \Gamma_{q,-k}^{\Lambda\Lambda_0} := (2\pi)^4 \frac{\delta^2 \Gamma^{\Lambda\Lambda_0}}{\delta \underline{\Phi}(-q) \delta \underline{\Phi}(k)}, \quad (67)$$

it is easy to see that

$$\int d^4 k W_{q,-k}^{\Lambda\Lambda_0} \hat{\mathbf{1}}_{\bar{c}} \Gamma_{k,-p}^{\Lambda\Lambda_0} \hat{\mathbf{1}}_c = \delta(q-p). \quad (68)$$

This implies that, using similar notations,

$$L_{q,-p}^{\Lambda\Lambda_0} = \int d^4 k \Gamma_{q,-k}^{\Lambda\Lambda_0} (\delta(k-p) + \hat{\mathbf{1}} \mathbf{C}^{\Lambda\Lambda_0}(p) \Gamma_{-p,k}^{\Lambda\Lambda_0})^{-1}. \quad (69)$$

Eventually, one obtains the FE for the functional Γ and also for the functional Γ_χ with one operator insertion of source χ

$$\dot{\Gamma} = \frac{\hbar}{2} \langle \dot{\mathbf{C}} \delta_\phi \delta_{\bar{\phi}} \Gamma, (1 + \hat{\mathbf{1}} \mathbf{C} \delta_{\bar{\phi}} \delta_\phi \Gamma)^{-1} \rangle, \quad (70)$$

$$\delta_\chi \dot{\Gamma} = \frac{\hbar}{2} \langle \dot{\mathbf{C}} (1 + \delta_{\bar{\phi}} \delta_\phi \Gamma \mathbf{C} \hat{\mathbf{1}})^{-1} \delta_\chi \delta_{\bar{\phi}} \delta_\phi \Gamma, (1 + \hat{\mathbf{1}} \mathbf{C} \delta_{\bar{\phi}} \delta_\phi \Gamma)^{-1} \rangle, \quad (71)$$

where $\phi, \bar{\phi} \in \{\underline{A}, \underline{B}, \underline{c}, \underline{\bar{c}}\}$ and we omit appropriate sums over field labels. Generalization to $\vec{\chi} = (\chi_1, \dots, \chi_{n_\chi})$ with $n_\chi > 1$ is straightforward. The FE for Γ in modern form has been introduced in [Wet93, BDM93, Mor94, KKS97, KM09]. Flow equations with composite operator insertions have been introduced in [KK92].

The mass dimension of a vertex function $\Gamma_{\vec{\chi}}^{\phi;w}(\vec{p})$ with N fields $\phi_i \in \{A, B, c, \bar{c}\}$, n_χ insertions of sources χ_i and $\|w\|$ momentum derivatives is $d := 4 - \sum_i^{n_\chi} [\chi_i] - \sum_i^N [\phi_i] - \|w\|$, where $[F]$ stands for the mass dimension of F in position space. We say that such a term is **irrelevant** if $d < 0$, as for example $\Gamma_\gamma^{AAc;w}$ for $[\gamma] = 2$, and **relevant** otherwise. Furthermore, we call a relevant term **marginal** if $d = 0$, as for example Γ^{AAAA} , or **strictly relevant** if $d > 0$.

Expanding in formal power series in \hbar we have

$$\Gamma^{\Lambda\Lambda_0}(\underline{\Phi}) = \sum_{l=0}^{\infty} \hbar^l \Gamma_l^{\Lambda\Lambda_0}(\underline{\Phi}). \quad (72)$$

We also note that the FE (70) and (71) admit an inductive structure in the loop number. This property allows us to prove statements by induction, first establishing them at tree-level, then proving that if they hold up to loop order $l - 1 \geq 0$ they are also valid at order l .

The proposition that follows proves that vertex functions Γ involving B fields do vanish. We use the notation $B\vec{\phi}$ to denote sequences of field labels with $\phi_i \in \{A, B, c, \bar{c}\}$.

Proposition 8 *Assume vanishing renormalization conditions for all relevant terms with at least one B field:*

$$\Gamma_l^{B\vec{\phi};0\Lambda_0;w}(\vec{q}) = 0, \quad (73)$$

where \vec{q} is nonexceptional for marginal terms and vanishing otherwise. Then, for all $B\vec{\phi}$, l , w , \vec{p} , and $0 < \Lambda \leq \Lambda_0$,

$$\Gamma_l^{B\vec{\phi};\Lambda\Lambda_0;w}(\vec{p}) = 0. \quad (74)$$

Rank-2 marginal terms are tensors with two vector indices, say μ and ν , and can be decomposed in the basis $\{\delta_{\mu\nu}, q_{i\mu}q_{j\nu}\}$, where q_i, q_j range over all momenta in \vec{q} other than q_0 . For rank-2 marginal terms only the coefficient of $\delta_{\mu\nu}$ is set to zero in (73).

Proof We prove the statement by induction, increasing in the loop order, $l - 1 \mapsto l$. Given l , we proceed by descending from w_{max} in the number of derivatives, $\|w\| \mapsto \|w\| - 1$. For fixed l and w , all possible terms $\Gamma_l^{B\vec{\phi};w}$ are considered. By construction, for fixed l and $B\vec{\phi}$, the inductive scheme deals first with the irrelevant terms and continues, if they exist, with the marginal terms, followed by more and more relevant terms. The identity $\Gamma_0^{B\vec{\phi};\Lambda\Lambda_0;w}(\vec{p}) = 0$ follows from the definition of Γ . Assume that the statement of the theorem holds up to loop order $l - 1 \geq 0$. It follows that at order l the rhs of the FE for vertex functions with B fields vanishes. Using the FE we integrate the irrelevant terms from Λ_0 downward to arbitrary $\Lambda > 0$. The boundary conditions $\Gamma_l^{B\vec{\phi};\Lambda_0\Lambda_0;w}(\vec{p}) = 0$ and the vanishing of the rhs of the FE imply that all irrelevant terms with B fields vanish at loop order l . The boundary conditions $\Gamma_l^{B\vec{\phi};0\Lambda_0;w}(\vec{q}) = 0$ and the vanishing of the rhs of the FE imply that all marginal terms vanish at their renormalization point for arbitrary $\Lambda > 0$. The derivatives wrt momenta of marginal terms are irrelevant terms. Consequently, we conclude that the marginal $\Gamma_l^{B\vec{\phi};\Lambda\Lambda_0;w}(\vec{p})$ vanish for all \vec{p} . Similar arguments hold for all strictly relevant terms. ■ In the following we will always adopt the renormalization conditions (73). Consequently, counterterms involving B fields are not generated.

Let us denote by \tilde{W}, \tilde{Z} the functional W, Z with b set to zero.

$$\tilde{W}(j, \bar{\eta}, \eta) := W(j, 0, \bar{\eta}, \eta), \quad \tilde{Z}(j, \bar{\eta}, \eta) := Z(j, 0, \bar{\eta}, \eta). \quad (75)$$

The covariance matrix $\tilde{\mathbf{C}}$ is obtained from $\mathbf{C}(43)$ by removing the column and row which correspond to b .

Proposition 9

$$W^{\Lambda\Lambda_0}(j, b, \bar{\eta}, \eta) = \frac{1}{2\xi} \langle b, b \rangle + \tilde{W}^{\Lambda\Lambda_0}(j - i\frac{1}{\xi} \partial b, \bar{\eta}, \eta). \quad (76)$$

Proof Using the definition of the partition function $Z^{\Lambda\Lambda_0}$ one computes

$$Z^{\Lambda\Lambda_0}(j + i\frac{1}{\xi}\partial b, b, \bar{\eta}, \eta) = e^{-\frac{1}{\hbar}\mathbb{L}^{\Lambda_0\Lambda_0}} \chi^{\Lambda\Lambda_0}(ij, ib, i\bar{\eta}, i\eta), \quad (77)$$

$$\mathbb{L}^{\Lambda_0\Lambda_0} := L^{\Lambda_0\Lambda_0}\left(\frac{\delta}{\delta j}, \frac{\delta}{\delta \bar{\eta}}, \frac{\delta}{\delta \eta}\right). \quad (78)$$

From definition (32) it follows

$$Z^{\Lambda\Lambda_0}(j + i\frac{1}{\xi}\partial b, b, \bar{\eta}, \eta) = e^{\frac{1}{2\hbar\xi}\langle b, b \rangle} e^{-\frac{1}{\hbar}\mathbb{L}^{\Lambda_0\Lambda_0}} e^{-\frac{1}{\hbar}\langle \bar{\eta}, S^{\Lambda\Lambda_0}\eta \rangle + \frac{1}{2\hbar}\langle j, C^{\Lambda\Lambda_0}j \rangle}. \quad (79)$$

Observing that the expression multiplying $e^{\frac{1}{2\hbar\xi}\langle b, b \rangle}$ is $\tilde{Z}^{\Lambda\Lambda_0}(j, \bar{\eta}, \eta)$ we obtain

$$Z^{\Lambda\Lambda_0}(j, b, \bar{\eta}, \eta) = e^{\frac{1}{2\hbar\xi}\langle b, b \rangle} \tilde{Z}^{\Lambda\Lambda_0}(j - i\frac{1}{\xi}\partial b, \bar{\eta}, \eta). \quad (80)$$

Taking the logarithm finishes the proof. \blacksquare

Substitution of $W^{\Lambda\Lambda_0}$ (76) into the definition of $\Gamma^{\Lambda\Lambda_0}$ (60) and integration by parts give

$$\Gamma^{\Lambda\Lambda_0}(\underline{A}, \underline{B}, \underline{\bar{c}}, \underline{c}) = \langle b, \underline{B} - i\frac{1}{\xi}\partial\underline{A} - \frac{1}{2}b \rangle + \tilde{\Gamma}^{\Lambda\Lambda_0}(\underline{A}, \underline{\bar{c}}, \underline{c}), \quad (81)$$

$$\tilde{\Gamma}^{\Lambda\Lambda_0}(\underline{A}, \underline{\bar{c}}, \underline{c}) := \langle j, \underline{A} \rangle + \langle \bar{\eta}, \underline{c} \rangle + \langle \underline{\bar{c}}, \eta \rangle - \tilde{W}^{\Lambda\Lambda_0}(j, \bar{\eta}, \eta) \Big|_{K^{\Lambda\Lambda_0}(\underline{\Phi})}. \quad (82)$$

From definition 6 it follows that $b = \xi\underline{B} - i\partial\underline{A}$. Consequently, the above expression becomes

$$\Gamma^{\Lambda\Lambda_0}(\underline{A}, \underline{B}, \underline{\bar{c}}, \underline{c}) = \frac{1}{2\xi}\langle \xi\underline{B} - i\partial\underline{A}, \xi\underline{B} - i\partial\underline{A} \rangle + \tilde{\Gamma}^{\Lambda\Lambda_0}(\underline{A}, \underline{\bar{c}}, \underline{c}). \quad (83)$$

Differentiation wrt \underline{A} yields an important identity:

Corollary 10

$$\frac{\delta\Gamma^{\Lambda\Lambda_0}}{\delta\underline{A}_\mu^a} = i\partial_\mu(\underline{B}^a - i\frac{1}{\xi}\partial\underline{A}^a) + \frac{\delta\tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta\underline{A}_\mu^a}. \quad (84)$$

For $\tilde{\Phi} := (A_\mu^a, c^a, \bar{c}^a)$ the functional $\tilde{\Gamma}^{\Lambda\Lambda_0}(\tilde{\Phi})$ is defined from $L^{\Lambda\Lambda_0}(\tilde{\Phi})$ in analogy with (62). For $0 < \Lambda < \Lambda_0$, it follows that

$$\tilde{\Gamma}^{\Lambda\Lambda_0}(\tilde{\Phi}) = \tilde{\Gamma}^{\Lambda\Lambda_0}(\tilde{\Phi}) + \frac{1}{2}\langle \tilde{\Phi}, \tilde{C}_{\Lambda\Lambda_0}^{-1} \tilde{\Phi} \rangle. \quad (85)$$

Substitution (63) and (85) into (83) yields $\Gamma^{\Lambda\Lambda_0}(\underline{\Phi}) = \tilde{\Gamma}^{\Lambda\Lambda_0}(\tilde{\underline{\Phi}})$. Note also that at $\Lambda = \Lambda_0$ we have $\underline{\Phi} = \tilde{\underline{\Phi}}$ and

$$\Gamma^{\Lambda_0\Lambda_0}(\underline{\Phi}) = L^{\Lambda_0\Lambda_0}(\Phi) = L^{\Lambda_0\Lambda_0}(\tilde{\Phi}) = \tilde{\Gamma}^{\Lambda_0\Lambda_0}(\tilde{\underline{\Phi}}) = \int d^4x (\mathcal{L}_0^{\Lambda_0\Lambda_0} + \mathcal{L}_{ct}^{\Lambda_0\Lambda_0}). \quad (86)$$

2.4 Violated Slavnov–Taylor identities

We are working in a framework where gauge invariance is broken already in the classical lagrangian due to the gauge fixing term. It has then been realized that invariance of the lagrangian under the BRST transformations ensures the gauge invariance of physical quantities to be calculated from the theory [Nie75, PS85]. On the level of correlation functions (Green’s functions in the relativistic theory) this invariance leads to a system of identities between different correlation functions which are called Slavnov Taylor identities (STI) [ZJ75, Sla72, Tay71]. These identities may be used to argue that physical quantities obtained from these functions, as for example the pole of the propagators for all physical fields of the Standard Model, are gauge invariant [GG00].

In our framework gauge invariance is also violated in an even more serious way by the presence of the regulators in (35), (36). We want to show that for a suitable class of renormalization conditions, which does not restrict the freedom in fixing the physical coupling constant and the normalization of the fields, gauge invariance can be recovered in the renormalized theory. This means we want to show that the STI hold once we take the limits $\Lambda \rightarrow 0$ and $\Lambda_0 \rightarrow \infty$.

The first step is then to write a system of violated STI suitable for our subsequent analysis of their restoration. To do so we thus analyse the behavior of the regularized generating functionals of the correlation functions under BRST transformations. The infinitesimal BRST transformations can be generated by composite operator insertions for which we also have a freedom of normalization, as encoded by the constants R_i introduced below [KM09].

We derive the *violated* STI setting $\Lambda = 0$. Remember the comments in the beginning of section 2.3 as regards the implications of this choice for the generating functionals on the IR side. We consider the functional $Z_{vst}^{0\Lambda_0}$ defined with the modified Lagrangian density

$$\mathcal{L}_{vst}^{\Lambda_0\Lambda_0} := \mathcal{L}^{\Lambda_0\Lambda_0} + \gamma\psi^{\Lambda_0} + \omega\Omega^{\Lambda_0}, \quad L_{vst}^{\Lambda_0\Lambda_0} := \int d^4x \mathcal{L}_{vst}^{\Lambda_0\Lambda_0}, \quad (87)$$

where γ, ω are external sources, and

$$\psi^{\Lambda_0} := R_1^{\Lambda_0} \partial c - ig R_2^{\Lambda_0} [A, c], \quad \Omega^{\Lambda_0} := \frac{1}{2i} g R_3^{\Lambda_0} \{c, c\}, \quad R_i^{\Lambda_0} = 1 + O(\hbar). \quad (88)$$

The requirement that at zero loop order ψ^{Λ_0} and Ω^{Λ_0} correspond to the classical BRST variation implies that the constants $R_i^{\Lambda_0}$ are equal to one at tree level. At higher orders we admit counterterms for the BRST transformation [KM09]. Performing the change of variables $\Phi \mapsto \Phi + \delta_\epsilon \Phi$ we obtain the identity

$$-\hbar \int \delta_\epsilon (d\mu_{0\Lambda_0} e^{-\frac{1}{\hbar} L_{vst}^{\Lambda_0 \Lambda_0}}) e^{\frac{1}{\hbar} K \cdot \Phi} = \int d\mu_{0\Lambda_0} e^{-\frac{1}{\hbar} L_{vst}^{\Lambda_0 \Lambda_0} + \frac{1}{\hbar} K \cdot \Phi} K \cdot \delta_\epsilon(\Phi), \quad (89)$$

where

$$\delta_\epsilon A := \epsilon \sigma_{0\Lambda_0} * \psi^{\Lambda_0}, \quad \delta_\epsilon c := -\epsilon \sigma_{0\Lambda_0} * \Omega^{\Lambda_0}, \quad \delta_\epsilon \bar{c} := \epsilon \sigma_{0\Lambda_0} * iB. \quad (90)$$

Using formula (39) for an infinitesimal change of variables and substituting the variation $\delta_\epsilon \Phi$ with its explicit form (90) we have

$$-\hbar \delta_\epsilon (d\mu_{0\Lambda_0}(\Phi)) = d\mu_{0\Lambda_0}(\Phi) \epsilon I_1^{\Lambda_0}(\Phi), \quad (91)$$

$$-\hbar \delta_\epsilon (e^{-\frac{1}{\hbar} L_{vst}^{\Lambda_0 \Lambda_0}(\Phi)}) = e^{-\frac{1}{\hbar} L_{vst}^{\Lambda_0 \Lambda_0}(\Phi)} \epsilon I_2^{\Lambda_0}(\Phi), \quad (92)$$

where

$$\begin{aligned} I_1^{\Lambda_0}(\Phi) &:= \langle A, C^{-1} \psi^{\Lambda_0} \rangle - \langle \bar{c}, S^{-1} \Omega^{\Lambda_0} \rangle - i \langle B, S^{-1} c \rangle \\ &\quad + i \langle \partial(B - i \frac{1}{\xi} \partial A), \sigma_{0\Lambda_0} * \psi^{\Lambda_0} \rangle, \end{aligned} \quad (93)$$

$$\begin{aligned} I_2^{\Lambda_0}(\Phi) &:= \langle \sigma_{0\Lambda_0} * \frac{\delta L^{\Lambda_0 \Lambda_0}}{\delta A}, \psi^{\Lambda_0} \rangle - \langle \sigma_{0\Lambda_0} * \frac{\delta L^{\Lambda_0 \Lambda_0}}{\delta c}, \Omega^{\Lambda_0} \rangle \\ &\quad + i \langle B, \sigma_{0\Lambda_0} * \frac{\delta L^{\Lambda_0 \Lambda_0}}{\delta \bar{c}} \rangle + \langle \gamma, Q_{\rho\gamma}^{\Lambda_0} \rangle + \langle \omega, Q_{\rho\omega}^{\Lambda_0} \rangle, \end{aligned} \quad (94)$$

$$\begin{aligned} Q_{\rho\gamma}^{\Lambda_0} &:= g(R_3^{\Lambda_0} - R_2^{\Lambda_0}) \epsilon^{abs} (\sigma_{0\Lambda_0} * \psi^{\Lambda_0})_\mu^b c^s \\ &\quad + g R_1^{\Lambda_0} R_3^{\Lambda_0} \epsilon^{abs} (c^s (\sigma_{0\Lambda_0} * \partial_\mu c^b) - \sigma_{0\Lambda_0} * (c^s \partial_\mu c^b)) \\ &\quad + g^2 R_2^{\Lambda_0} R_3^{\Lambda_0} \epsilon^{ase} \epsilon^{ebt} ((\sigma_{0\Lambda_0} * (A_\mu^b c^t)) c^s - A_\mu^b (\sigma_{0\Lambda_0} * (c^t c^s))), \end{aligned} \quad (95)$$

$$Q_{\rho\omega}^{\Lambda_0} := (g R_3^{\Lambda_0})^2 c^s (\sigma_{0\Lambda_0} * (c^a c^s)). \quad (96)$$

The terms $Q_{\rho\gamma}^{\Lambda_0}$ and $Q_{\rho\omega}^{\Lambda_0}$ originate from the variations $\delta_\epsilon(\gamma\psi^{\Lambda_0})$, $\delta_\epsilon(\omega\Omega^{\Lambda_0})$, see (87). Substituting $B \mapsto \frac{1}{\xi}(b + i\partial A)$, see (45), in both terms $I_i^{\Lambda_0}$ we get

$$\int d\mu_{0\Lambda_0} I_i^{\Lambda_0} e^{-\frac{1}{\hbar}L_{vst}^{\Lambda_0\Lambda_0} + \frac{1}{\hbar}K\cdot\Phi} = \int d\mu_{0\Lambda_0} J_i^{\Lambda_0} e^{-\frac{1}{\hbar}L_{vst}^{\Lambda_0\Lambda_0} + \frac{1}{\hbar}K\cdot\Phi}, \quad (97)$$

where

$$\begin{aligned} J_1^{\Lambda_0}(\Phi) &:= \langle A, C^{-1}\psi^{\Lambda_0} \rangle - \langle \bar{c}, S^{-1}\Omega^{\Lambda_0} \rangle + \frac{1}{\xi} \langle \partial A, S^{-1}c \rangle \\ &\quad - i\frac{1}{\xi} \langle b, S^{-1}c + \sigma_{0\Lambda_0} * \partial\psi^{\Lambda_0} \rangle, \end{aligned} \quad (98)$$

$$\begin{aligned} J_2^{\Lambda_0}(\Phi) &:= \langle \sigma_{0\Lambda_0} * \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta A}, \psi^{\Lambda_0} \rangle - \langle \sigma_{0\Lambda_0} * \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta c}, \Omega^{\Lambda_0} \rangle \\ &\quad + i\frac{1}{\xi} \langle b + i\partial A, \sigma_{0\Lambda_0} * \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta \bar{c}} \rangle + \langle \gamma, Q_{\rho\gamma}^{\Lambda_0} \rangle + \langle \omega, Q_{\rho\omega}^{\Lambda_0} \rangle. \end{aligned} \quad (99)$$

Introducing the operators $Q_\rho^{\Lambda_0}$ and $Q_\beta^{\Lambda_0}$,

$$\begin{aligned} Q_\rho^{\Lambda_0} &:= \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta A_\mu} \sigma_{0\Lambda_0} * \psi_\mu^{\Lambda_0} - \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta c} \sigma_{0\Lambda_0} * \Omega^{\Lambda_0} - \frac{1}{\xi} \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta \bar{c}} \sigma_{0\Lambda_0} * \partial A \\ &\quad + AC^{-1}\psi^{\Lambda_0} - \bar{c}S^{-1}\Omega^{\Lambda_0} + \frac{1}{\xi} \partial AS^{-1}c, \end{aligned} \quad (100)$$

$$Q_\beta^{\Lambda_0} := \sigma_{0\Lambda_0} * \left(\frac{\delta L^{\Lambda_0\Lambda_0}}{\delta \bar{c}} - \partial\psi^{\Lambda_0} \right) - S^{-1}c, \quad (101)$$

we have

$$J_1^{\Lambda_0} + J_2^{\Lambda_0} = i\frac{1}{\xi} \langle b, Q_\beta^{\Lambda_0} \rangle + \langle \gamma, Q_{\rho\gamma}^{\Lambda_0} \rangle + \langle \omega, Q_{\rho\omega}^{\Lambda_0} \rangle + \int d^4x Q_\rho^{\Lambda_0}(x). \quad (102)$$

We may now express the lhs of (89) as

$$\epsilon \int d\mu_{0\Lambda_0} (J_1^{\Lambda_0} + J_2^{\Lambda_0}) e^{-\frac{1}{\hbar}L_{vst}^{\Lambda_0\Lambda_0} + \frac{1}{\hbar}K\cdot\Phi} = -\epsilon \mathcal{D}Z_{aux}^{0\Lambda_0} \Big|_{\rho, \beta=0}, \quad (103)$$

where

$$\mathcal{D} := \int d^4x \hbar \frac{\delta}{\delta \rho(x)} + i\frac{1}{\xi} \langle b, \hbar \frac{\delta}{\delta \beta} \rangle, \quad (104)$$

$$Z_{aux}^{0\Lambda_0} := \int d\mu_{0\Lambda_0}(\Phi) e^{-\frac{1}{\hbar}L_{aux}^{\Lambda_0\Lambda_0} + \frac{1}{\hbar}K\cdot\Phi}, \quad (105)$$

$$\mathcal{L}_{aux}^{\Lambda_0\Lambda_0} := \mathcal{L}_{vst}^{\Lambda_0\Lambda_0} + \rho Q_\rho^{\Lambda_0} + \rho \gamma_\mu^a Q_{\rho\gamma_\mu^a}^{\Lambda_0} + \rho \omega^a Q_{\rho\omega^a}^{\Lambda_0} + \beta Q_\beta^{\Lambda_0}. \quad (106)$$

Defining

$$\mathcal{S} := \langle j, \sigma_{0\Lambda_0} * \hbar \frac{\delta}{\delta \gamma} \rangle + \langle \bar{\eta}, \sigma_{0\Lambda_0} * \hbar \frac{\delta}{\delta \omega} \rangle - i \langle \sigma_{0\Lambda_0} * \hbar \frac{\delta}{\delta b}, \eta \rangle, \quad (107)$$

we write the rhs of (89) in the following form

$$\int d\mu_{0\Lambda_0} e^{-\frac{1}{\hbar} L_{vst}^{0\Lambda_0} + \frac{1}{\hbar} K \cdot \Phi} K \cdot \delta_\epsilon \Phi = -\epsilon \mathcal{S} Z_{vst}^{0\Lambda_0} = -\epsilon \mathcal{S} Z_{aux}^{0\Lambda_0} \Big|_{\rho, \beta=0}. \quad (108)$$

Using equations (103) and (108) we write identity (89) as follows

$$\mathcal{D} Z_{aux}^{0\Lambda_0} \Big|_{\rho, \beta=0} = \mathcal{S} Z_{aux}^{0\Lambda_0} \Big|_{\rho, \beta=0}. \quad (109)$$

Then we perform the Legendre transform on the auxiliary functional $W_{aux} := \hbar \log Z_{aux}$. This defines the functional Γ_{aux} , from which are obtained $\tilde{\Gamma}_{aux}$, see (83), and the functionals

$$\tilde{\Gamma}_\chi^{0\Lambda_0} := \frac{\delta \tilde{\Gamma}_{aux}^{0\Lambda_0}}{\delta \chi} \Big|_{\chi=0}, \quad \tilde{\Gamma}_1^{0\Lambda_0} := \int d^4 x \tilde{\Gamma}_{\rho(x)}^{0\Lambda_0}, \quad (110)$$

where $\chi \in \{\rho, \beta\}$ and γ, ω are arbitrary. Equation (109) then gives

$$\tilde{\Gamma}_1^{0\Lambda_0} + \langle (i\underline{B} + \frac{1}{\xi} \partial \underline{A}), \tilde{\Gamma}_\beta^{0\Lambda_0} \rangle = \frac{1}{2} \mathcal{S} \underline{\Gamma}^{0\Lambda_0}, \quad (111)$$

where

$$\underline{\Gamma}^{0\Lambda_0} := i \langle \underline{B}, \bar{\omega} \rangle + \tilde{\Gamma}^{0\Lambda_0}, \quad \tilde{\Gamma}^{0\Lambda_0} := \tilde{\Gamma}_1^{0\Lambda_0} + \frac{1}{2\xi} \langle \underline{A}, \partial \partial \underline{A} \rangle, \quad (112)$$

and $\mathcal{S} := \mathcal{S}_{\tilde{c}} + \mathcal{S}_A - \mathcal{S}_c$, with

$$\mathcal{S}_A := \left\langle \frac{\delta \underline{\Gamma}^{0\Lambda_0}}{\delta \underline{A}}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \gamma} \right\rangle + \left\langle \frac{\delta \underline{\Gamma}^{0\Lambda_0}}{\delta \gamma}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \underline{A}} \right\rangle, \quad (113)$$

$$\mathcal{S}_{\tilde{c}} := \left\langle \frac{\delta \underline{\Gamma}^{0\Lambda_0}}{\delta \tilde{c}}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \bar{\omega}} \right\rangle + \left\langle \frac{\delta \underline{\Gamma}^{0\Lambda_0}}{\delta \bar{\omega}}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \tilde{c}} \right\rangle, \quad (114)$$

$$\mathcal{S}_c := \left\langle \frac{\delta \underline{\Gamma}^{0\Lambda_0}}{\delta \underline{c}}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \omega} \right\rangle + \left\langle \frac{\delta \underline{\Gamma}^{0\Lambda_0}}{\delta \omega}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \underline{c}} \right\rangle, \quad (115)$$

$$\frac{\delta}{\delta \tilde{c}} := \frac{\delta}{\delta \underline{c}} - \partial \frac{\delta}{\delta \gamma}. \quad (116)$$

We rewrite equation (111) in the following form

$$\langle i\underline{B}, \tilde{\Gamma}_\beta^{0\Lambda_0} \rangle = \frac{1}{2} \mathcal{S}_{\tilde{c}} \Gamma^{0\Lambda_0} = \langle i\underline{B}, \sigma_{0\Lambda_0} * \frac{\delta}{\delta \tilde{c}} \tilde{\Gamma}^{0\Lambda_0} \rangle, \quad (117)$$

$$\tilde{F}_1^{0\Lambda_0} = \frac{1}{2} \tilde{S} \tilde{\Gamma}^{0\Lambda_0}, \quad (118)$$

where

$$\tilde{S} := \mathcal{S}_A - \mathcal{S}_c, \quad \tilde{F}_1^{0\Lambda_0} := \tilde{\Gamma}_1^{0\Lambda_0} + \frac{1}{\xi} \langle \partial A, \tilde{\Gamma}_\beta^{0\Lambda_0} \rangle. \quad (119)$$

The introduction of the functional $\tilde{\Gamma}$ leads to relation (118) and to the consistency condition given in (121) below. They are important in the analysis of the renormalization conditions for $\tilde{\Gamma}_1^{0\Lambda_0; \vec{\phi}}$, see section 5.

An algebraic computation shows that

$$(\mathcal{S}_i \mathcal{S}_j + \mathcal{S}_j \mathcal{S}_i) \Gamma^{0\Lambda_0} = 0, \quad \forall i, j \in \{A, c, \tilde{c}\}. \quad (120)$$

Consequently $\tilde{S}^2 \tilde{\Gamma}^{0\Lambda_0} = 0$. Thus application of the operator \tilde{S} to equation (118) yields

$$\tilde{S} \tilde{F}_1^{0\Lambda_0} = 0. \quad (121)$$

Using again (120), we also have

$$\frac{1}{2} (\tilde{S} \mathcal{S}_{\tilde{c}} + \mathcal{S}_{\tilde{c}} \tilde{S}) \Gamma^{0\Lambda_0} = 0, \quad \text{and thus} \quad \tilde{S} \tilde{\Gamma}_\beta + \sigma_{0\Lambda_0} * \frac{\delta}{\delta \tilde{c}} \tilde{F}_1^{0\Lambda_0} = 0. \quad (122)$$

Finally we set $\gamma, \omega = 0$ in (117), (118) to get the AGE and the STI:

$$\tilde{\Gamma}_\beta^{0\Lambda_0} = \sigma_{0\Lambda_0} * \left(\frac{\delta \tilde{\Gamma}^{0\Lambda_0}}{\delta \tilde{c}} - \partial \tilde{\Gamma}_\gamma^{0\Lambda_0} \right) \quad (AGE), \quad (123)$$

$$\tilde{F}_1^{0\Lambda_0} = \left\langle \frac{\delta \tilde{\Gamma}^{0\Lambda_0}}{\delta \underline{A}}, \sigma_{0\Lambda_0} * \tilde{\Gamma}_\gamma^{0\Lambda_0} \right\rangle - \left\langle \frac{\delta \tilde{\Gamma}^{0\Lambda_0}}{\delta \tilde{c}}, \sigma_{0\Lambda_0} * \tilde{\Gamma}_\omega^{0\Lambda_0} \right\rangle \quad (STI). \quad (124)$$

The goal is to show that $\tilde{\Gamma}_\beta^{0\infty} = 0$ and $\tilde{\Gamma}_1^{0\infty} = 0$, in the sense of theorems 3, 4. In appendix G we give a version of these two equations valid for $0 < \Lambda < \Lambda_0$.

2.4.1 The Antighost Equation

In this section we extend the AGE for arbitrary $\Lambda \in [0, \Lambda_0]$. In notations (11), (75), (110) and with $\mathcal{L}_{aux}^{\Lambda_0\Lambda_0}$ given in (106) we have

$$- \tilde{W}_\beta^{\Lambda\Lambda_0} \tilde{Z}_{vst}^{\Lambda\Lambda_0} = -\hbar \frac{\delta}{\delta\beta} \tilde{Z}_{aux}^{\Lambda\Lambda_0} \Big|_{\rho,\beta=0} = \int d\tilde{\mu} Q_\beta^{\Lambda_0} e^{-\frac{1}{\hbar} L_{vst}^{\Lambda_0\Lambda_0} + \frac{1}{\hbar} \tilde{K} \cdot \tilde{\Phi}}, \quad (125)$$

where $d\tilde{\mu} := d\mu_{\Lambda\Lambda_0}(\tilde{\Phi})$, $\tilde{\Phi} := (A, c, \bar{c})$, $\tilde{K} := (j, \bar{\eta}, \eta)$. Integration by parts, see (345), yields

$$\int d\tilde{\mu} \left(\frac{\delta L_{vst}^{\Lambda_0\Lambda_0}}{\delta \bar{c}} - S_{\Lambda\Lambda_0}^{-1} c - \eta \right) e^{-\frac{1}{\hbar} L_{vst}^{\Lambda_0\Lambda_0} + \frac{1}{\hbar} \tilde{K} \cdot \tilde{\Phi}} = 0. \quad (126)$$

From definition (87) we also have

$$\frac{\delta L_{vst}^{\Lambda_0\Lambda_0}}{\delta \bar{c}} = \frac{\delta L^{\Lambda_0\Lambda_0}}{\delta \bar{c}}. \quad (127)$$

Substituting Q_β from (101) into (125), then using (126), (127) and putting $\gamma, \omega = 0$, we obtain

$$- \tilde{W}_\beta^{\Lambda\Lambda_0} = \sigma_{0\Lambda_0} * \left(S_{\Lambda\Lambda_0}^{-1} \frac{\delta \tilde{W}^{\Lambda\Lambda_0}}{\delta \bar{\eta}} + \eta + \partial \tilde{W}_\gamma^{\Lambda\Lambda_0} \right) - S^{-1} \frac{\delta \tilde{W}^{\Lambda\Lambda_0}}{\delta \bar{\eta}}. \quad (128)$$

Finally we perform the Legendre transform

$$\begin{aligned} \tilde{\Gamma}_\beta^{\Lambda\Lambda_0} &= \sigma_{0\Lambda_0} * \left(S_{\Lambda\Lambda_0}^{-1} \underline{c} + \frac{\delta \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \bar{c}} - \partial \tilde{\Gamma}_\gamma^{\Lambda\Lambda_0} \right) - S^{-1} \underline{c} \\ &= \sigma_{0\Lambda_0} * \left(\frac{\delta \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \bar{c}} - \partial \tilde{\Gamma}_\gamma^{\Lambda\Lambda_0} \right) - S^{-1} \underline{c}. \end{aligned} \quad (129)$$

One can also prove identity (129) by noting that its lhs and the Γ -functionals on the rhs obey FEs of form (71). Moreover, identity (123) implies that both sides have the same boundary condition at $\Lambda = 0$.

3 Momentum bounds

From now on we use the following conventions:

- $0 < \Lambda \leq \Lambda_0$, unless otherwise stated.

- M is a fixed mass parameter such that $0 < M \leq \Lambda_0$.
- We omit the tilda for all vertex functions and insertions, for example $\tilde{\Gamma}^{\vec{\phi};w} \mapsto \Gamma^{\vec{\phi};w}$, $\tilde{\Gamma}^{\vec{\phi};w} \mapsto \Gamma^{\vec{\phi};w}$, $\tilde{\Gamma}_{\chi}^{\vec{\phi};w} \mapsto \Gamma_{\chi}^{\vec{\phi};w}$.
- We use A , c , \bar{c} instead of \underline{A} , \underline{c} , $\underline{\bar{c}}$, respectively.
- A tensor monomial is a tensor product of Kronecker δ 's and momentum variables in $p := (p_1, \dots, p_{n-1})$, for example $\delta_{\mu\nu} p_{1\rho} p_{2\sigma}$. Let $\{\delta^s p^k\}$ be the set of all monomials being a product of s Kronecker δ 's and k momenta p_i and let $\{\delta^s p^k\}_r$ be the union of the sets $\{\delta^s p^k\}$ such that r equals the rank of the monomials: $r = 2s + k$. For example $\{\delta^2\}_4 = \{\delta_{\mu\nu} \delta_{\rho\sigma}, \delta_{\mu\sigma} \delta_{\rho\nu}, \delta_{\mu\rho} \delta_{\nu\sigma}\}$.
- For $\vec{p} \in \mathbb{P}_n$ we define the η -function by

$$\eta(\vec{p}) := \min_{S \in \wp_{n-1} \setminus \{\emptyset\}} (|\sum_{i \in S} p_i|, M). \quad (130)$$

where \wp_{n-1} denotes the power set of $[n-1]$ (the sum does not include $p_0 = -\sum_1^{n-1} p_i$). A momentum configuration \vec{p} is nonexceptional iff $\eta(\vec{p}) \neq 0$ and exceptional otherwise.

- For a fixed constant c such that $0 < c < 1$ we define

$$\mathbb{M}_n := \{\vec{p} \in \mathbb{P}_n : \eta(\vec{p}) > cM \text{ and } p_i^2 \leq M^2 \forall i \in [n-1]\}. \quad (131)$$

Every $\vec{p} \in \mathbb{M}_n$ is nonexceptional.

- $\forall n \geq 2$, a momentum configuration $\vec{p} \in \mathbb{M}_n$ is **symmetric** iff $\vec{p} \in \mathbb{M}_n^s$,

$$\mathbb{M}_n^s := \{\vec{p} \in \mathbb{M}_n : p_i p_j = \frac{M^2}{n-1} (n \delta_{ij} - 1) \forall i, j \in [n-1]\}. \quad (132)$$

- $\forall n \geq 3$, a momentum configuration $\vec{p} \in \mathbb{M}_n$ is **coplanar** iff $\vec{p} \in \mathbb{M}_n^{cp}$,

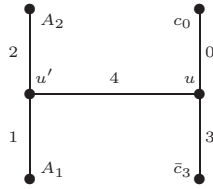
$$\mathbb{M}_n^{cp} := \{\vec{p} \in \mathbb{M}_n : \dim(\text{span}(p_0, \dots, p_{n-1})) = 2\}. \quad (133)$$

- In the following, a renormalization point is denoted by $\vec{q} \in \mathbb{P}_n$. See appendix H for the list of all relevant terms and their renormalization points.

3.1 Weighted trees

The bounds on the vertex functions presented in section 3.2 are expressed in terms of sets of weighted trees that are introduced by definitions 11, 12 below. As seen from (152), to each edge e of a weighted tree is associated a factor $(|p_e| + \Lambda)^{-\theta(e)}$, p_e being the momentum traversing the edge and $\theta(e)$ being the θ -weight of the edge, expressed as a sum of the ρ and σ -weights of the edge, see (139). The relation (141) expresses the fact that the total θ -weight of a tree is in agreement with power counting. Nonvanishing σ -weights are introduced in order to define viable tree bounds for momentum derived vertex functions. The definition of the σ -weight is inspired by how momentum derivatives are distributed along a tree, taking care of momentum conservation. Before giving the definition of the weighted trees we set up some necessary notations.

- A tree τ is a connected graph with no cycles. The sets of vertices and edges of a tree τ are denoted respectively by $V(\tau)$ and $E(\tau)$, or shortly V , E . In the following, the terms “edge“ and “line“ are equivalent.
- Let V_m be the set of vertices of valence m . Then, $V = \bigcup_{m \geq 1} V_m$.
- Let E_1 be the set of edges incident to vertices of valence 1. In other words, E_1 is the set of external edges.
- Each tree has a bijection $\psi : \{0, \dots, n-1\} \rightarrow V_1$, $i \mapsto v_i$ and a sequence of n field labels $\vec{\varphi} = (\varphi_0, \dots, \varphi_{n-1})$, $\varphi_i \in \{A, c, \bar{c}, \gamma, \omega, \beta\}$ and $n = |V_1|$. The field label φ_i defines the type of the vertex $v_i \in V_1$. Let $V_\varphi \subseteq V_1$ be the set of all vertices of type φ , for example V_A . Furthermore, let $E_\varphi \subseteq E_1$ be the set of all edges incident to vertices in V_φ , for example E_\varkappa with $\varkappa \in \{\gamma, \omega\}$.



When needed, the edges are labeled by integers and the vertices by symbols. The edge incident to a vertex $v_i \in V_1$ has the same index i . As an example, for the tree above we have:

$$\begin{aligned}
V &= \{c_0, A_1, A_2, \bar{c}_3, u, u'\}, & V_3 &= \{u', u\}, \\
V_1 &= \{c_0, A_1, A_2, \bar{c}_3\}, & V_A &= \{A_1, A_2\}, & V_{\bar{c}} &= \{\bar{c}_3\}, & V_c &= \{c_0\}, \\
E_1 &= \{0, 1, 2, 3\}, & E_A &= \{1, 2\}, & E_{\bar{c}} &= \{3\}, & E_c &= \{0\}.
\end{aligned}$$

- Recalling the definition of \mathbb{P}_n in (16), for every edge $e \in E$ and vertex $v \in V_1$, the momentum assignments p_e, p_v are functions from the set \mathbb{P}_n to \mathbb{R}^4 , with $n = |V_1|$, defined by the following construction:

- label the vertices in V_1 by means of $\psi : i \mapsto v_i$ and set $p_{v_i}(\vec{p}) := p_i$,
- apply momentum conservation to all vertices to get $p_e(\vec{p})$.

We use similar notations for multi-indices: $w_v := w_{\psi^{-1}(v)}$ for $w \in \mathbb{W}_n$ and $v \in V_1$. Given the momentum assignments, a set-valued function K on E is defined by

$$K_e := \{v \in V_1 \setminus \{v_0\} : \frac{\partial p_e}{\partial p_v} \neq 0\}. \quad (134)$$

For the tree given above we have $p_{c_0} = -(p_{A_1} + p_{A_2} + p_{\bar{c}_3})$ and

$$K_0 = \{A_1, A_2, \bar{c}_3\}, \quad K_1 = \{A_1\}, \quad K_2 = \{A_2\}, \quad (135)$$

$$K_3 = \{\bar{c}_3\}, \quad K_4 = \{A_1, A_2\}. \quad (136)$$

Some additional structure is needed, always in view of the bounds.

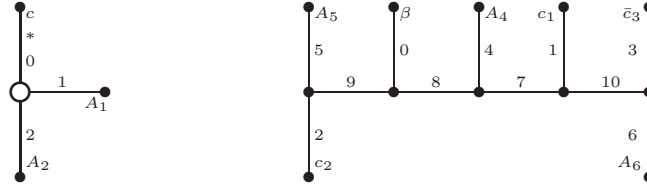
- The vertices in V_3 are additionally labeled either as ”**regular**” (\bullet) or as ”**hollow**” (\circ). The sets of regular and hollow vertices are respectively denoted by V_\bullet and V_\circ , hence $V_3 = V_\bullet \cup V_\circ$. In terms of our bounds, regular vertices do change the ρ -weight of incident edges, while hollow vertices do not, see definition 13. We use hollow vertices for the bounds on 3-point functions and in the proof of the theorems, see for example section 4.2 on the junction of weighted trees and definition 23.
- Edges carry zero or more labels ”*”. Edges are referred to as ”***-edges**” if they have one or more labels ”*”, and as ”**regular edges**” otherwise. The set of all *-edges is denoted by E_* . The *-edges play a special role in our bounds, because to each $e \in E_*$ is associated a supplementary factor $|p_e| + \Lambda$, see (152) and theorem 4.

Definition 11 Let be given a sequence of $n \geq 3$ field labels, $\vec{\varphi} = (\varphi_0, \dots, \varphi_{n-1})$, with $\varphi_i \in \{A, c, \bar{c}, \gamma, \omega, \beta\}$. Let $\mathbb{T}_{\vec{\varphi}}$ denote the set of all trees that satisfy the following rules:

- There is a bijection $\psi : \{0, \dots, n-1\} \rightarrow V_1$. Each $v_i \in V_1$ has type φ_i .
- $V = V_1 \cup V_3$.
- If $n = 3$ then $V_3 = V_0$.
- $|E_*| \in \mathbb{N}$.

Definition 12 In the notations of definition 11, let $\mathcal{T}_{\vec{\varphi}}^{(s)}$ denote the set of all trees in $\mathbb{T}_{\vec{\varphi}}$ with total number of labels "*" equal to s and such that $V_3 = V_0$ whenever $n > 3$. For shortness we set $\mathcal{T}_{\vec{\varphi}} := \mathcal{T}_{\vec{\varphi}}^{(0)}$ and $\mathcal{T}_{1\vec{\varphi}} := \mathcal{T}_{\vec{\varphi}}^{(1)}$.

As an example, below we show two trees, $\tau_3 \in \mathcal{T}_{1cAA}$ and $\tau_7 \in \mathcal{T}_{\beta c c \bar{c} A A A}$.



Definition 13 Fix a tree from $\mathbb{T}_{\vec{\varphi}}$. A ρ -weight is a function $\rho : E \rightarrow \{0, 1, 2\}$ with the following properties:

1. $\forall e \in E_1, \rho(e) = 0$.
2. There exists a map $\chi : V_0 \rightarrow E \setminus E_1$ such that
 - a) if $\chi(v) = e$, then e is incident to v ,
 - b) $\forall e \in E \setminus E_1, \rho(e) = 2 - |\chi^{-1}(\{e\})|$.

Definition 14 Let be given a tree from $\mathbb{T}_{\vec{\varphi}}$ and $w \in \mathbb{W}_n$, with $n = |V_1|$. A σ -weight is a function $\sigma : E \rightarrow \mathbb{N}$ defined by

$$\sigma(e) := \sum_{v \in V_1} \sigma_v(e), \quad (137)$$

where $(\sigma_v : E \rightarrow \mathbb{N})_{v \in V_1}$ is a family of functions such that

$$\sum_{e \in E} \sigma_v(e) = w_v, \quad \sigma_v(e) = 0 \text{ if } v \notin K_e. \quad (138)$$

By definition (21), $w_0 = 0$ for every $w \in \mathbb{W}_n$. Hence, $\sigma_{v_0}(e) = 0 \forall e \in E$.

Definition 15 Let be given a tree $\tau \in \mathbb{T}_{\vec{\varphi}}$ and $w \in \mathbb{W}_n$, with $n = |V_1|$. A θ -weight is a function $\theta : E \rightarrow \mathbb{N}$ defined by

$$\theta(e) := \rho(e) + \sigma(e), \quad (139)$$

where ρ and σ are a ρ -weight and a σ -weight corresponding to w , respectively. The pair (τ, θ) is a weighted tree. The total θ -weight of (τ, θ) is

$$\theta(\tau) := \sum_{e \in E} \theta(e). \quad (140)$$

The set of all θ -weights corresponding to given τ and w is denoted by Θ_{τ}^w .

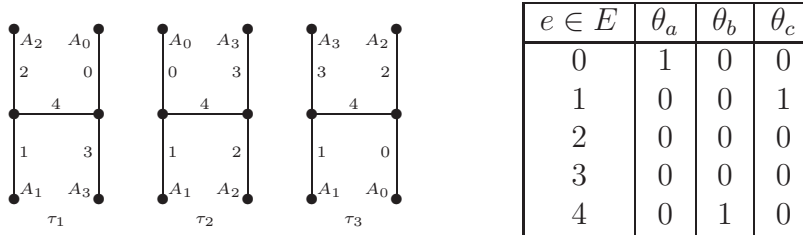
For every tree $\tau \in \mathcal{T}_{\vec{\varphi}}^{(s)}$ with $n \geq 4$ the total θ -weight is given by the formula

$$\theta(\tau) = n + \|w\| - 4. \quad (141)$$

This relation follows from definitions 12, 13, 14, which give the sum rule $\sum_{e \in E} \theta(e) = \|w\| + 2|E \setminus E_1| - |V_3|$, and from the relations $|E \setminus E_1| - |V_3| + 1 = 0$ and $|V_3| = n - 2$.

As an example we consider three trees $\tau_1, \tau_2, \tau_3 \in \mathcal{T}_{AAAA}$. We give three different weights $\theta_a, \theta_b, \theta_c$, which all correspond to the derivative wrt the momentum p_1 , literally $w_1 = 1$ and $w = (0, 1, 0, 0)$. We find a family of weighted trees $\{(\tau_i, \theta) : \theta \in \Theta_{\tau_i}^w\}_{i \in \{1, 2, 3\}}$, where

$$\Theta_{\tau_3}^w = \Theta_{\tau_1}^w = \{\theta_a, \theta_b, \theta_c\}, \quad \Theta_{\tau_2}^w = \{\theta_a, \theta_c\}. \quad (142)$$



3.2 Theorems

We always assume that the renormalization constants are independent of Λ_0 (though weakly Λ_0 -dependent "renormalization constants" can also be accommodated for, see [KKS92]). From now on we denote $\vec{\varkappa} := (\varkappa_1, \dots, \varkappa_{n_{\varkappa}})$ with $\varkappa_i \in \{\gamma, \omega\}$, $n_{\varkappa} \geq 0$.

Hypothesis RC1 We impose on all strictly relevant terms vanishing renormalization conditions at zero momentum and $\Lambda = 0$:

$$\Gamma_{\vec{z}}^{0\Lambda_0; \vec{\phi}; w}(0) = 0, \quad \text{if } 2n_{\varkappa} + N + \|w\| < 4. \quad (143)$$

Hypothesis RC2 On the following marginal terms we impose renormalization conditions at zero momentum and $\Lambda = M$:

$$\Gamma^{M\Lambda_0; c\bar{c}\bar{c}}(0) = 0, \quad \Gamma^{M\Lambda_0; c\bar{c}AA}(0) = 0, \quad \partial_A \Gamma^{M\Lambda_0; c\bar{c}A}(0) = 0, \quad (144)$$

for the notation see (23).

Remark that Bose–Fermi symmetry and translation invariance imply that $\partial_A \Gamma^{M\Lambda_0; c\bar{c}A}(0) = 0$ iff $\partial_c \Gamma^{M\Lambda_0; A\bar{c}c}(0) = 0$. To prove proposition 16 and theorem 1 all remaining marginal renormalization constants are chosen at $\Lambda = 0$ arbitrarily but in agreement with the global symmetries of the regularized theory: $SU(2)$, Euclidean isometries $ISO(4)$, ghost number conservation. For instance, all renormalization conditions must comply with the vanishing of the ghost number violating functions, like $\Gamma^{\Lambda, \Lambda_0; cccc}$ or $\Gamma_{(\varkappa_0, \varkappa_1)}^{\Lambda, \Lambda_0}$ for $\varkappa_i \in \{\gamma, \omega\}$. The list of the remaining marginal renormalization constants follows literally from (511) and appendix F.

Proposition 16 Assume the validity of hypotheses RC1 and RC2. For all sequences of $N \geq 3$ field labels in $\{A, c, \bar{c}\}$ with $\phi_{N-1} = \bar{c}$, denoted by $\vec{\phi}\bar{c}$, all $w = (w', 0) \in \mathbb{W}_N$, all $(\vec{p}, 0) \in \mathbb{P}_N$, and all positive Λ, Λ_0 s.t. $\max(\Lambda, M) \leq \Lambda_0$,

$$\Gamma^{\Lambda\Lambda_0; \vec{\phi}\bar{c}; w}(\vec{p}, 0) = 0. \quad (145)$$

Note that in (145) the momentum of the indicated antighost \bar{c} vanishes, and there is no derivative wrt this momentum.

Proof We prove the statement by induction, increasing in the loop order, $l - 1 \mapsto l$. Given l , we proceed by descending from w_{max} in the number of derivatives, $\|w\| \mapsto \|w\| - 1$. For fixed l and w , all possible terms $\Gamma_l^{\vec{\phi}\bar{c}; w}$ are considered. By construction, for fixed l and $\vec{\phi}\bar{c}$, the inductive scheme deals first with the irrelevant terms and continues, if they exist, with the marginal terms, followed by more and more relevant terms. Since the momentum of the antighost has been assumed to vanish, the statement holds at loop order $l = 0$. The validity of the statement for all loop orders smaller than l implies

that $\dot{\Gamma}_l^{\Lambda\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0$. The irrelevant terms have vanishing boundary conditions, hence $\Gamma_l^{\Lambda_0\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0$. Integrating the FE from Λ_0 downwards to arbitrary $\Lambda > 0$, we get $\Gamma_l^{\Lambda\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0$ for the irrelevant terms. Next we consider the marginal terms. Since the corresponding irrelevant terms have already been shown to vanish at vanishing antighost momentum, we use the Taylor formula to extend (144) to arbitrary momenta $(\vec{p}, 0)$, still preserving the vanishing antighost momentum. Then, we integrate the FE from M to arbitrary $\Lambda > 0$, which completes the proof that $\Gamma_l^{\Lambda\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0$ for marginal terms. Similar arguments hold for all the strictly relevant terms $\Gamma_l^{\vec{\phi}\vec{c};w}$. ■

Corollary 17 *The following counterterms vanish, see (30),*

$$r_1^{0,\bar{c}cAA} = 0, \quad r^{0,\bar{c}c\bar{c}c} = 0, \quad r_2^{0,\bar{c}cAA} = 0, \quad r_2^{0,A\bar{c}c} = 0. \quad (146)$$

Proof Using (145) we have, for all $p_2, p_3 \in \mathbb{R}^4$ and $\Lambda = \Lambda_0$ (omitted),

$$\Gamma^{c\bar{c}c\bar{c}}(0, p_2, p_3) = 0, \quad \Gamma^{c\bar{c}AA}(0, p_2, p_3) = 0, \quad \partial_{p_2}\Gamma^{A\bar{c}c}(0, p_2) = 0. \quad (147)$$

Recall (86) and (30). ■

Corollary 18 *For all $X \in \{\beta, 1\}$, $\vec{\varkappa}$, $\vec{\phi}\vec{c}$, $w = (w', 0)$, \vec{p} , and all positive Λ , Λ_0 s.t. $\max(\Lambda, M) \leq \Lambda_0$:*

$$\Gamma_{X\vec{\varkappa}}^{\Lambda\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0. \quad (148)$$

Proof It follows from the definitions of the inserted functions given in (87), (88), (106) that at tree level $\Gamma_{X\vec{\varkappa};l=0}^{\Lambda\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0$. Then using equation (146) one shows that for all these terms we have vanishing boundary conditions, $\Gamma_{X\vec{\varkappa};l}^{\Lambda_0\Lambda_0;\vec{\phi}\vec{c};w}(\vec{p}, 0) = 0 \forall l$. Assuming that the statement is true at the loop order $l - 1 \geq 0$, by induction in l , using (145) and integrating the FE from Λ_0 to arbitrary Λ one shows that it holds at the loop order l . ■

From now on, for simplicity of notation we write $\mathcal{P}_s^{(k)}$ to denote polynomials with nonnegative coefficients and degree s where the superscript is a

label to make one polynomial different from another. We define:

$$P_s^{\lambda_1 \lambda_2}(\vec{p}) := \mathcal{P}_s^{(0)}\left(\log_+ \frac{\max(|\vec{p}|, M)}{\lambda_1 + \eta(\vec{p})}\right) + \mathcal{P}_s^{(1)}\left(\log_+ \frac{\lambda_2}{M}\right), \quad (149)$$

$$P_s^\lambda(\vec{p}) := P_s^{\lambda \lambda}(\vec{p}), \quad (150)$$

$$\Pi_{\tau, \theta}^\lambda(\vec{p}) := \prod_{e \in E} (\lambda + |p_e|)^{-\theta(e)}, \quad (151)$$

$$Q_\tau^{\Lambda; w}(\vec{p}) := \frac{\prod_{e \in E_*} (\Lambda + |p_e|)}{\prod_{e \in E_\varkappa} (\Lambda + |p_e|)} \begin{cases} \inf_{i \in \mathbb{I}} \sum_{\theta \in \Theta_\tau^{w'(i)}} \Pi_{\tau, \theta}^\Lambda(\vec{p}), & |V_1| = 3, \\ \sum_{\theta \in \Theta_\tau^w} \Pi_{\tau, \theta}^\Lambda(\vec{p}), & \text{otherwise,} \end{cases} \quad (152)$$

where $\tau \in \mathbb{T}_{\vec{\phi}}$, $w'(i)$ is obtained from w by diminishing w_i by one unit, and, for nonvanishing w , $\mathbb{I} := \{i : w_i > 0\}$. The following sets are also used in theorems 1–4:

$$\mathbb{Y}_n^+ := \{(\Lambda, \Lambda_0) : 0 < \Lambda \leq \Lambda_0 \text{ and } \Lambda_0 \geq M\} \times \mathbb{P}_n, \quad (153)$$

$$\mathbb{Y}_n := \mathbb{Y}_n^+ \cup \{(0, \Lambda_0) : \Lambda_0 \geq M\} \times \{\vec{p} \in \mathbb{P}_n : \eta(\vec{p}) \neq 0\}. \quad (154)$$

Theorem 1 *There exists a collection of regular vertex functions $\Gamma_{\vec{z}; l}^{\vec{\phi}}$ on $\mathbb{Y}_{N+n_\varkappa}^+$, complying with the global symmetries of the theory, satisfying the FE and the renormalization conditions given by hypotheses RC1 and RC2, and with irrelevant terms vanishing at $\Lambda = \Lambda_0$. Furthermore, for all $\vec{\phi}$, \vec{z} , all $l \in \mathbb{N}$, $w \in \mathbb{W}_{N+n_\varkappa}$, the following bounds hold on $\mathbb{Y}_{N+n_\varkappa}^+$:*

a) $d \geq 0$ or $N + n_\varkappa = 2$

$$|\Gamma_{\vec{z}; l}^{\Lambda \Lambda_0; \vec{\phi}; w}(\vec{p})| \leq (\Lambda + |\vec{p}|)^d P_r^\Lambda(\vec{p}), \quad (155)$$

b) $d < 0$

$$|\Gamma_{\vec{z}; l}^{\Lambda \Lambda_0; \vec{\phi}; w}(\vec{p})| \leq \sum_{\tau \in \mathcal{T}_{\vec{z}\vec{\phi}}} Q_\tau^{\Lambda; w}(\vec{p}) P_r^\Lambda(\vec{p}). \quad (156)$$

Here $d := 4 - 2n_\varkappa - N - \|w\|$. If $l = 0$ then $r := 0$, otherwise r stands for $r(d, l)$.

$$r(d, l) := \begin{cases} 2l, & d \geq 0, \\ 2l - 1, & d < 0. \end{cases} \quad (157)$$

Theorem 1 shows in particular that the functions $\Gamma_{\vec{z}}^{\Lambda\Lambda_0;\vec{\phi};w}$ are bounded uniformly in Λ_0 . To prove convergence in the limit $\Lambda_0 \rightarrow \infty$ we establish the following bounds for their derivatives wrt Λ_0 .

Theorem 2 *Let be given a collection of vertex functions $\Gamma_{\vec{z};l}^{\vec{\phi}}$ as in theorem 1. Then, for all $\vec{\phi}$, \vec{z} , all $l \in \mathbb{N}$, $w \in \mathbb{W}_{\mathbb{N}+n_{\mathcal{X}}}$, the following bounds hold on $\mathbb{Y}_{\mathbb{N}+n_{\mathcal{X}}}^+$:*

a) $d \geq 0$ or $\mathbb{N} + n_{\mathcal{X}} = 2$

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \frac{\Lambda + M + |\vec{p}|}{\Lambda_0^2} (\Lambda + |p|)^d P_r^{\Lambda\Lambda_0}(\vec{p}), \quad (158)$$

b) $d < 0$

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \frac{\Lambda + M + |\vec{p}|}{\Lambda_0^2} \sum_{\tau \in \mathcal{T}_{\vec{z}\vec{\phi}}} Q_{\tau}^{\Lambda;w}(\vec{p}) P_r^{\Lambda\Lambda_0}(\vec{p}). \quad (159)$$

See theorem 1 for the definition of d and r .

Convergence of the limit $\Lambda \rightarrow 0^+$ of the terms $\Gamma_{\vec{z}}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})$, $\partial_{\Lambda_0} \Gamma_{\vec{z}}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})$ when \vec{p} is nonexceptional (or $d > 0$) follows from the Cauchy criterion

$$|\partial_{\Lambda_0}^k \Gamma_{\vec{z}}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p}) - \partial_{\Lambda_0}^k \Gamma_{\vec{z}}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \int_{\Lambda}^{\Lambda'} d\lambda |\partial_{\lambda} \partial_{\Lambda_0}^k \Gamma_{\vec{z}}^{\lambda\Lambda_0;\vec{\phi};w}(\vec{p})|, \quad (160)$$

and the bounds from theorems 1, 2. Convergence of the limit $\Lambda_0 \rightarrow \infty$ of the terms $\Gamma_{\vec{z}}^{0\Lambda_0;\vec{\phi};w}(\vec{p})$ when \vec{p} is nonexceptional (or $d > 0$) follows from Cauchy criterion and the bounds from theorem 2,

$$|\Gamma_{\vec{z}}^{0\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_{\vec{z}}^{0\Lambda'_0;\vec{\phi};w}(\vec{p})| \leq \int_{\Lambda_0}^{\Lambda'_0} d\lambda_0 |\partial_{\lambda_0} \Gamma_{\vec{z}}^{0\lambda_0;\vec{\phi};w}(\vec{p})|. \quad (161)$$

In the following we will consider the functions $\Gamma_{1\vec{z}}^{0\Lambda_0;\vec{\phi};w}$, $\Gamma_{\beta\vec{z}}^{0\Lambda_0;\vec{\phi};w}$ which appear on the lhs respectively of the ST identities (124) and of the AGE (123). The goal of theorems 3, 4 is to show that $\Gamma_{1\vec{z}}^{0\Lambda_0;\vec{\phi};w}$ and $\Gamma_{\beta\vec{z}}^{0\Lambda_0;\vec{\phi};w}$ vanish in the limit $\Lambda_0 \rightarrow \infty$, which restores the STI and AGE. The renormalization

conditions for these functions at $\Lambda = 0$ are obtained from the rhs of the STI and AGE. In section 5 we show that the required boundary marginal terms $\Gamma_{1\vec{z}}^{0\Lambda_0;\vec{\phi};w}$, $\Gamma_{\beta\vec{z}}^{0\Lambda_0;\vec{\phi};w}$ satisfy the bounds of the theorems under the conditions specified in hypothesis RC3.

Hypothesis RC3 *We allow R^{AAA} , r_1^{AA} , $r^{\bar{c}c}$ to be chosen arbitrarily but the remaining marginal renormalization constants must satisfy a set of equations: $R_1(308)$, $R_2(327)$, $R_3(330)$, $r_2^{AA}(324)$, $R_1^{A\bar{c}c}(314)$, $R_{1,2}^{AAAA}(334)$, see appendix F and (511) in appendix E for notations.*

For shortness we also introduce the following definition

$$F_s^{\Lambda\Lambda_0}(\vec{p}) := \frac{M + |\vec{p}| + \Lambda}{\Lambda_0} \left(1 + \left(\frac{|\vec{p}|}{\Lambda_0} \right)^{w_{max}} \right) \mathcal{P}_s^{(2)} \left(\frac{|\vec{p}|}{\Lambda + M} \right). \quad (162)$$

Theorem 3 *Let be given a collection of vertex functions $\Gamma_{\vec{z};l}^{\vec{\phi}}$, regular on $\mathbb{Y}_{N+n_\varkappa}$, complying with the global symmetries of the theory, satisfying the hypotheses RC1, RC2, RC3 and the bounds of theorems 1,2. Let $\Gamma_{\beta,\vec{z};l}^{\vec{\phi}}$ be a collection of vertex functions with one insertion of the operator $Q_\beta^{\Lambda_0}$ (101), regular on $\mathbb{Y}_{1+N+n_\varkappa}$, complying with the global symmetries of the theory, satisfying the FE, and s.t. the AGE (123) holds. Then, for all $\vec{\phi}$, \vec{z} , all $l \in \mathbb{N}$, $w \in \mathbb{W}_{1+N+n_\varkappa}$, the following bounds hold on $\mathbb{Y}_{1+N+n_\varkappa}$:*

a) $d \geq 0$ or $N + n_\varkappa = 1$

$$|\Gamma_{\beta\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq (\Lambda + |\vec{p}|)^d F_{s_\beta}^{\Lambda\Lambda_0}(\vec{p}) P_{r_\beta}^{\Lambda\Lambda_0}(\vec{p}), \quad (163)$$

b) $d < 0$

$$|\Gamma_{\beta\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \sum_{\tau \in \mathcal{T}_{\beta\vec{z}\vec{\phi}}} Q_\tau^{\Lambda;w}(\vec{p}) F_{s_\beta}^{\Lambda\Lambda_0}(\vec{p}) P_{r_\beta}^{\Lambda\Lambda_0}(\vec{p}). \quad (164)$$

Here $d := 3 - 2n_\varkappa - N - \|w\|$ and $s_\beta := 0$. If $l = 0$ then $r_\beta := 0$ otherwise r_β stands for $r_\beta(d, l)$.

$$r_\beta(d, l) := \begin{cases} 2l, & d \geq 0, \\ 2l - 1, & d < 0. \end{cases} \quad (165)$$

We note that the 1-point vertex function with integrated insertion (110) vanishes, $\Gamma_1^{\Lambda\Lambda_0;c;w} = 0$, e.g. due to SU(2) symmetry.

Theorem 4 *Let be given a collection of vertex functions $\Gamma_{\vec{x};l}^{\vec{\phi}}$, $\Gamma_{\beta\vec{x};l}^{\vec{\phi}}$ as in theorem 3. Let $\Gamma_{1,\vec{x};l}^{\vec{\phi}}$ be a collection of vertex functions with one integrated insertion of the appropriate operator among $Q_{\rho}^{\Lambda_0}$ (100), $Q_{\rho\gamma}^{\Lambda_0}$ (95), $Q_{\rho\omega}^{\Lambda_0}$ (96), regular on $\mathbb{Y}_{N+n_{\varkappa}}$, complying with the global symmetries of the theory, and satisfying the FE. Assume that the STI (123) and consistency conditions (121) and (122) do hold. Then, for all $\vec{\phi}$, \vec{x} , all $l \in \mathbb{N}$, $w \in \mathbb{W}_{N+n_{\varkappa}}$, the following bounds hold on $\mathbb{Y}_{N+n_{\varkappa}}$:*

a) $d > 0$ or $N + n_{\varkappa} = 2$

$$|\Gamma_{1\vec{x};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq (\Lambda + |\vec{p}|)^d F_{s_1}^{\Lambda\Lambda_0}(\vec{p}) P_{r_1}^{\Lambda\Lambda_0}(\vec{p}), \quad (166)$$

b) $d \leq 0$

$$|\Gamma_{1\vec{x};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \sum_{\tau \in \mathcal{T}_{1\vec{x}\vec{\phi}}} Q_{\tau}^{\Lambda;w}(\vec{p}) F_{s_1}^{\Lambda\Lambda_0}(\vec{p}) P_{r_1}^{\Lambda\Lambda_0}(\vec{p}). \quad (167)$$

Here $d := 5 - 2n_{\varkappa} - N - \|w\|$. If $l = 0$ then $r_1 := 0$, $s_1 := 0$ otherwise r_1, s_1 stand respectively for $r_1(d, l)$, $s_1(d, l)$.

$$r_1(d, l) := \begin{cases} 3l, & d > 0, \\ 3l - 1, & d = 0, \\ 3l - 2, & d < 0, \end{cases} \quad s_1(d, l) := \begin{cases} l, & d \geq 0, \\ l - 1, & d < 0. \end{cases} \quad (168)$$

4 Proof of Theorems 1–4

In this section we will prove theorems 1, 2, 3 and 4 in this order. We proceed by induction in the loop order l . We first verify that they hold at tree level $l = 0$. Afterwards we assume that they hold true up to loop order $l - 1 \geq 0$, and we will verify the induction step from $l - 1$ to l .

Put $D_X := 4$ for all vertex functions $\Gamma_{\vec{x}}^{\vec{\phi}}$. For all inserted functions $\Gamma_{X\vec{x}}^{\vec{\phi};w}$ with $X \in \{\beta, 1\}$ and $\|w\| \leq w_{max}$ let

$$D_X := \begin{cases} 3, & X = \beta, \\ 5, & X = 1, \end{cases} \quad d_X := D_X - 2n_{\varkappa} - N - \|w\|. \quad (169)$$

Note that at zero loop order $\Gamma_{l=0}^{\Lambda\Lambda_0} = \Gamma_{l=0}^{\Lambda_0\Lambda_0}$. Using (86), (110) and the definition of \mathcal{L}_{aux} in (106) one finds that in momentum space

$$\Gamma_1^{\Lambda_0\Lambda_0;\vec{\phi}} = Q_{\rho(0)}^{\Lambda_0;\vec{\phi}}, \quad \Gamma_{1\varkappa}^{\Lambda_0\Lambda_0;\vec{\phi}} = Q_{\rho(0)\varkappa}^{\Lambda_0;\vec{\phi}}, \quad \Gamma_{\beta}^{\Lambda_0\Lambda_0;\vec{\phi}} = Q_{\beta}^{\Lambda_0;\vec{\phi}}, \quad (170)$$

where the momentum variable corresponding to the source ρ is set to zero. Everywhere in the following $Q_\rho^{\Lambda_0}$ will stand for $Q_{\rho(0)}^{\Lambda_0}$. From definition of $Q_\rho^{\Lambda_0}$ in (100) it follows that the vertex functions $Q_{\rho;l=0}^{\Lambda_0;\vec{\phi}}$ with $N = 2$ vanish. The nonvanishing functions $Q_{\rho;l=0}^{\Lambda_0;\vec{\phi}}$, $Q_{\beta;l=0}^{\Lambda_0;\vec{\phi}}$ have the form $h_s(p, q)(1 - \sigma_{0\Lambda_0}(p^2))$ where h_s is a homogeneous tensor polynomial of degree $s \leq 2$ in the momentum variables $p, q \in \mathbb{R}^4$ which depends at most linearly on the momentum q . From the definitions of $Q_{\rho\omega}^{\Lambda_0}$ (96) and $Q_{\rho\gamma}^{\Lambda_0}$ (95) we obtain that $Q_{\rho\omega;l=0}^{\Lambda_0;\vec{\phi}}$ has the form $h_s(p)(\sigma_{0\Lambda_0}((p+q)^2) - \sigma_{0\Lambda_0}(p^2))$ with $s \leq 1$. For $\|w\| \leq s$ (relevant terms) using inequalities (429), (435), (440) we have

$$|\Gamma_{X\vec{z};l=0}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq c \Lambda_0^{s-\|w\|} \left(\frac{|\vec{p}|}{\Lambda_0} \right)^{s+1-\|w\|} \leq c \frac{|\vec{p}|}{\Lambda_0} (|\vec{p}| + \Lambda)^{s-\|w\|}. \quad (171)$$

For $\|w\| > s$ (irrelevant terms) the same inequalities yield

$$\begin{aligned} |\Gamma_{X\vec{z};l=0}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| &\leq c \Lambda_0^{s-\|w\|} \left(\frac{|\vec{p}|}{\Lambda_0} + 1 \right) \\ &\leq c (|\vec{p}| + \Lambda)^{s-\|w\|} \frac{|\vec{p}| + \Lambda}{\Lambda_0} \left(\frac{|\vec{p}|}{\Lambda_0} + 1 \right)^{\|w\|-s}. \end{aligned} \quad (172)$$

Since $s - \|w\|$ is the dimension d_X we have the following bounds

$$|\Gamma_{X\vec{z};l=0}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p})| \leq (\Lambda + |\vec{p}|)^{d_X} F_0^{\Lambda\Lambda_0}(\vec{p}), \quad X \in \{\beta, 1\}. \quad (173)$$

Thus the statements of theorems 1–4 hold at loop number $l = 0$. The proof proceeds by induction on l and on the number of derivatives $\|w\|$, ascending in l and, for fixed l , descending in $\|w\|$ from w_{max} to 0.

4.1 Chains of vertex functions

Definition 19 *A division in m parts of a finite set \mathbb{I} is a sequence $S := (s_j)_{j \in [m]}$ of m disjoint sets $s_j \subseteq \mathbb{I}$, possibly empty, such that $\bigcup_{j \in [m]} s_j = \mathbb{I}$. An ordered partition is a division with all s_j nonempty. Given a division S as above stated, a division of a sequence $\vec{\Psi} = (\Psi_i)_{i \in \mathbb{I}}$ is the sequence of elements $\vec{\Psi}_j := (\Psi_i)_{i \in s_j}$, with $j \in [m]$.*

Definition 20 *Let S be a division in m parts of a finite set \mathbb{I} , and $\vec{\Psi} = (\Psi_i)_{i \in \mathbb{I}}$ be a sequence of labels $\Psi_i \in \{A, c, \bar{c}, \gamma, \omega, \beta, 1\}$. Denote by $(\vec{\Psi}_j)$ the*

division of $\vec{\Psi}$ induced by S . A **chain** of vertex functions is then defined by the expression

$$\mathcal{F}_S^{\zeta_1 \vec{\Psi} \bar{\zeta}_m} := \Gamma_{\zeta_1 \bar{\Psi}_1 \bar{\zeta}_1} \prod_{j=2}^m \mathbf{C}_{\zeta_j \bar{\zeta}_{j-1}} \Gamma_{\zeta_j \bar{\Psi}_j \bar{\zeta}_j}, \quad (174)$$

where the repeated field labels $\zeta_j, \bar{\zeta}_j$ belong to $\{A, c, \bar{c}\}$ and are summed over (as usual).

Using this definition the FE (70) has the form

$$\dot{\Gamma}^{\vec{\Psi}} = \frac{\hbar}{2} \sum_S (-)^{\pi_a} \langle \dot{\mathcal{C}} \mathcal{F}_S^{A \vec{\Psi} A} + \dot{S} (\mathcal{F}_S^{c \vec{\Psi} c} - \mathcal{F}_S^{c \vec{\Psi} \bar{c}}) \rangle. \quad (175)$$

The sum above runs over all possible divisions of $[0 : n - 1]$, n being the number of components of $\vec{\Psi}$. The symbol π_a denotes the number of transpositions *mod* 2 of the anticommuting variables $\{c, \bar{c}, \beta, \gamma, 1\}$ in the permutation $i \mapsto \pi(i)$ such that $(\Psi_{\pi(0)}, \dots, \Psi_{\pi(n-1)}) = \vec{\Psi}_1 \oplus \dots \oplus \vec{\Psi}_m$, where $(a_1, \dots, a_p) \oplus (a_{p+1}, \dots, a_q) = (a_1, \dots, a_q)$.

A preliminary step toward the proof of theorem 1 is to bound $\partial^w (\mathbf{C}_{\zeta \bar{\zeta}} \Gamma^{\zeta \bar{\zeta}})$ with $\|w\| \leq w_{max}$.

Proposition 21 For all $0 < k < l$, $0 \leq w \leq w_{max}$, $p \in \mathbb{R}^4$

$$\left| \left(\prod_{i=0}^w \frac{\partial}{\partial p_{\mu_i}} \right) \left(\Gamma_k^{\zeta \bar{\zeta}; \Lambda \Lambda_0}(p) \mathbf{C}_{\zeta \bar{\zeta}}^{\Lambda \Lambda_0}(p) \right) \right| \leq \frac{P_{2k}^\Lambda}{(|p| + \Lambda)^w}. \quad (176)$$

Proof Using inequality (4) we see that

$$|\partial^w (\Gamma_k^{\zeta \bar{\zeta}; \Lambda \Lambda_0} \mathbf{C}_{\zeta \bar{\zeta}}^{\Lambda \Lambda_0})| \leq \sum_{w_1=0}^w \frac{w!}{w_1!(w-w_1)!} |\partial^{w_1} \Gamma_k^{\zeta \bar{\zeta}; \Lambda \Lambda_0}| |\partial^{w-w_1} \mathbf{C}_{\zeta \bar{\zeta}}^{\Lambda \Lambda_0}|. \quad (177)$$

Setting $w_2 = w - w_1$ it follows from (482) and the bounds of theorem 1 already proved inductively for $k < l$ that

$$\frac{1}{w_2!} |\partial^{w_1} \Gamma_k^{\zeta \bar{\zeta}; \Lambda \Lambda_0}| |\partial^{w_2} \mathbf{C}_{\zeta \bar{\zeta}}^{\Lambda \Lambda_0}| \leq \frac{c_\xi d^{w_2} P_{2k}^\Lambda}{(|p| + \Lambda)^{w_1+w_2}}. \quad (178)$$

Because $w \leq w_{max}$ the constants c_ξ, d^{w_2} may be absorbed in P_{2k}^Λ . ■

Definition 22 *Let be given*

1. a sequence $\vec{\Psi} = (\Psi_i)_{i \in \mathbb{I}}$ as in Definition 20;
2. an ordered partition $S = (s_j)_{j \in [m]}$ of \mathbb{I} ;
3. the sequences of field labels $(\zeta_j)_{j \in [m]}$, $(\bar{\zeta}_j)_{j \in [m]}$;
4. a multi-index $w \in \mathbb{W}_n$ and a sequence $\mathbf{w} := (\mathbf{w}_j)_{j \in [m]}$ such that $\mathbf{w}_j \in \mathbb{W}_n$ and $\sum_{j \in [m]} \mathbf{w}_j = w$.

Then, we define a **reduced chain** of vertex functions as

$$\Gamma^{\zeta_1 \vec{\Psi}_1 \bar{\zeta}_1; \mathbf{w}_1} \prod_{j=2}^m \mathbf{C}_{\zeta_j \bar{\zeta}_{j-1}} \Gamma^{\zeta_j \vec{\Psi}_j \bar{\zeta}_j; \mathbf{w}_j}, \quad (179)$$

where $\Gamma^{\zeta_j \vec{\Psi}_j \bar{\zeta}_j; \mathbf{w}_j} = \partial^{\mathbf{w}_j} (\Gamma^{\zeta_j \vec{\Psi}_j \bar{\zeta}_j})$ are derivatives wrt the external momenta appearing in $\dot{\Gamma}^{\vec{\Psi}}$, and the sequences $(\zeta_j)_{j \in [m]}$, $(\bar{\zeta}_j)_{j \in [m]}$ are **fixed**. Introducing the auxiliary quantities $\dot{\Upsilon} := (\dot{\Upsilon}_j)_{j \in [m]}$ and $\dot{\Upsilon}_j := \zeta_j \vec{\Psi}_j \bar{\zeta}_j$, we denote (179) by $\mathcal{S}^{\dot{\Upsilon}; \mathbf{w}}$ or, with some abuse of notations, by $\mathcal{S}_S^{\zeta_1 \vec{\Psi}_1 \bar{\zeta}_1; \mathbf{w}}$.

The adjective "reduced" indicates that the chains contain neither Γ^{AA} , $\Gamma^{c\bar{c}}$ nor derivatives applied to the propagators \mathbf{C} .

It follows from inequalities (482), (176) and theorems 1–4 proved in loop order $l - 1$ that there exists a common bound for the terms $\Gamma_{l_1+l_2}^{\zeta \vec{\Psi} \bar{\zeta}; w_1+w_2+w_3} \mathbf{C}$ and $\Gamma_{l_1}^{\zeta \vec{\Psi} \bar{\zeta}; w_1} \partial^{w_2} (\mathbf{C} \Gamma_{l_2}^{\zeta \bar{\zeta}}) \partial^{w_3} \mathbf{C}$. This property basically follows from an explicit check of the degree of the polynomials in the bounds and from the fact that, for w_2, w_3 supported on p , a factor $(|p| + \Lambda)^{-\|w_2+w_3\|}$ times a weighted tree (τ, θ) in the bound of $\Gamma_{l_1}^{\zeta \vec{\Psi} \bar{\zeta}; w_1}$ may be bounded by the weighted tree (τ, θ') in the bound of $\Gamma_{l_1+l_2}^{\zeta \vec{\Psi} \bar{\zeta}; w_1+w_2+w_3}$, with θ' equal to θ everywhere but on the edge $e \in E_1$ carrying momentum p , for which $\theta'(e) = \theta(e) + \|w_2 + w_3\|$. Hence to bound $\dot{\Gamma}_l^{\vec{\Psi}}$ it is enough to consider a loop integral with a reduced chain

$$\int \frac{d^4 k}{(2\pi)^4} \dot{\mathbf{C}}_{\zeta \bar{\zeta}}(k) \mathcal{S}_{S; l-1}^{\zeta \vec{\Psi} \bar{\zeta}; \mathbf{w}}(k, \vec{p}_S, -k), \quad (180)$$

where $\vec{p}_{s_j} = (p_i)_{i \in s_j}$. As an example we give in Appendix B the complete list of chains for $\Gamma^{AAc\bar{c}}$. The appellative "reduced" may be omitted in the following, since it is always clear from the context whether a chain is reduced or not.

4.2 Junction of weighted trees

Given a reduced chain $\mathcal{S}^{\hat{\Upsilon};w}$ we define its amplitude $\hat{\mathcal{S}}^{\hat{\Upsilon};w}$ by substituting the vertex functions and propagators with their corresponding bounds taken from theorems 1–4 and inequalities (482). Recalling that $\vec{\Upsilon}_j := \zeta_j \vec{\Psi}_j \bar{\zeta}_j$, we then set

$$\hat{\mathcal{S}}^{\Lambda;\hat{\Upsilon};w} := \hat{\mathcal{S}}^{\Lambda;\vec{\Upsilon}_1;w_1} \prod_{j=2}^m \frac{1}{(\Lambda + |p_{\zeta_j}|)^2} \hat{\mathcal{S}}^{\Lambda;\vec{\Upsilon}_j;w_j} \quad (181)$$

$$\hat{\mathcal{S}}^{\Lambda;\vec{\Upsilon}_j;w} := \begin{cases} (\Lambda + |\vec{p}|)^{d_X}, & \text{case } a, \\ \sum_{\tau \in \mathcal{T}_{\vec{\Upsilon}_j}} Q_{\tau}^{\Lambda;w}, & \text{case } b, . \end{cases} \quad (182)$$

Here the cases a and b refer to the respective parts in theorems 1–4.

The tree structure of the bound is spoiled if there exists an interval $\mathbb{J}^a := [j^a : j^a + m^a - 1] \subset [m]$ such that all $\hat{\mathcal{S}}^{\Lambda;\vec{\Upsilon}_j;w}$ for $j \in \mathbb{J}^a$ correspond to a strictly relevant contribution, associated to the cases a in the theorems. A workaround for this difficulty will start with the following definition. For every tree $\tau \in \mathbb{T}_{\vec{\varphi}}$, set $E_{1;v} := \{e \in E_1 : e \text{ incident to } v\}$.

Definition 23 Let $\vec{\varphi}$ be an arbitrary sequence with $\varphi_i \in \{A, c, \bar{c}, \beta, \gamma, \omega\}$. A tree $f \in \mathbb{T}_{\vec{\varphi}}$ is a **fragment** if

- a) there exists $s \in \{0, 1\}$ s.t. $|V_{\circ}| + s$ equals the total number of "*" labels,
- b) $\forall v \in V_{\circ} : |E_{1;v}| \geq 2$ and $E_{1;v} \cap E_* \neq \emptyset$.

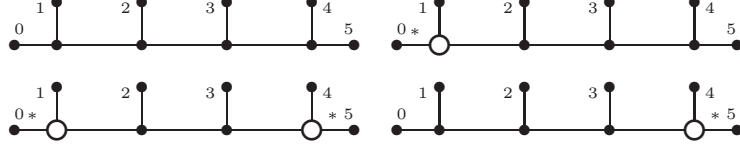
The set of all such fragments is denoted by $\mathcal{F}_{\vec{\varphi}}^{(s)}$. Moreover we set $\mathcal{F}_{\vec{\varphi}} := \mathcal{F}_{\vec{\varphi}}^{(0)}$ and $\mathcal{F}_{1\vec{\varphi}} := \mathcal{F}_{\vec{\varphi}}^{(1)}$.

For each $s \in \{0, 1\}$, $\mathcal{T}_{\vec{\varphi}}^{(s)} \subset \mathcal{F}_{\vec{\varphi}}^{(s)} \subset \mathbb{T}_{\vec{\varphi}}$.

Let us now state a useful property which relies on the above definition. Let be given a subsequence $\mathbb{J}^a := [j^a : j^a + m^a - 1] \subset [m]$ with any number of elements m^a . Let S and $\vec{\Psi}$ be as in definition 22. Define the following restrictions: $S^a = S|_{\mathbb{J}^a}$, $\mathbf{w}^a = \mathbf{w}|_{\mathbb{J}^a}$. Set $w^a := \sum_{j \in \mathbb{J}^a} w_j$ and $\vec{\Psi}^a := (\Psi_i)$ with $i \in \cup_{j \in \mathbb{J}^a} s_j$. Then, there exists a set of fragments $\mathcal{F}_{\zeta \vec{\Psi}^a \bar{\zeta}}$ such that

$$\hat{\mathcal{S}}_{S^a}^{\Lambda;\zeta \vec{\Psi}^a \bar{\zeta};w^a} \leq 2^{\frac{|S^a|}{2}} \sum_{f \in \mathcal{F}_{\zeta \vec{\Psi}^a \bar{\zeta}}} Q_f^{\Lambda;w^a}. \quad (183)$$

Here we give a proof of (183) for an example, generalisation is clear. Consider an amplitude $\hat{S}_{\zeta_1 A^4 \bar{\zeta}_4; w}$ composed of four elements $\hat{S}_{\zeta_j A \bar{\zeta}_j; w}$ where $j \in \{1, \dots, 4\}$ and $w = 0$. First let us define the following set of fragments \mathcal{F} (here $s = 0$):



In each fragment the $*$ -edge corresponds to a factor $\Lambda + |p_e|$ in the corresponding amplitude $Q_{f \in \mathcal{F}}^{\Lambda; w=0}$. One shows the following bound:

$$\begin{array}{c} \begin{array}{cccc} \begin{array}{c} 1 \\ \bullet \\ | \\ \oplus \\ 0 \end{array} & \begin{array}{c} 2 \\ \bullet \\ | \\ \oplus \\ |C| \end{array} & \begin{array}{c} 3 \\ \bullet \\ | \\ \oplus \\ |C| \end{array} & \begin{array}{c} 4 \\ \bullet \\ | \\ \oplus \\ |C| \end{array} \\ \hat{S}_{\zeta_1 A \bar{\zeta}_1} & \hat{S}_{\zeta_2 A \bar{\zeta}_2} & \hat{S}_{\zeta_3 A \bar{\zeta}_3} & \hat{S}_{\zeta_4 A \bar{\zeta}_4} \end{array} \end{array} \leq 2^{\frac{4}{2}} \sum_{f \in \mathcal{F}} Q_f^{\Lambda; w=0}, \quad (184)$$

or more explicitly:

$$\frac{\prod_{v \in V_3} (\Lambda + |\vec{p}_v|)}{\prod_{e \in E \setminus E_1} (\Lambda + |p_e|)^2} \leq 2^{\frac{|V_3|}{2}} \sum_{\{\chi\}} \frac{\prod_{v \in V_3} (\Lambda + |p_{\chi(v)}|)}{\prod_{e \in E \setminus E_1} (\Lambda + |p_e|)^2}. \quad (185)$$

Here the \oplus -vertices stand each for a corresponding $\hat{S}_{\zeta_j A \bar{\zeta}_j; w=0}$, see (182); the set of \oplus -vertices is identified with V_3 in (185); \vec{p}_v indicates the set of incoming momenta of the vertex v ; p_e denotes the momentum corresponding to an edge e ; the sum runs over the set of functions $\chi : V_3 \rightarrow E \setminus E_1 \cup \{e_0, e_5\}$ which map every vertex $v \in V_3$ to an edge incident to v ; the $|C|$'s stand for the usual bounds on the corresponding propagators.

Let be given a sequence of fragments $\vec{f} := (f_1, \dots, f_{m'})$ with $f_j \in \mathcal{F}_{\zeta_j \bar{\psi}_j \bar{\zeta}_j}$ and a sequence $w' = (w'_j)_{j \in [m']}$ with $w'_j \in \mathbb{W}_n$. We define the amplitude $\hat{Q}_{\vec{f}}^{\Lambda; w'}$ by

$$\hat{Q}_{\vec{f}}^{\Lambda; w'} := Q_{f_1}^{\Lambda; w'_1} \prod_{j=2}^{m'} \frac{1}{(\Lambda + |p_{\zeta_j}|)^2} Q_{f_j}^{\Lambda; w'_j}. \quad (186)$$

Lemma 24 *Given an amplitude $\hat{Q}_{\vec{f}}^{\Lambda; w'}$ as above there exists a fragment $f \in \mathcal{F}_{\zeta \bar{\psi} \bar{\zeta}}$ such that*

$$\hat{Q}_{\vec{f}}^{\Lambda; w'} \leq Q_f^{\Lambda; w'}. \quad (187)$$

Proof We proceed by induction in m' . If there are no joints, $m' = 1$, the statement is evident. Assume it is true for some $m' - 1 \geq 0$ and consider a sequence of m' fragments. Let $v_l, v_r \in V_3$ be the left and right vertices of a joint. Recall definition in (151).

- $v_l, v_r \in V_3 \setminus V_o$.

$$\begin{array}{c} | \quad v_l \quad \cdot \quad \text{---} \overset{C}{\text{---}} \quad \cdot \quad \text{---} \quad v_r \quad | \\ | \quad \underline{0} \quad \cdot \quad \text{---} \quad \cdot \quad \text{---} \quad | \end{array} \leq \begin{array}{c} | \quad v_l \quad \quad v_r \quad | \\ | \quad \underline{2} \quad \quad | \end{array} \quad (188)$$

$$\Pi_{f_l, \theta_l}^\Lambda \frac{1}{(\Lambda + |p|)^2} \Pi_{f_r, \theta_r}^\Lambda = \Pi_{f, \theta}^\Lambda \quad (189)$$

Here the corresponding external edges have been merged together to form a new internal edge with θ -weight equal to 2. The θ -weight of all other edges is unchanged.

- $v_l \in V_3 \setminus V_o, v_r \in V_o$ or vice versa.

$$\begin{array}{c} | \quad v_l \quad \cdot \quad \text{---} \overset{C}{\text{---}} \quad \cdot \quad \overset{*}{\text{---}} \quad v_r \quad | \\ | \quad \underline{0} \quad \cdot \quad \text{---} \quad \cdot \quad \text{---} \quad | \end{array} \leq \begin{array}{c} | \quad v_l \quad \quad v_r \quad | \\ | \quad \underline{1} \quad \quad | \end{array} \quad (190)$$

$$\Pi_{f_l, \theta_l}^\Lambda \frac{1}{(\Lambda + |p|)^2} (\Lambda + |p|) \Pi_{f_r, \theta_r}^\Lambda = \Pi_{f, \theta}^\Lambda \quad (191)$$

$$\begin{array}{c} | \quad v_l \quad \cdot \quad \text{---} \overset{C}{\text{---}} \quad \cdot \quad \overset{*}{\text{---}} \quad v_r \quad | \\ | \quad \underline{0} \quad \cdot \quad \text{---} \quad \cdot \quad \text{---} \quad | \end{array} \leq \begin{array}{c} | \quad v_l \quad \quad v_r \quad | \\ | \quad \underline{2} \quad \quad | \end{array} \quad (192)$$

$$\Pi_{f_l, \theta_l}^\Lambda \frac{1}{(\Lambda + |p|)^2} (\Lambda + |p_*|) \Pi_{f_r, \theta_r}^\Lambda = (\Lambda + |p_*|) \Pi_{f, \theta}^\Lambda \quad (193)$$

- $v_l, v_r \in V_o$.

$$\begin{array}{c} \cdot \quad | \quad v_l \quad \overset{*}{\text{---}} \quad \cdot \quad \text{---} \overset{C}{\text{---}} \quad \cdot \quad \overset{*}{\text{---}} \quad v_r \quad | \\ \cdot \quad | \quad \underline{0} \quad \cdot \quad \text{---} \quad \cdot \quad \text{---} \quad | \end{array} \leq \begin{array}{c} \cdot \quad | \quad v_l \quad \quad v_r \quad | \\ \cdot \quad | \quad \underline{1} \quad \quad | \end{array} \quad (194)$$

$$\Pi_{f_l, \theta_l}^\Lambda (\Lambda + |p|) \frac{1}{(\Lambda + |p|)^2} (\Lambda + |p_*|) \Pi_{f_r, \theta_r}^\Lambda = (\Lambda + |p_*|) \Pi_{f, \theta}^\Lambda \quad (195)$$

$$\begin{array}{c} \cdot \quad | \quad v_l \quad \overset{*}{\text{---}} \quad \cdot \quad \text{---} \overset{C}{\text{---}} \quad \cdot \quad \overset{*}{\text{---}} \quad v_r \quad | \\ \cdot \quad | \quad \underline{0} \quad \cdot \quad \text{---} \quad \cdot \quad \text{---} \quad | \end{array} \leq \begin{array}{c} \cdot \quad | \quad v_l \quad \quad v_r \quad | \\ \cdot \quad | \quad \underline{0} \quad \quad | \end{array} \quad (196)$$

$$\Pi_{f_l, \theta_l}^\Lambda (\Lambda + |p|) \frac{1}{(\Lambda + |p|)^2} (\Lambda + |p|) \Pi_{f_r, \theta_r}^\Lambda = \Pi_{f, \theta}^\Lambda \quad (197)$$

where $p_\varkappa + p_1 + p_2 = 0$, one realizes that for the marginal vertex function $\Gamma_{\varkappa;l'}^{\Lambda\Lambda_0;\phi_1\phi_2}$ the following inequality holds

$$|\Gamma_{\varkappa;l'}^{\Lambda\Lambda_0;\phi_1\phi_2}(\vec{p})| \leq \frac{1}{\Lambda + |p_\varkappa|} \left(\begin{array}{c} \bullet \\ | \\ \circ \\ | \\ \bullet \end{array} \begin{array}{c} \varkappa \\ | \\ \bullet \end{array} + \begin{array}{c} \bullet \\ | \\ \circ \\ | \\ \bullet \end{array} \begin{array}{c} \varkappa \\ | \\ \bullet \end{array} \right) P_{r_\varkappa(l')}^\Lambda. \quad (203)$$

Similarly, substituting the relevant terms $\Gamma_{\beta;l'}^{\Lambda\Lambda_0;\vec{\phi};w}$, $N + \|w\| \leq 3$, at loop number $l' < l$ with the bounds of theorem 3 we have

$$|\Gamma_{\beta;l'}^{\Lambda\Lambda_0;\phi_1\phi_2\phi_3}| \leq \begin{array}{c} \beta \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} F_0^{\Lambda\Lambda_0} P_{r_\beta(l')}^{\Lambda\Lambda_0}, \quad (204)$$

$$|\Gamma_{\beta;l'}^{\Lambda\Lambda_0;\phi_1\phi_2}| \leq \left(\begin{array}{c} \beta \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} + \begin{array}{c} \beta \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_2 \\ | \\ \bullet \\ | \\ \bullet \end{array} \right) F_0^{\Lambda\Lambda_0} P_{r_\beta(l')}^{\Lambda\Lambda_0}, \quad (205)$$

$$|\partial \Gamma_{\beta;l'}^{\Lambda\Lambda_0;\phi_1\phi_2}| \leq \begin{array}{c} \beta \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} F_0^{\Lambda\Lambda_0} P_{r_\beta(l')}^{\Lambda\Lambda_0}. \quad (206)$$

Furthermore using the bounds of theorem 4 which are assumed to be true for any loop order $l' < l$, we obtain the following inequalities for strictly relevant terms $\Gamma_{1;l'}^{\Lambda\Lambda_0;\vec{\phi};w}$, $N + \|w\| < 5$,

$$|\Gamma_{1;l'}^{\Lambda\Lambda_0;\phi_0\phi_1\phi_2\phi_3}| \leq \left(\begin{array}{c} \phi_0 \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} + \begin{array}{c} \phi_0 \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} + \dots \right) F_{s_1(l')}^{\Lambda\Lambda_0} P_{r_1(l')}^{\Lambda\Lambda_0}, \quad (207)$$

$$|\Gamma_{1;l'}^{\Lambda\Lambda_0;\phi_0\phi_1\phi_2}| \leq \left(\begin{array}{c} \phi_0 \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} + \begin{array}{c} \phi_0 \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_2 \\ | \\ \bullet \\ | \\ \bullet \end{array} + \dots \right) F_{s_1(l')}^{\Lambda\Lambda_0} P_{r_1(l')}^{\Lambda\Lambda_0}, \quad (208)$$

$$|\partial \Gamma_{1;l'}^{\Lambda\Lambda_0;\phi_0\phi_1\phi_2}| \leq \left(\begin{array}{c} \phi_0 \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_1 \\ | \\ \bullet \\ | \\ \bullet \end{array} + \begin{array}{c} \phi_0 \\ | \\ \bullet \\ | \\ \bullet \end{array} \begin{array}{c} \phi_2 \\ | \\ \bullet \\ | \\ \bullet \end{array} \right) F_{s_1(l')}^{\Lambda\Lambda_0} P_{r_1(l')}^{\Lambda\Lambda_0}, \quad (209)$$

$$|\Gamma_{1;l'}^{\Lambda\Lambda_0;\phi_0\phi_1;w} \mathbf{C}^{\Lambda\Lambda_0}| \leq \frac{\Lambda + |p|}{(\Lambda + |p|)\|w\|} F_{s_1(l')}^{\Lambda\Lambda_0} P_{r_1(l')}^{\Lambda\Lambda_0}, \quad (210)$$

where the dots stand for omitted fragments obtained by permuting the "*" over the external edges. Substituting the bounds for $\Gamma_{1;\varkappa}^{\vec{\phi};w}$ with those for $\partial_{\vec{c}} \Gamma_1^{\vec{\phi};w}$ yields a similar result for the relevant terms $\Gamma_{1;\varkappa}^{\vec{\phi};w}$ with $n_\varkappa > 0$. Consequently the amplitude $\hat{S}_{1;\vec{\varkappa}}^{\Lambda;\vec{\phi};w}$ is bounded by a sum of amplitudes $Q_f^{\Lambda;w}$ with $f \in \mathcal{F}_{1;\vec{\varphi}}$.

4.3 Irrelevant terms

The irrelevant terms at arbitrary $0 < \Lambda \leq \Lambda_0$ are reconstructed by using the FE:

$$\Gamma_{X\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p}) = \Gamma_{X\vec{z};l}^{\Lambda_0\Lambda_0;\vec{\phi};w}(\vec{p}) + \int_{\Lambda_0}^{\Lambda} d\lambda \dot{\Gamma}_{X\vec{z};l}^{\lambda\Lambda_0;\vec{\phi};w}(\vec{p}). \quad (211)$$

At $\Lambda = \Lambda_0$ for all loop orders $l > 0$, for all $n_x \geq 0$ and $d_X < 0$ we have

$$|\Gamma_{X\vec{z};l}^{\Lambda_0\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \begin{cases} (|\vec{p}| + \Lambda_0)^{d_X} F_{s_X}^{\Lambda_0\Lambda_0} P_{r_X}^{\Lambda_0\Lambda_0} & X \in \{\beta, 1\}, \\ 0 & otherwise, \end{cases} \quad (212)$$

where the upper inequality is obtained using the bounds on relevant terms from theorem 1.

To integrate the FE from the boundary Λ_0 to Λ we substitute the chains $\mathcal{S}^{\zeta\phi_0\phi_1\bar{\zeta}}$, $\mathcal{S}^{\zeta_x\phi\bar{\zeta}}$ with the bounds given in theorem 1 and use inequalities (381), (391). Eventually we get

$$|\Gamma_l^{\Lambda\Lambda_0;\phi_0\phi_1;w}(p)| \leq (\Lambda + |p|)^{2-\|w\|} P_{2l-1}^{\Lambda}, \quad \|w\| > 2, \quad (213)$$

$$|\Gamma_{z;l}^{\Lambda\Lambda_0;\phi;w}(p)| \leq (\Lambda + |p|)^{1-\|w\|} P_{2l-1}^{\Lambda}, \quad \|w\| > 1. \quad (214)$$

In a similar way substitution of $\mathcal{S}^{\zeta\beta\phi\bar{\zeta}}$, $\mathcal{S}_1^{\zeta\phi_0\phi_1\bar{\zeta}}$ with the bounds of theorems 3, 4 and then integration from Λ_0 to Λ , using (405), give

$$|\Gamma_{\beta;l}^{\Lambda\Lambda_0;\phi;w}(p) - \Gamma_{\beta;l}^{\Lambda_0\Lambda_0;\phi;w}(p)| < \frac{F_0^{\Lambda\Lambda_0} P_{2l-1}^{\Lambda\Lambda_0}}{(\Lambda + |p|)^{\|w\|-2}}, \quad \|w\| > 2, \quad (215)$$

$$|\Gamma_{1;l}^{\Lambda\Lambda_0;\phi_0\phi_1;w}(p) - \Gamma_{1;l}^{\Lambda_0\Lambda_0;\phi_0\phi_1;w}(p)| < \frac{F_{l-1}^{\Lambda\Lambda_0} P_{3l-2}^{\Lambda\Lambda_0}}{(\Lambda + |p|)^{\|w\|-3}}, \quad \|w\| > 3. \quad (216)$$

Proposition 26 *Let $\tau_i \in \mathcal{T}_{\zeta\bar{\zeta}\bar{\phi}}$, $\phi_j \in \{A, c, \bar{c}, \gamma, \omega\}$, $|V_1| > 4$, $k \in \mathbb{N}$ and $\theta(\tau_i) > 2$. Then $\exists \tau_f \in \mathcal{T}_{\bar{\phi}}$ such that*

$$\int_{\Lambda}^{\Lambda_0} d\lambda \lambda \mathcal{R}_{\zeta\bar{\zeta}}(Q_{\tau_i}^{\lambda;w}) P_k^{\lambda}|_{\eta=0} \leq Q_{\tau_f}^{\Lambda;w} P_{k+1}^{\Lambda}. \quad (217)$$

Proof Denote by $v, \bar{v} \in V \setminus V_1$ the vertices incident to $e_\zeta, e_{\bar{\zeta}} \in E_1$. First we assume that $v \neq \bar{v}$. The result of the restriction $\mathcal{R}_\zeta(Q_{\tau_i}^{\lambda;w})$ is the amplitude of the tree τ_i with two edges $e_1 = \{u_1, v\}$, $e_2 = \{u_2, v\}$ carrying opposite momenta. Here $\{u, v\}$ denotes the edge which links the vertices u, v . Without restriction we assume that $e_1 = \chi(v)$. Furthermore we define a subtree τ' of the initial tree τ_i by the substitution $e_2 \mapsto \{u_1, u_2\}$ such that it does not include the vertices v, ζ nor the edges e_1, e_ζ .

$$\rho(e_1) = 1.$$

The identity $p_{e_1} = -p_{e_2}$ implies

$$\lambda \frac{1}{(\lambda + |p_{e_1}|)^{\theta(e_1)}} \leq \frac{1}{(\lambda + |p_{e_1}|)^{\sigma(e_1)}} = \frac{1}{(\lambda + |p_{e_2}|)^{\sigma(e_1)}}. \quad (218)$$

Consequently, $\lambda \mathcal{R}_\zeta(Q_{\tau_i}^{\lambda;w})$ is bounded by the amplitude of the subtree τ' with $\sigma'(e_2) = \sigma(e_2) + \sigma(e_1)$.

$$\rho(e_1) = 0 \text{ and } \rho(e_2) > 0.$$

$$\lambda \frac{1}{(\lambda + |p_{e_2}|)^{\theta(e_2)}} \leq \frac{1}{(\lambda + |p_{e_2}|)^{\theta(e_2)-1}}. \quad (219)$$

Thus, $\lambda \mathcal{R}_\zeta(Q_{\tau_i}^{\lambda;w})$ is bounded by the amplitude of the subtree τ' with $\sigma'(e_2) = \sigma(e_2) + \sigma(e_1)$, $\rho'(e_2) = \rho(e_2) - 1$ and $\chi'(u_1) = e_2$.

$$\rho(e_1) = 0 \text{ and } \rho(e_2) = 0.$$

This implies $e_2 \in E_1$. Because $|V_1| > 4$ the vertex u_1 is incident to an edge $e_u \in E \setminus E_1$ such that $\rho(e_u) > 0$.

$$\lambda \frac{1}{(\lambda + |p_{e_u}|)^{\theta(e_u)}} \leq \frac{1}{(\lambda + |p_{e_u}|)^{\theta(e_u)-1}}. \quad (220)$$

It follows that $\lambda \mathcal{R}_\zeta Q_{\tau_i}^{\lambda;w}$ is bounded by the amplitude of the subtree τ' with $\sigma'(e_2) = \sigma(e_2) + \sigma(e_1)$, $\rho'(e_u) = \rho(e_u) - 1$ and $\chi'(u_1) = e_u$.

For the subtree τ' the total weight satisfies $\theta(\tau') > 1$. Consequently, we repeat the above reduction by substituting λ -multiplication with integration. Then using inequalities (381), (391) we have

$$\int_\Lambda^{\Lambda_0} \frac{d\lambda P_k^\lambda|_{\eta=0}}{(\lambda + |p_a|)(\lambda + |p_b|)} \leq \frac{P_{k+1}^\Lambda}{\Lambda + |p_a| + |p_b|} \leq \frac{P_{k+1}^\Lambda}{\Lambda + |\hat{p}|}, \quad (221)$$

where $\hat{p} \in \{p_a, p_b, p_a + p_b\}$. Here p_a denotes the momentum of the edge used for the reduction, and $\hat{p} = p_b$ is the momentum of an arbitrary edge with nonvanishing θ -weight of the final tree τ_f . Because of the previous edge reduction we can reach the threshold $|V'_1| = 4$. In this case we can choose freely \hat{p} to equal p_a , p_b or $p_a + p_b$ thus keeping the σ -weight on the corresponding edge, see (381) and (391).

The case $v = \bar{v}$ follows directly from (381), (391). To see this we denote by $e \in E \setminus E_1$ the edge incident to v and introduce a vertex $u \in V_3$ adjacent to v . Hence, $\chi(v) = e = \{v, u\}$. If $\tilde{\theta}, \theta$ are two θ -weights where the only difference is that $\chi(u) = e$ for the first and $\chi(u) \neq e$ for the second, then

$$\Pi_{\tau_i, \tilde{\theta}}^\lambda \leq \Pi_{\tau_i, \theta}^\lambda = \frac{1}{\lambda} \frac{1}{\lambda + p_{e \in E(\tau_f)}} \Pi_{\tau_f, \theta}^\lambda. \quad (222)$$

The final tree τ_f is a subtree of τ_i which does not include vertices $\zeta, \bar{\zeta}, v$ nor the corresponding edges. Moreover, $\theta(\tau_i) > 2 \implies \theta(\tau_f) > 0$. \blacksquare

Inequality (217) can be applied to bound $\mathcal{S}^{\Lambda\Lambda_0; \vec{\zeta}\vec{\phi}; w}$. In this case $\tau_i \in \mathcal{T}_{\zeta\bar{\zeta}\vec{\zeta}\vec{\phi}}$ and the total weight satisfies $\theta(\tau_i) > 2 - n_\varkappa$. But in theorem 1 for each edge $e_\varkappa \in E_\varkappa$ we have $(\Lambda + |p_\varkappa|)$ as a denominator which is equivalent to an additional σ -weight of the edge e_\varkappa . An effective tree $\tilde{\tau}$ with $\tilde{\sigma}(e_\varkappa) = \sigma(e_\varkappa) + 1$ has $\theta(\tilde{\tau}) > 2$ and satisfies the conditions of proposition 26.

Using (405) a similar inequality follows for the irrelevant functions $\Gamma_{\beta\vec{\zeta}; l}^{\Lambda\Lambda_0; \vec{\phi}; w}$. For any $\tau_i \in \mathcal{T}_{\zeta\bar{\zeta}\beta\vec{\zeta}\vec{\phi}}$ there exists $\tau_f \in \mathcal{T}_{\beta\vec{\zeta}\vec{\phi}}$ such that

$$\int_{\Lambda}^{\Lambda_0} d\lambda \lambda \mathcal{R}_{\zeta\bar{\zeta}}(Q_{\tau_i}^{\lambda; w}) F_0^{\lambda\Lambda_0} P_{2(l-1)}^{\lambda\Lambda_0} |_{\eta=0} \leq Q_{\tau_f}^{\Lambda; w} F_0^{\Lambda\Lambda_0} P_{2l-1}^{\Lambda\Lambda_0}. \quad (223)$$

Before application of (405) to the irrelevant terms $\Gamma_{1\vec{\zeta}; l}^{\Lambda\Lambda_0; \vec{\phi}; w}$ we need a minor change in (221)

$$\int_{\Lambda}^{\Lambda_0} d\lambda \left(\frac{\lambda + |p_*|}{\lambda + |p_e|} \right) \frac{F_s^{\lambda\Lambda_0} P_k^{\lambda\Lambda_0} |_{\eta=0}}{(\lambda + |p_a|)(\lambda + |p_b|)} \leq \left(\frac{\Lambda + |p'_*|}{\Lambda + |p_e|} \right) \frac{F_s^{\Lambda\Lambda_0} P_{k+1}^{\Lambda\Lambda_0}}{\Lambda + |p_b|}, \quad (224)$$

where p'_* is one out of p_* or p_e in such way that $|p'_*| := \max(|p_*|, |p_e|)$. If $p_* = p_e$ the label "*" is moved to edge e . Finally, $\forall \tau_i \in \mathcal{T}_{1\zeta\bar{\zeta}\vec{\zeta}\vec{\phi}}$

$$\int_{\Lambda}^{\Lambda_0} d\lambda \lambda \mathcal{R}_{\zeta\bar{\zeta}}(Q_{\tau_i}^{\lambda; w}) F_{l-1}^{\lambda\Lambda_0} P_{3(l-1)}^{\lambda\Lambda_0} |_{\eta=0} \leq \sum_{\tau \in \mathcal{T}_{1\vec{\zeta}\vec{\phi}}} Q_{\tau}^{\Lambda; w} F_{l-1}^{\Lambda\Lambda_0} P_{3l-2}^{\Lambda\Lambda_0}. \quad (225)$$

4.4 Marginal terms

Lemma 27 *Let $\Gamma_{\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}$ denote a marginal term, $n_{\varkappa} \leq 1$. Then $\forall \Lambda < \eta(\vec{p})$*

$$|\Gamma_{\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_{\vec{z};l}^{\eta\Lambda_0;\vec{\phi};w}(\vec{p})| \leq P_{2l-2}^\Lambda(\vec{p}). \quad (226)$$

Proof Note that in theorem 1 for all $l' < l$ the bounds for $\Gamma_{\varkappa;l'}^{\Lambda\Lambda_0;\vec{\phi};w}$ are more restrictive than the ones for $\partial_{\bar{c}}\Gamma_{l'}^{\Lambda\Lambda_0;\bar{c}\vec{\phi};w}$. So we will only treat the case $n_{\varkappa} = 0$ explicitly. In this case we integrate the FE from η to Λ

$$\begin{aligned} |\Gamma_l^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_l^{\eta\Lambda_0;\vec{\phi};w}(\vec{p})| &\leq \int_{\Lambda}^{\eta} d\lambda |\dot{\Gamma}_l^{\lambda\Lambda_0;\vec{\phi}}(\vec{p})| \\ &\leq \sum_{\tau \in \mathcal{T}_{\zeta\bar{\zeta}\vec{\phi}}\Lambda} \int_{\Lambda}^{\eta} d\lambda \lambda \mathcal{R}_{\zeta\bar{\zeta}}(Q_{\tau}^{\lambda;w}) P_{2l-2}^{\lambda}|_{\eta_{\tau}=0}, \end{aligned} \quad (227)$$

where $\theta(\tau) = 2$. Denoting by q_a, q_b the momenta of the edges with nonvanishing θ -weight and using equations (397),(399) for each term of the sum we obtain

$$\int_{\Lambda}^{\eta} d\lambda \frac{\lambda P_k^{\lambda}|_{\eta_{\tau}=0}}{(\lambda + |q_a|)(\lambda + |q_b|)} \leq \int_{\Lambda}^{\eta} d\lambda \frac{P_k^{\lambda}|_{\eta_{\tau}=0}}{\lambda + |q_b|} \leq P_k^{\Lambda}, \quad k = 2l - 2. \quad (228)$$

■

4.4.1 $\Gamma^{c\bar{c}AA}$ and $\Gamma^{c\bar{c}\bar{c}}$

The renormalization condition is $\Gamma_l^{M\Lambda_0;\vec{\phi}}(0) = 0$. For $p_2, p_3 \in \mathbb{R}^4$ and $\Lambda' := \max(\Lambda, \eta(\vec{p}))$ equation (145) gives $\Gamma_l^{\Lambda'\Lambda_0;\vec{\phi}}(0_{\bar{c}}, p_2, p_3) = 0$, where the subscript \bar{c} indicates the momentum of the antighost.

$$|\Gamma_l^{\Lambda'\Lambda_0;\vec{\phi}}(\vec{p})| \leq \int_0^1 dt |p_{\bar{c}}| |\partial_{\bar{c}}\Gamma_l^{\Lambda'\Lambda_0;\vec{\phi}}(tp_{\bar{c}}, p_2, p_3)|. \quad (229)$$

Substituting $|\partial_{\bar{c}}\Gamma_l^{\Lambda'\Lambda_0;\vec{\phi}}|$ with the bound of theorem 1 and using inequality (413) we obtain

$$|\Gamma_l^{\Lambda'\Lambda_0;\vec{\phi}}(\vec{p})| \leq \left(1 + \log_+ \frac{|\vec{p}|}{\Lambda' + \eta}\right) P_{2l-1}^{\Lambda'\Lambda'}(\vec{p}) \Big|_{\eta=0} \leq P_{2l}^{\Lambda}(\vec{p}). \quad (230)$$

If $\Lambda' = \Lambda$ the proof is finished. Otherwise we use (226).

4.4.2 $\partial_A \Gamma^{c\bar{c}A}$

The renormalization condition is $\partial_A \Gamma_l^{M\Lambda_0;c\bar{c}A}(0) = 0$. For $p_A \in \mathbb{R}^4$ and $\Lambda' = \max(\Lambda, \eta(\vec{p}))$ equation (145) gives $\partial_A \Gamma_l^{\Lambda'\Lambda_0;c\bar{c}A}(0, p_A) = 0$.

$$|\partial_A \Gamma_l^{\Lambda'\Lambda_0;c\bar{c}A}(\vec{p})| \leq \int_0^1 dt |p_{\bar{c}}| |\partial_{\bar{c}} \partial_A \Gamma_l^{\Lambda'\Lambda_0;c\bar{c}A}(tp_{\bar{c}}, p_A)|. \quad (231)$$

We substitute $|\partial_{\bar{c}} \partial_A \Gamma_l^{\Lambda'\Lambda_0;c\bar{c}A}(\vec{p}_t)|$ with the bound from theorem 1 with the choice $w'(2) = (0, 1, 0)$, see after (152) for the definition of w' . Then

$$Q_{c\bar{c}A}^{\Lambda';(0,1,1)} \in \left\{ \frac{1}{\Lambda' + t|p_{\bar{c}}|}, \frac{1}{\Lambda' + |tp_{\bar{c}} + p_A|} \right\}, \quad (232)$$

and using (413) we obtain the inequality

$$|\partial_A \Gamma_l^{\Lambda'\Lambda_0;c\bar{c}A}(\vec{p})| \leq P_{2l}^{\Lambda'\Lambda'}(\vec{p}) \leq P_{2l}^{\Lambda}(\vec{p}). \quad (233)$$

If $\Lambda' = \Lambda$ the proof is finished. If not we use (226).

4.4.3 Renormalization at $\Lambda = 0$

First we consider the marginal terms $\Gamma_{\omega;l}^{cc}$, $\Gamma_{\gamma;l}^{cA}$, Γ_l^{AAAA} , $\partial_{\bar{c}}\Gamma_l^{c\bar{c}A}$, $\partial\Gamma_l^{AAA}$, $\partial\partial\Gamma_l^{AA}$, $\partial\partial\Gamma_l^{c\bar{c}}$, here denoted by $\Gamma_{\vec{z};l}^{\vec{\phi};w}$. The marginal terms $\Gamma_{1\vec{z};l}^{\vec{\phi}}$ and $\Gamma_{\beta;l}^{\vec{\phi}}$ will be discussed later.

Let $\{\delta^s \underline{e}\}_r$ be a basis at the renormalization point \vec{q} . We define $\Gamma_{\vec{z};l}^{\vec{\phi};w}(\vec{q})$ in the following way:

$$\Gamma_{\vec{z};l}^{0\Lambda_0;\vec{\phi};w}(\vec{q}) := \sum_{t \in \{\delta^s\}_r} r_t t + \sum_{t\alpha \in \{\delta^s \underline{e}^k > 0\}_r} \zeta_t t, \quad (234)$$

where the coefficients r_t are fixed by the renormalization conditions, see appendix H and hypothesis RC3, and the remaining coefficients ζ_t are defined using lemma 30. Then, from the bounds on the irrelevant terms and lemma 31, see (367), it follows that $\Gamma_{\vec{z};l}^{\vec{\phi};w}(\vec{q})$ complies with theorem 1 at loop order l . Let $\Lambda' := \max(\Lambda, \eta(\vec{p}))$. It is easy to verify the following inequalities

$$\int_0^{\Lambda'} d\lambda \frac{\log^k \frac{M}{\lambda}}{\lambda + M} < k! + \log_+ \frac{\Lambda'}{M}, \quad \int_0^{\Lambda'} d\lambda \frac{\log^k \frac{\lambda}{M}}{\lambda + M} < 2 \left(1 + \log_+^{k+1} \frac{\Lambda'}{M} \right). \quad (235)$$

Recalling (198), (200) we obtain the following bound

$$|\dot{\Gamma}_{\vec{z};l}^{\lambda\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \lambda \frac{\mathcal{P}_{2l-2}^{(0)}(\log_+ \frac{M}{\lambda}) + \mathcal{P}_{2l-2}^{(1)}(\log_+(\frac{\lambda}{M}))}{(\lambda + M)^2}. \quad (236)$$

This implies that

$$|\Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{q}) - \Gamma_{\vec{z};l}^{0\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \int_0^{\Lambda'} d\lambda |\dot{\Gamma}_{\vec{z};l}^{\lambda\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \mathcal{P}_{2l-1}^{(1)} \left(\log_+ \frac{\Lambda'}{M} \right). \quad (237)$$

Using inequality (413) we get

$$|\Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0) - \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \sum_{j=1}^{n-1} I_j(\vec{q}) \leq P_{2l}^{\Lambda'}(0), \quad (238)$$

$$|\Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0)| \leq \sum_{j=1}^{n-1} I_j(\vec{p}) \leq P_{2l}^{\Lambda'}(\vec{p}), \quad (239)$$

where $I_j(\vec{q})$ is the interpolation along the vector q_j ,

$$I_j(\vec{q}) = \int_0^1 dt |q_j| |\partial_j \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\sum_{i=1}^{j-1} \vec{q}_i + t\vec{q}_j)|. \quad (240)$$

Here to each vector $q_i \in \mathbb{R}^4$ is associated $\vec{q}_i \in \mathbb{P}_n$ whose components are $(\vec{q}_i)_k = -q_i \delta_{k,0} + q_i \delta_{k,i}$. Once again we have substituted $|\partial_j \Gamma_{\vec{z};l}^{M\Lambda_0;\vec{\phi};w}|$ with the bound from theorem 1. If $\Lambda = \Lambda'$ we stop here. If not we use (226).

The remaining marginal terms $\Gamma_{X\vec{x};l}^{\Lambda\Lambda_0;\vec{\phi};w}$ with $X \in \{\beta, 1\}$ can be treated similarly. Note that the bound from theorem 1 for terms of the type $\partial_{\vec{c}}\Gamma_{1;l}^{\vec{c}\vec{\phi};w}$ is the same as the one for the corresponding terms $\Gamma_{1\vec{x};l}^{\vec{\phi};w}$. Consequently the proof of the bounds for the marginal terms $\Gamma_{1\gamma;l}^{ccA}$, $\Gamma_{1\omega;l}^{ccc}$, $\partial\Gamma_{1\gamma;l}^{cc}$ is the same as the proof for respectively $\partial_{\vec{c}}\Gamma_{1;l}^{\vec{c}ccA}$, $\partial\partial_{\vec{c}}\Gamma_{1;l}^{\vec{c}cc}$ which we shall consider now.

Let us denote by $\Gamma_{X;l}^{0\Lambda_0;\vec{\phi};w}$ any marginal term without insertions γ or ω , and by \vec{q} the corresponding renormalization point as given in appendix H. We now anticipate the important fact that the relevant renormalization constants comply with the bounds, which will be proven in section 5. Then using the bounds on irrelevant terms from theorems 3, 4 and lemma 31, see (367), we obtain

$$|\Gamma_{X;l}^{0\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \frac{M}{\Lambda_0} \mathcal{P}_{r_X}^{(1)}\left(\log_+ \frac{\Lambda_0}{M}\right), \quad (241)$$

in agreement with theorems 3, 4. Set $\Lambda' := \max(\Lambda, M)$. To integrate the FE from 0 to Λ' we substitute the chain of vertex functions with the trees from theorems 3, 4. Then we have

$$|\Gamma_{X;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{q}) - \Gamma_{X;l}^{0\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \frac{\Lambda + M}{\Lambda_0} \mathcal{P}_{r_X}^{(1)}\left(\log_+ \frac{\Lambda_0}{M}\right). \quad (242)$$

Using inequality (413) we obtain

$$|\Gamma_{X;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0) - \Gamma_{X;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \frac{\Lambda + M}{\Lambda_0} P_{r_X}^{\Lambda'\Lambda_0}(0), \quad (243)$$

$$|\Gamma_{X;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0) - \Gamma_{X;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{p})| \leq F_{s_X}^{\Lambda\Lambda_0}(\vec{p}) P_{r_X}^{\Lambda\Lambda_0}(\vec{p}). \quad (244)$$

Note here that as compared to the case $X = \beta$, in the case $X = 1$ there appears an additional factor

$$\frac{\Lambda' + |p_*|}{\Lambda' + |p_a|} \leq 1 + 2 \frac{|\vec{p}|}{\Lambda + M}, \quad (245)$$

in the tree bounds from theorem 4. This factor leads to the polynomial $\mathcal{P}_{s_1}^{(2)}$ in the bounds of theorem 4.

If $\Lambda = \Lambda'$ the proof is finished. If not we integrate downwards using the FE, substitute the chain with the tree bound from theorems 3, 4 and

use (401) to get

$$|\Gamma_{\beta;l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_{\beta;l}^{M\Lambda_0;\vec{\phi};w}(\vec{p})| \leq F_0^{\Lambda\Lambda_0}(\vec{p}) P_{2l-1}^{\Lambda\Lambda_0}(\vec{p}), \quad (246)$$

$$|\Gamma_{1;l}^{\Lambda\Lambda_0;\vec{\phi};w}(0) - \Gamma_{1;l}^{M\Lambda_0;\vec{\phi};w}(\vec{p})| \leq \sum_{\tau \in \overline{T}_{1\vec{\phi}}} Q_{\tau}^{\Lambda;w} F_{l-1}^{\Lambda\Lambda_0} P_{3l-2}^{\Lambda\Lambda_0}(\vec{p}). \quad (247)$$

4.5 Strictly relevant terms

If $n_{\mathcal{Z}} = 0$ the notation $\Gamma_{\vec{z};l}^{\vec{\phi};w}$ stands for $\Gamma_l^{c\bar{c}A}$, Γ_l^{AAA} , $\partial\Gamma_l^{AA}$, $\partial\Gamma_l^{c\bar{c}}$. In the case $n_{\mathcal{Z}} = 1$, it stands for $\Gamma_{\gamma;l}^c$. Moreover we impose $\Gamma_{\vec{z};l}^{0\Lambda_0\vec{\phi};w}(0) = 0$ and denote by \vec{p} arbitrary momenta with corresponding $\eta(\vec{p})$. We integrate the FE upwards from 0 to Λ and substitute $|\dot{\Gamma}_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}|$ with the tree bound of theorem 1. Then

$$|\Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(0)| \leq \int_0^{\Lambda} d\lambda |\dot{\Gamma}_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(0)| \leq \int_0^{\Lambda} d\lambda \lambda^{d-1} P_{2l-2}^{\Lambda\Lambda}(0) \leq \int_0^{\Lambda+\eta} d\lambda \lambda^{d-1} P_{2l-2}^{\Lambda\Lambda}(0),$$

where $d > 0$. Inequality (410) then gives

$$|\Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(0)| \leq (\Lambda + \eta)^d P_{2l-2}^{\Lambda+\eta\Lambda}(0). \quad (248)$$

Substituting $\partial\Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}$ with the bound from theorem 1 we obtain

$$\begin{aligned} |\Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(\vec{p}) - \Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(0)| &\leq \int_0^1 dt |p_i| |\partial_i \Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(t\vec{p})| \\ &\leq |\vec{p}| \int_0^1 dt (\Lambda + t|\vec{p}|)^{d-1} P_{2l}^{\Lambda}(t\vec{p}) \leq I + I', \end{aligned} \quad (249)$$

where

$$I := |\vec{p}| (\Lambda + |\vec{p}|)^{d-1} \int_0^1 dt \mathcal{P}_{2l}^{(0)} \left(\log_+ \frac{\max(M, |\vec{p}|)}{\Lambda + t\eta(\vec{p})} \right), \quad (250)$$

$$I' := (\Lambda + |\vec{p}|)^d \mathcal{P}_{2l}^{(1)} \left(\log_+ \frac{\Lambda}{M} \right). \quad (251)$$

The calculations given in (410) yield

$$I \leq |\vec{p}|(\Lambda + |\vec{p}|)^{d-1} \mathcal{P}_{2l}^{(0)}(\log_+ \frac{\max(M, |\vec{p}|)}{\Lambda + \eta(\vec{p})}). \quad (252)$$

This implies

$$|\Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(\vec{p}) - \Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(0)| \leq (\Lambda + |\vec{p}|)^d P_{2l}^\Lambda(\vec{p}). \quad (253)$$

Combining (248) and (253) proves the bounds of theorem 1 to loop order l .

4.5.1 $\partial\Gamma_1^{\phi^3}$, $\Gamma_1^{\phi^4}$ and $\Gamma_{1\gamma}^{cc}$

The goal of this section is to explain the expression for the polynomial degree r_1 . Subsequently we denote by $\Gamma_1^{\vec{\phi};w}$ the following terms: $\partial\Gamma_1^{cAA}$, $\partial\Gamma_1^{cc\bar{c}}$, Γ_1^{cAAA} , $\Gamma_1^{cAc\bar{c}}$ and we impose vanishing renormalization conditions at the origin, $\Gamma_1^{0\Lambda_0;\vec{\phi};w}(0) = 0$. From the bounds of theorem 4 one realises that the analysis for $\Gamma_{1\gamma}^{cc}$ is similar to $\partial_{\bar{c}}\Gamma_1^{cc\bar{c}}$. For an arbitrary $\vec{p} \in \mathbb{P}_n$ let $\Lambda' := \max(\Lambda, \eta(\vec{p}))$. Using (198), (199)

$$|\Gamma_{1;l}^{\Lambda'\Lambda_0\vec{\phi};w}(0)| \leq (\Lambda + \eta) \frac{\Lambda + M}{\Lambda_0} P_{3(l-1)}^{\Lambda+\eta\Lambda_0}(0). \quad (254)$$

We extend $\Gamma_{1;l}^{\Lambda'\Lambda_0;\vec{\phi};w}$ from 0 to \vec{p} using the usual path given in equation (240) (with $\vec{q} \mapsto \vec{p}$), where we substitute $\partial\Gamma_{1;l}^{\Lambda'\Lambda_0\vec{\phi};w}$ with the tree bound of theorem 4. The rational factor in this bound makes these terms different from other strictly relevant terms, for example $\partial\partial\Gamma_1^{cA}$, Γ_β^{cA} or $\partial\Gamma_\beta^c$. Then using inequality (413) we have

$$|\Gamma_{1;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_{1;l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0)| \leq (\Lambda + |\vec{p}|) \left(1 + \log_+ \frac{|\vec{p}|}{\Lambda + \eta}\right) F_l^{\Lambda\Lambda_0}(\vec{p}) P_{3l-1}^{\Lambda\Lambda_0}(\vec{p}).$$

If $\Lambda = \Lambda'$ the proof is finished. If not we repeat the steps used to show (226) integrating the FE from η to Λ and using (399)

$$|\Gamma_{1;l}^{\Lambda\Lambda_0;\vec{\phi};w}(\vec{p}) - \Gamma_{1;l}^{\eta\Lambda_0;\vec{\phi};w}(\vec{p})| \leq |\vec{p}| F_{l-1}^{\Lambda\Lambda_0}(\vec{p}) P_{3l-2}^{\Lambda\Lambda_0}(\vec{p}).$$

4.5.2 $\Gamma_{\beta}^{\vec{\phi};w}$, $\Gamma_1^{\phi^2;w}$ and $\Gamma_1^{\phi^3}$

In this section we briefly discuss the remaining strictly relevant terms $\Gamma_{\beta}^{\vec{\phi};w}$, $\Gamma_1^{\phi^3}$ and $\Gamma_1^{\phi^2;w}$ denoting all of them by $\Gamma_X^{\vec{\phi};w}$ with $X \in \{\beta, 1\}$. We impose renormalization conditions $\Gamma_X^{0\Lambda_0; \vec{\phi}; w}(0) = 0$ and integrate the FE from 0 to Λ . We use the bounds of theorems 3,4 and then for an arbitrary \vec{p} extend the integration up to $\Lambda + \eta(\vec{p})$

$$|\Gamma_X^{\Lambda\Lambda_0; \vec{\phi}; w}(0)| \leq (\Lambda + \eta)^{d_X} \frac{\Lambda + M}{\Lambda_0} P_{r_X}^{\Lambda + \eta \Lambda_0}(0). \quad (255)$$

Integration along the path given in equation (240) and using (410) yields

$$|\Gamma_X^{\Lambda\Lambda_0; \vec{\phi}; w}(\vec{p}) - \Gamma_X^{\Lambda\Lambda_0; \vec{\phi}; w}(0)| \leq (\Lambda + |\vec{p}|)^{d_X} F_{s_X}(\vec{p}) P_{s_X}^{\Lambda\Lambda_0}(\vec{p}). \quad (256)$$

4.6 Convergence

We first prove the bounds for $\partial_{\Lambda_0} \Gamma^{\vec{\phi}}$. Then we proceed with the other functions $\partial_{\Lambda_0} \Gamma_{\vec{z}}^{\vec{\phi}}$ ascending in the number of insertions $n_{\mathcal{z}}$. We use the same inductive scheme as before, based on the FE.

We start with the irrelevant terms integrating the FE from Λ_0 to Λ , using the boundary conditions $\Gamma_l^{\Lambda_0\Lambda_0; \vec{\phi}; w} = 0$ and applying the derivative wrt Λ_0 .

$$\partial_{\Lambda_0} \Gamma_{\vec{z}; l}^{\Lambda\Lambda_0; \vec{\phi}; w} = -\dot{\Gamma}_{\vec{z}; l}^{\Lambda_0\Lambda_0; \vec{\phi}; w} + \int_{\Lambda_0}^{\Lambda} d\lambda \partial_{\Lambda_0} \dot{\Gamma}_{\vec{z}; l}^{\lambda\Lambda_0; \vec{\phi}; w}. \quad (257)$$

To bound the first term of the expression we substitute into the FE the irrelevant tree bound

$$|\dot{\Gamma}_{\vec{z}; l}^{\Lambda_0\Lambda_0; \vec{\phi}; w}| \leq \frac{1}{\Lambda_0} \sum_{\tau \in \mathcal{T}_{\vec{z}\vec{\phi}}} Q_{\tau}^{\Lambda_0; w} P_{2(l-1)}^{\Lambda_0\Lambda_0}. \quad (258)$$

If $n_{\mathcal{z}} = 0$ then $\tau \in \mathcal{T}_{\vec{\phi}}$ and $\theta(\tau) > 0$. Consequently, recalling (151)

$$\Pi_{\tau, \theta}^{\Lambda_0}(\vec{p}) \leq \frac{\Lambda + |\vec{p}|}{\Lambda_0} \Pi_{\tau, \theta}^{\Lambda}(\vec{p}), \quad \text{and thus} \quad Q_{\tau}^{\Lambda_0; w} \leq \frac{\Lambda + |\vec{p}|}{\Lambda_0} Q_{\tau}^{\Lambda; w}. \quad (259)$$

Otherwise, the denominator $\Lambda_0 + |p_e|$ with $e_x \in E_x$ gives the inequality

$$\frac{1}{\Lambda_0 + |p_x|} \leq \frac{\Lambda + |\vec{p}|}{\Lambda_0} \frac{1}{\Lambda + |p_x|}. \quad (260)$$

In both cases this yields

$$|\dot{\Gamma}_{\vec{z};l}^{\Lambda_0\Lambda_0;\vec{\phi};w}| \leq \frac{\Lambda + |\vec{p}|}{\Lambda_0^2} \sum_{\tau \in \mathcal{T}_{\vec{z}\vec{\phi}}} Q_\tau^{\Lambda;w} P_{2(l-1)}^{\Lambda\Lambda_0}. \quad (261)$$

To analyse the second term we apply $\partial_p^w \partial_{\Lambda_0}$ to the chain of vertex functions given in definition 20. This gives a chain with the element $\partial_{\Lambda_0}((\partial_p^{w_1} \mathbf{C})\Gamma_{l'}^{w_2})$ $l' < l$ which we bound using (475), (482) and theorem 2.

$$\begin{aligned} |\partial_{\Lambda_0}((\partial_p^{w_1} \mathbf{C})\Gamma_{\vec{z};l'}^{\lambda\Lambda_0;\zeta\vec{\phi}\bar{\zeta};w_2})| &\leq |\partial_{\Lambda_0} \partial_p^{w_1} \mathbf{C}| |\Gamma_{\vec{z};l'}^{\lambda\Lambda_0;\zeta\vec{\phi}\bar{\zeta};w_2}| + |\partial_p^{w_1} \mathbf{C}| |\partial_{\Lambda_0} \Gamma_{\vec{z};l'}^{\lambda\Lambda_0;\zeta\vec{\phi}\bar{\zeta};w_2}| \\ &\leq \frac{c}{(\lambda + |p_\zeta|)^2 + \|w_1\|} \left(\frac{\lambda + |p_\zeta|}{\Lambda_0^2} |\Gamma_{\vec{z};l'}^{\lambda\Lambda_0;\zeta\vec{\phi}\bar{\zeta};w_2}| + |\partial_{\Lambda_0} \Gamma_{\vec{z};l'}^{\lambda\Lambda_0;\zeta\vec{\phi}\bar{\zeta};w_2}| \right) \\ &\leq \frac{1}{(\lambda + |p_\zeta|)^2} \frac{\lambda + M + |\vec{p}|}{\Lambda_0^2} \sum_{\tau \in \mathcal{T}_{\vec{z}\zeta\vec{\phi}\bar{\zeta}}} Q_\tau^{\lambda;w_2+w_1} P_{r(l')}^{\lambda\Lambda_0}(\vec{p}). \end{aligned} \quad (262)$$

We proceed as for the proof of inequality (217) substituting expression (221) with the integral

$$\int_{\Lambda}^{\Lambda_0} d\lambda \frac{(\lambda + M + |\vec{p}|) P_{2l-2}^{\lambda\Lambda_0} |_{\eta_\tau=0}}{\Lambda_0^2 (\lambda + |p_a|)(\lambda + |p_b|)} \leq \frac{\Lambda + M + |\vec{p}|}{\Lambda_0^2} \frac{1}{\Lambda + |\vec{p}|} P_{2l-1}^{\Lambda\Lambda_0}(\vec{p}), \quad (263)$$

see the explanation after (221). If $\tau \in \mathcal{T}_{x\phi\phi}$ we always choose $\hat{p} = p_x$ in order to preserve the factor $\Lambda + |p_x|$ in the denominator of Q_τ^Λ .

For the marginal terms we shall integrate the FE upwards from 0 to Λ . For the terms with the antighost we use renormalization conditions (144)

$$\Gamma_l^{M\Lambda_0;c\bar{c}AA}(0) = 0, \quad \Gamma_l^{M\Lambda_0;c\bar{c}c\bar{c}}(0) = 0, \quad \partial_A \Gamma_l^{M\Lambda_0;c\bar{c}A}(0) = 0. \quad (264)$$

Using equation (145) we obtain that for these terms at $\Lambda' = \max(\Lambda, \eta(\vec{p}))$

$$\begin{aligned} |\partial_{\Lambda_0} \Gamma_l^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{p})| &\leq \int_0^1 dt |p_{\bar{e}}| |\partial_{\Lambda_0} \partial_{\bar{e}} \Gamma_l^{\Lambda'\Lambda_0;\vec{\phi};w}(tp_{\bar{e}}, \dots)| \\ &\leq \frac{\Lambda' + M + |\vec{p}|}{\Lambda_0^2} \sum_{\tau \in \mathcal{T}_{\vec{\phi}}^0} \int_0^1 dt |p_{\bar{e}}| Q_\tau^{\Lambda';w+1_{\bar{e}}} P_{2l-1}^{\Lambda'\Lambda_0} \leq \frac{\Lambda + M + |\vec{p}|}{\Lambda_0^2} P_{2l}^{\Lambda\Lambda_0}, \end{aligned} \quad (265)$$

where we have substituted $\partial_{\Lambda_0} \partial_{\vec{e}} \Gamma_l^{\vec{\phi};w}$ with the bound of theorem 2 and applied inequality (413).

The remaining marginal terms Γ^{AAAA} , $\partial_{\vec{e}} \Gamma^{c\bar{c}A}$, $\Gamma_{\vec{z}}^{\phi\phi}$, $\partial\partial\Gamma^{\phi\phi}$, $\partial\Gamma_{\gamma}^c$, are renormalized at $\Lambda = 0$ and nonvanishing momentum \vec{q} , chosen in \mathbb{M}_n^s in all cases but Γ^{AAAA} for which $\vec{q} \in \mathbb{M}_4^{cp}$. See appendix H for the list of all relevant terms and their renormalization points. Since the renormalization constants are independent of Λ_0 , their derivative wrt Λ_0 vanishes: it follows that the coefficients of δ -tensors in the decomposition of $\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{0\Lambda_0;\vec{\phi};w}(\vec{q})$ vanish. Hence using lemma 31 and the bounds on irrelevant terms we have

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{0\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \frac{M}{\Lambda_0^2} \mathcal{P}_{2l-1}^{(1)}(\log_+ \frac{\Lambda_0}{M}). \quad (266)$$

We integrate the FE from 0 to Λ' and substitute the chain with the tree bound. Using inequalities (235), (236) it is easy to get the following bound

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{q}) - \partial_{\Lambda_0} \Gamma_{\vec{z};l}^{0\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \frac{\Lambda + M}{\Lambda_0^2} \mathcal{P}_{2l-1}^{(1)}(\log_+ \frac{\Lambda_0}{M}). \quad (267)$$

Integrating back and forth along the path given in equation (240), substituting the irrelevant term $\partial_{\Lambda_0} \partial \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0\vec{\phi};w}$ with its bounds and using inequality (413) we obtain

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0) - \partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{q})| \leq \frac{\Lambda + M}{\Lambda_0^2} P_{2l}^{\Lambda'\Lambda_0}(0), \quad (268)$$

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(\vec{p}) - \partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda'\Lambda_0;\vec{\phi};w}(0)| \leq \frac{\Lambda + M + |\vec{p}|}{\Lambda_0^2} P_{2l}^{\Lambda\Lambda_0}(\vec{p}). \quad (269)$$

If $\Lambda = \Lambda'$ the proof of the bounds on the marginal terms is complete. Otherwise we integrate the FE downwards from η to Λ and repeat the arguments given to prove inequality (226) with a minor change in the integrand

$$\int_{\Lambda}^{\eta} d\lambda \frac{\lambda + M + |\vec{p}|}{\Lambda_0^2} \frac{\lambda P_k^{\Lambda\Lambda_0}|_{\eta\tau=0}}{(\lambda + |p_a|)(\lambda + |p_b|)} \leq \frac{M + |\vec{p}|}{\Lambda_0^2} \int_{\Lambda}^{\eta} d\lambda \frac{2P_k^{\Lambda\Lambda_0}|_{\eta\tau=0}}{\lambda + |p_b|}. \quad (270)$$

For the strictly relevant terms we integrate the FE from 0 to Λ substituting the vertex functions and propagators with their bounds and extending the upper limit of integration to $\Lambda + \eta$

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda\Lambda_0\vec{\phi};w}(0)| \leq \frac{\Lambda + M}{\Lambda_0^2} \int_0^{\Lambda+\eta} d\lambda \lambda^{d-1} P_{2l-2}^{\Lambda\Lambda_0}(0), \quad (271)$$

where $d > 0$. Using inequality (410) we obtain

$$|\partial_{\Lambda_0} \Gamma_{\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}(0)| \leq \frac{\Lambda + M}{\Lambda_0^2} (\Lambda + \eta)^d P_{2l-2}^{\Lambda+\eta\Lambda_0}(0). \quad (272)$$

To extend to momentum \vec{p} we proceed as in (249), the only change being an additional factor of $\frac{\Lambda+M+|\vec{p}|}{\Lambda_0^2}$.

5 Restoration of the STI

As mentioned before theorems 3, 4 we now consider all nontrivial marginal terms $\Gamma_{X\vec{z};l}^{\Lambda\Lambda_0;\vec{\phi};w}$ with $X \in \{\beta, 1\}$ at $\Lambda = 0$. We want to show that these terms verify the bounds of theorems 3, 4. In this section, since $\Lambda = 0$, we will omit the parameters Λ, Λ_0 in the notations wherever this is not ambiguous, i.e. we write $\Gamma_{\vec{\phi}}$ for $\Gamma^{0\Lambda_0;\vec{\phi}}$.

The subsequent relations are obtained by projecting the AGE (123) and the STI (124) on the respective monomial in the fields (for example, c in section 5.1) to read off the lhs from the rhs, taken at the renormalization point. We will establish appropriate relations in order to make the coefficients of the δ -tensors wrt the monomial basis at the renormalization point satisfy the bounds of these two theorems. In this analysis we make particular use of the consistency conditions, see 5.0.3. In section 5.0.4 we prove the existence of a solution for the above mentioned system of relations that does not depend on the UV cutoff. In the remaining sections we treat the different marginal terms one by one.

5.0.1 Smallness relations

It is helpful to introduce the notion of *small* terms, which vanish in the limit $\Lambda_0 \rightarrow \infty$. For fixed loop order l and $X \in \{\beta, 1\}$, a homogeneous function $f(\vec{p}, M, \Lambda_0)$ of mass dimension $[f]$ is said *small* on a subset $Y \subset \mathbb{M}_n$, and denoted by $f \stackrel{X,Y,l}{\sim} 0$, if for all $w \in \mathbb{W}_n$ with $\|w\| \leq [f]$, there exists a polynomial $\mathcal{P}_{r_X}^{(1)}$ of degree $r_X([f] - \|w\|, l)$, see theorems 3 and 4, such that the following bound holds for all $\Lambda_0 \geq M$ and all $\vec{p} \in Y$, see (131):

$$|\partial_p^w f(\vec{p}, M, \Lambda_0)| \leq \frac{M^{1+[f]-\|w\|}}{\Lambda_0} \mathcal{P}_{r_X}^{(1)}(\log_+ \frac{\Lambda_0}{M}). \quad (273)$$

Furthermore, $f \stackrel{X,Y,l}{\sim} g$ iff $[f] = [g]$ and $f - g \stackrel{X,Y,l}{\sim} 0$. Because both relations $\stackrel{1,Y,l}{\sim}$ and $\stackrel{\beta,Y,l}{\sim}$ only differ by the degree of polynomials we have $f \stackrel{\beta,Y,l}{\sim} g \implies f \stackrel{1,Y,l}{\sim} g$. Since the loop order l and the renormalization point \vec{q} are evident from the context, we write:

$$f \stackrel{Y}{\sim} g \quad \text{for} \quad f \stackrel{1,Y,l}{\sim} g, \quad f \stackrel{\beta,Y}{\sim} g \quad \text{for} \quad f \stackrel{\beta,Y,l}{\sim} g, \quad (274)$$

$$f \sim g \quad \text{for} \quad f \stackrel{1,\{\vec{q}\}}{\sim} g, \quad f \stackrel{\beta}{\sim} g \quad \text{for} \quad f \stackrel{\beta,\{\vec{q}\}}{\sim} g. \quad (275)$$

Theorem 1 implies that for every vertex function $\Gamma_{\vec{z}}^{\vec{\phi}}(\vec{p})$ there exists a constant c such that $\forall w \in \mathbb{W}_n, \forall \vec{p} \in \mathbb{M}_n, \forall \Lambda_0 \geq M$

$$|\partial^w \Gamma_{\vec{z}}^{\vec{\phi}}(\vec{p})| \leq c M^{4-2n_{\vec{z}}-N-\|w\|}. \quad (276)$$

Using also that

$$|\partial^w (\sigma_{0\Lambda_0} - 1)| < c_w \frac{1}{\Lambda_0^w}, \quad |\sigma_{0\Lambda_0} - 1| < c_0 \frac{M}{\Lambda_0}, \quad (277)$$

the terms on rhs of the STI and the AGE satisfy the relations:

$$\Gamma_{\vec{z}_1;w_1}^{\vec{\phi}_1} \partial^{w_2} (\sigma_{0\Lambda_0} \Gamma_{\vec{z}_2}^{\vec{\phi}_2}) \stackrel{\beta,Y}{\sim} \Gamma_{\vec{z}_1;w_1}^{\vec{\phi}_1} \Gamma_{\vec{z}_2;w_2}^{\vec{\phi}_2}, \quad (278)$$

$$\partial^w (\sigma_{0\Lambda_0} \Gamma_{\vec{z}}^{\vec{\phi}}) \stackrel{\beta,Y}{\sim} \Gamma_{\vec{z};w}^{\vec{\phi}}, \quad (279)$$

$$\partial^w (p \sigma_{0\Lambda_0} \Gamma_{\vec{\gamma}}^{\vec{\phi}}) \stackrel{\beta,Y}{\sim} p \Gamma_{\vec{\gamma};w}^{\vec{\phi}}. \quad (280)$$

This fact will be useful in the calculations underlying the following sections.

5.0.2 The functional $F_{1,rel}$

In this section we introduce the notation for the renormalization constants for the functional F_1 (119). For this purpose we define the auxiliary functional

$$\begin{aligned} F_{1,rel}^{\Lambda_0} := & \epsilon^{dab} \langle U_{\sigma}^{\gamma cc} | \check{\gamma}_{\sigma}^d \check{c}^a \check{c}^b \rangle + \epsilon^{sab} \epsilon^{sde} u^{\gamma Acc} \langle \check{\gamma}_{\sigma}^a \check{A}_{\sigma}^b \check{c}^d \check{c}^e \rangle \\ & + \langle U_{\sigma}^{cA} | \check{c}^d \check{A}_{\sigma}^d \rangle + \epsilon^{dab} \langle U^{\bar{c}cc} | \check{c}^d \check{c}^a \check{c}^b \rangle + \epsilon^{dab} \langle U_{\mu\nu}^{cAA} | \check{c}^d \check{A}_{\mu}^a \check{A}_{\nu}^b \rangle \\ & + \langle U_{\rho\mu\nu}^{cAAA} | \check{c}^d \check{A}_{\rho}^d \check{A}_{\mu}^a \check{A}_{\nu}^a \rangle + \langle U_{1;\mu}^{\bar{c}ccA} | \check{c}^b \check{c}^b \check{c}^a \check{A}_{\mu}^a \rangle \\ & + \langle U_{2;\mu}^{\bar{c}ccA} | \check{c}^a \check{c}^b \check{c}^b \check{A}_{\mu}^a \rangle + \epsilon^{bed} u_1^{\bar{c}ccAA} \langle \check{c}^b \check{c}^e \check{c}^d \check{A}_{\mu}^a \check{A}_{\mu}^a \rangle \\ & + \epsilon^{bed} u_2^{\bar{c}ccAA} \langle \check{c}^b \check{c}^e \check{c}^d \check{A}_{\mu}^a \check{A}_{\mu}^e \rangle + \epsilon^{bed} u_3^{\bar{c}ccAA} \langle \check{c}^a \check{c}^e \check{c}^d \check{A}_{\mu}^a \check{A}_{\mu}^b \rangle, \end{aligned} \quad (281)$$

where

$$U_\sigma^{\gamma cc}(l, p, q) := i(p + q)_\sigma u^{\gamma cc}, \quad (282)$$

$$U_\sigma^{cA}(l, p) := ip_\sigma p^2 u^{cA}, \quad (283)$$

$$U^{\bar{c}cc}(l, p, q) := (p^2 + q^2)u_1^{\bar{c}cc} + 2pq u_2^{\bar{c}cc}, \quad (284)$$

$$U_{\mu\nu}^{cAA}(l, p, q) := (p_\mu p_\nu - q_\mu q_\nu)u_1^{cAA} + \delta_{\mu\nu}(p^2 - q^2)u_2^{cAA}, \quad (285)$$

$$\begin{aligned} U_{\rho\mu\nu}^{cAAA}(l, k, p, q) &:= i(u_1^{cAAA}k_\rho + u_2^{cAAA}p_\rho + u_2^{cAAA}q_\rho)\delta_{\mu\nu} \\ &\quad + i(u_3^{cAAA}k_\nu + u_4^{cAAA}q_\nu + u_5^{cAAA}p_\nu)\delta_{\rho\mu} \\ &\quad + i(u_3^{cAAA}k_\mu + u_4^{cAAA}p_\mu + u_5^{cAAA}q_\mu)\delta_{\rho\nu}, \end{aligned} \quad (286)$$

$$U_{1;\mu}^{\bar{c}ccA}(l, k, q, p) := ip_\mu u_1^{\bar{c}ccA} + ik_\mu u_2^{\bar{c}ccA} + iq_\mu u_3^{\bar{c}ccA}, \quad (287)$$

$$U_{2;\mu}^{\bar{c}ccA}(l, k, q, p) := i(k_\mu - q_\mu)u_4^{\bar{c}ccA}, \quad (288)$$

and the u 's are functions of Λ_0 defined by the marginal renormalization conditions

$$F_{1;\vec{z}}^{0\Lambda_0;\vec{\phi};w}(\vec{q}) = F_{1,rel;\vec{z}}^{\Lambda_0;\vec{\phi};w}(\vec{q}) + \sum_{t \in \{\delta^s \underline{e}^k > 0\}_r} \zeta_t^{\Lambda_0} t. \quad (289)$$

Here $2n_\varkappa + N + \|w\| = 5$, \vec{q} is the renormalization point defined in section H, $\underline{e} = (e_i)_{i \in [m]}$ is an orthogonal basis for the linear span of \vec{q} , r is the tensor rank of $F_{1;\vec{z}}^{0\Lambda_0;\vec{\phi};w}(\vec{q})$. The $\zeta_t^{\Lambda_0}$ are the uniquely defined coefficients of tensors t . Note that we implicitly set to zero all constants associated to strictly relevant renormalization conditions for F_1 . These constants are not needed because, thanks to hypothesis RC1, the RHS of the STI and AGE at the current loop order vanish at zero momenta.

The renormalization points \vec{q} are chosen in agreement with the hypotheses of lemma 31. From lemma 31, theorem 3, and the irrelevant bounds of theorem 4, for the marginal terms one has

$$F_{1;\vec{z}}^{0\Lambda_0;\vec{\phi};w}(\vec{q}) \sim F_{1,rel;\vec{z}}^{\Lambda_0;\vec{\phi};w}(\vec{q}). \quad (290)$$

5.0.3 Consistency conditions

Here we establish the consistency conditions implied by the nilpotency, see (121) and (122). Below we will rely on the validity of theorem 4 at loop orders $l' < l$ for all terms and at the current order l only for irrelevant terms: these properties are true in our inductive scheme. Recall definitions (116),

(131) and (274). Using the AGE (123), the bounds of theorems 1-3, and (122) we get

$$\left(S\Gamma_\beta\right)^{\vec{\phi}} \stackrel{\mathbb{M}_n}{\sim} 0, \quad \text{and thus} \quad \left(\frac{\delta}{\delta\bar{c}}F_1\right)^{\vec{\phi}} \stackrel{\mathbb{M}_n}{\sim} 0, \quad (291)$$

where $\vec{\phi} = (\phi_1, \dots, \phi_{n-1})$ and $\phi_i \in \{A, c, \bar{c}\}$. Equation (291), theorems 1, 2, and the bounds of theorem 4 for irrelevant terms yield

$$\left(\frac{\delta}{\delta\bar{c}}F_{1,rel}\right)^{\vec{\phi}} \stackrel{\mathbb{M}_n}{\sim} 0, \quad \vec{\phi} \in \{(c, c), (c, c, A)\}. \quad (292)$$

See section 5.0.2 for the definition of $F_{1,rel}$ and of the constants u^Φ . In section 5.5 it will be shown that $u_{1,2,3}^{\bar{c}c c A A} \sim 0$. Equation (292) then gives

$$u^{\gamma cc} \sim u_1^{\bar{c}cc} \sim u_2^{\bar{c}cc}, \quad -2u^{\gamma Acc} \sim u_1^{\bar{c}c c A} \sim u_2^{\bar{c}c c A} \sim u_3^{\bar{c}c c A}, \quad u_4^{\bar{c}c c A} \sim 0. \quad (293)$$

Let us exploit (121) to obtain more constraints on the renormalization constants u^Φ . At loop order l

$$SF_{1;l} = S_0F_{1;l} + \sum_{l' < l} S_{l-l'}F_{1;l'}. \quad (294)$$

By induction $(S_{l-l'}F_{1;l'})^{\vec{\phi}} \stackrel{\mathbb{M}_n}{\sim} 0$ for all loop orders $l' < l$. Then equation (121) implies that

$$(S_0F_{1;l})^{\vec{\phi}} \stackrel{\mathbb{M}_n}{\sim} 0, \quad (295)$$

where

$$\begin{aligned} S_0 = & \langle \tilde{\delta}_{A_\sigma^d} \Gamma_0, \sigma_{0\Lambda_0} \tilde{\delta}_{\gamma_\sigma^d} \rangle + \langle \tilde{\delta}_{\gamma_\sigma^d} \Gamma_0, \sigma_{0\Lambda_0} \tilde{\delta}_{A_\sigma^d} \rangle \\ & - \langle \tilde{\delta}_{c^d} \Gamma_0, \sigma_{0\Lambda_0} \tilde{\delta}_{\omega^d} \rangle - \langle \tilde{\delta}_{\omega^d} \Gamma_0, \sigma_{0\Lambda_0} \tilde{\delta}_{c^d} \rangle, \end{aligned} \quad (296)$$

and recalling notation for $\tilde{\delta}_\phi$ from (24)

$$\begin{aligned} \tilde{\delta}_{A_\sigma^d(q)} \Gamma_0 &= g\epsilon^{adb} (\langle \tilde{\gamma}_\sigma^a \tilde{c}^b; q \rangle + \langle ip_\sigma \tilde{c}^a(p) \tilde{c}^b; q \rangle) + A_{\sigma'}^d(q) (\delta_{\sigma'\sigma} q^2 - q_{\sigma'} q_\sigma) \sigma_{0\Lambda_0}^{-1}(q^2) \\ &\quad + 3\epsilon^{dab} \langle F_{\sigma\mu\nu}^{AAA}(q, \cdot, \cdot) | \check{A}_\mu^a \check{A}_\nu^b; q \rangle + 4R_{\sigma\rho\mu\nu}^{AAAA} \langle \check{A}_\rho^d \check{A}_\mu^a \check{A}_\nu^a; q \rangle, \\ \tilde{\delta}_{\gamma_\sigma^d(q)} \Gamma_0 &= iq_\sigma c^d(q) + g\epsilon^{dab} \langle \check{A}_\sigma^a \tilde{c}^b; q \rangle, \\ \tilde{\delta}_{c^d(q)} \Gamma_0 &= g\epsilon^{dab} \langle ip_\mu \tilde{c}^a(p) \check{A}_\mu^b + \check{\gamma}_\mu^a \check{A}_\mu^b + \tilde{\omega}^b \tilde{c}^a; q \rangle - iq_\mu \gamma_\mu^d(q) - \bar{c}^d(q) q^2 \sigma_{0\Lambda_0}^{-1}(q^2), \\ \tilde{\delta}_{\omega^d(q)} \Gamma_0 &= \frac{1}{2} g\epsilon^{dab} \langle \tilde{c}^a \tilde{c}^b; q \rangle. \end{aligned}$$

Here the notation $\langle ip_\sigma \check{c}^a(p) \check{c}^b; q \rangle$ corresponds to $\langle \phi_1 \phi_2; q \rangle$ with $\phi_1(p) = ip_\sigma \check{c}^a(p)$, $\phi_2 = \check{c}^b$.

For all $\vec{\phi}$, $\vec{\varkappa}$ and w such that $N + 2n_\varkappa + \|w\| = 6$ we have

$$(S_0 F_{1;l})_{\vec{\varkappa}}^{\vec{\phi};w} = (S_0 F_{1,rel;l})_{\vec{\varkappa}}^{\vec{\phi};w} + \sum_{\pi} (-)^{\pi_a} S_{0,\vec{\varkappa}_1}^{\vec{\phi}_1;w_1} \Delta_{\vec{\varkappa}_2;l}^{\vec{\phi}_2;w_2}, \quad (297)$$

where $\Delta_l^{\Lambda\Lambda_0} := F_{1;l}^{\Lambda\Lambda_0} - F_{1,rel;l}^{\Lambda\Lambda_0}$, the sum runs over the permutations $\pi = (\pi_\phi, \pi_\varkappa, \pi_w)$ such that $\vec{\phi}_{\pi_\phi} = \vec{\phi}_1 \oplus \vec{\phi}_2$, $\vec{\varkappa}_{\pi_\varkappa} = \vec{\varkappa}_1 \oplus \vec{\varkappa}_2$, $w_{\pi_w} = w_1 + w_2$, and π_a is the number of transpositions *mod* 2 of anticommuting variables in the permutation π . Using (296), for the terms in the sum on the rhs of (297) we have

$$\begin{aligned} |S_{0,\vec{\varkappa}_1}^{\vec{\phi}_1;w_1} \Delta_{\vec{\varkappa}_2;l}^{\vec{\phi}_2;w_2}| &\leq |\Gamma_{\vec{\varkappa}_1;0}^{A\vec{\phi}_1;w'_1}, \sigma_{0\Lambda_0}^{w''_1} \Delta_{\vec{\varkappa}_2\gamma;l}^{\vec{\phi}_2;w_2}| + |\Gamma_{\vec{\varkappa}_1;0}^{c\vec{\phi}_1;w'_1}, \sigma_{0\Lambda_0}^{w''_1} \Delta_{\vec{\varkappa}_2\omega;l}^{\vec{\phi}_2;w_2}| \\ &\quad + |\Gamma_{\vec{\varkappa}_1\gamma;0}^{\vec{\phi}_1;w'_1}, \sigma_{0\Lambda_0}^{w''_1} \Delta_{\vec{\varkappa}_2;l}^{A\vec{\phi}_2;w_2}| + |\Gamma_{\vec{\varkappa}_1\omega;0}^{\vec{\phi}_1;w'_1}, \sigma_{0\Lambda_0}^{w''_1} \Delta_{\vec{\varkappa}_2;l}^{c\vec{\phi}_2;w_2}|. \end{aligned} \quad (298)$$

Let us show that the lhs of (298) is *small* on \mathbb{M}_n . Using section 5.0.2 and the bounds on irrelevant terms of theorems 3, 4 we see that for all marginal terms

$$\Delta_{\vec{\varkappa};l}^{\vec{\phi};w} \sim 0, \quad \text{and thus} \quad \Delta_{\vec{\varkappa};l}^{\vec{\phi};w} \stackrel{\mathbb{M}_n}{\sim} 0. \quad (299)$$

The relation on the rhs can be obtained by adapting the interpolation in equations (242)-(244). Define $N_i := |\vec{\phi}_i|$ and $n_{\varkappa_i} := |\vec{\varkappa}_i|$. Consider the sum of the first and second term on the rhs. If $2n_{\varkappa_2} + N_2 + \|w_2\| \geq 3$ then the bounds of theorems 1, 3, 4 and (299) imply that the sum is *small*. On the other hand if $2n_{\varkappa_1} + N_1 + \|w_1\| > 3$ then $\|w''_1\| > 0$ and the bounds of theorems 1, 3 also give that the sum is *small*. The analysis of the sum of the third and fourth term on the rhs is similar. If $2n_{\varkappa_2} + N_2 + \|w_2\| \geq 4$ then the bounds of theorems 1, 3, 4 and (299) imply that the sum is *small*. If $2n_{\varkappa_1} + N_1 + \|w_1\| > 2$ then $\|w''_1\| > 0$ and using the bounds of theorems 1, 3 we obtain again that the sum is *small*. It follows that the lhs of (298) is *small*. This fact and (295) imply that $(S_0 F_{1,rel;l})_{\vec{\varkappa}}^{\vec{\phi};w} \stackrel{\mathbb{M}_n}{\sim} 0$ for all marginal terms, which leads to the following equations

$$gu^{\gamma cc} \sim -u^{\gamma Acc}, \quad u_1^{cAA} + u_2^{cAA} \sim \frac{gu^{cA}}{2}, \quad (300)$$

$$u_1^{cAAA} \sim u_2^{cAAA}, \quad u_3^{cAAA} \sim u_4^{cAAA} \sim u_5^{cAAA}. \quad (301)$$

5.0.4 Existence of a constant solution

By our convention (which is the standard one) the renormalization constants that are solutions of the relations listed in RC3 are supposed not to depend on Λ_0 . We give here a proof of this property, which is not evident because these relations contain nontrivial functions of Λ_0 , here denoted by $\zeta_{\Phi}^{\Lambda_0}$. The relations corresponding to the marginal terms $\Gamma_{1\vec{z}}^{\vec{\phi};w}$, $\Gamma_{\beta\vec{z}}^{\vec{\phi};w}$ have respectively the general form

$$c^{\Phi} + \varrho C_1^{\Phi} \varrho + \zeta_{\Phi}^{\Lambda_0} \sim 0, \quad \zeta_{\Phi}^{\Lambda_0} := \varrho C_2^{\Phi} \zeta^{\Lambda_0} + \zeta^{\Lambda_0} C_3^{\Phi} \zeta^{\Lambda_0}, \quad (302)$$

$$c^{\Phi} + V_1^{\Phi} \varrho + \zeta_{\Phi}^{\Lambda_0} \stackrel{\beta}{\sim} 0, \quad \zeta_{\Phi}^{\Lambda_0} := V_2^{\Phi} \zeta^{\Lambda_0}. \quad (303)$$

Here $\varrho = (r_i^{\phi\cdots}, R_i, \Sigma_L^{AA}, \Sigma_T^{AA}, \Sigma^{\bar{c}c})$ denotes the relevant terms for vertex functions, see (511) for the list of $r_i^{\phi\cdots}$ and appendixes E, F for the remaining terms. The sequence ζ^{Λ_0} stands for the irrelevant terms listed in appendixes E, F and for the derivative of $r_i^{\phi\cdots}$, R_i wrt scalar products of momenta. Finally, c^{Φ} is a constant, $V_{1,2}^{\Phi}$ are constant vectors, and $C_{1,2,3}^{\Phi}$ are constant matrices.

At loop order l , the terms $\zeta_{\Phi,l}^{\Lambda_0}$ depend only on $\varrho_{l'}$ of loop order $l' < l$: this property holds because each $\zeta_{\Phi,l}^{\Lambda_0}$ is at least linear in the ζ^{Λ_0} and because all the $\zeta_{l=0}^{\Lambda_0}$ vanish. Moreover, at order l for each relation we have a distinct renormalization constant. Consequently, the aforementioned relations have a solution. The existence of a solution ϱ_l independent of Λ_0 follows immediately if the limit $\lim_{\Lambda_0 \rightarrow \infty} \zeta_{\Phi,l}^{\Lambda_0}$ exists: in this case it is enough to choose a solution of the following equations

$$c^{\Phi} + \varrho C_1^{\Phi} \varrho + \zeta_{\Phi}^{\infty} = 0, \quad c^{\Phi} + V_1^{\Phi} \varrho + \zeta_{\Phi}^{\infty} = 0. \quad (304)$$

The convergence of $\zeta_{\Phi,l}^{\Lambda_0}$ relies on the validity of the bounds of theorem 2 up to order l for all irrelevant terms $\Gamma_{\vec{z}}^{\vec{\phi};w}$ and up to order $l - 1$ for all the relevant ones. This property holds because in our inductive scheme at fixed loop order the irrelevant terms are treated before the relevant ones.

5.1 Γ_{β}^c

The renormalization point is $\vec{q} = (-\bar{q}, \bar{q}) \in \mathbb{M}_2^s$, see (132).

$$\Gamma_{\beta^b}^{c^a}(p) \stackrel{\beta}{\sim} \sigma_{0\Lambda_0}(p^2) \Gamma^{c^a \bar{c}^b}(p) + ip_{\mu} \Gamma_{\gamma_{\mu}^b}^{c^a}(p) \stackrel{\beta}{\sim} -\delta^{ab} f(p^2), \quad (305)$$

$$f(x) := x(1 + \Sigma^{\bar{c}c}(x) - R_1(x)). \quad (306)$$

For the marginal term we obtain

$$-\frac{\delta^{ab}}{3}\Gamma_{\beta^b}^{c^a;p_\mu p_\nu}(p)\stackrel{\beta}{\sim}2\delta_{\mu\nu}f'(p^2)+4p_\mu p_\nu f''(p^2). \quad (307)$$

The coefficient of $\delta_{\mu\nu}$ is *small* at the renormalization point iff

$$f' = 1 - r^{\bar{c}c} - R_1 - \zeta_{\beta^c}^{\Lambda_0} \stackrel{\beta}{\sim} 0, \quad \zeta_{\beta^c}^{\Lambda_0}(p^2) := p^2 \frac{\partial R_1(p^2)}{\partial p^2}. \quad (308)$$

This gives the renormalization condition for R_1 .

5.2 Γ_{β}^{cA}

The renormalization point is $\vec{q} = (\bar{k}, \bar{p}, \bar{q}) \in \mathbb{M}_3^s$, see (132). With $k = -p - q$,

$$\Gamma_{\beta^b}^{c^a A^d_\mu}(p, q) \stackrel{\beta}{\sim} \Gamma^{A^d_\mu c^a \bar{c}^b}(p, k) - ik_\rho \Gamma_{\gamma^b}^{c^a A^d_\mu}(p, q) \quad (309)$$

$$= i\epsilon^{dab} \mathcal{I}_\mu^{0\Lambda_0}(p, q), \quad (310)$$

$$\mathcal{I}_\mu^{0\Lambda_0}(p, q) = k_\mu R_1^{A\bar{c}c}(k, p) + p_\mu r_2^{A\bar{c}c}(k, p) - gk_\rho F_{\rho\mu}^{\gamma Ac}(q, p). \quad (311)$$

Let $\Delta_{\mu\nu}^{\Lambda\Lambda_0} := \mathcal{I}_\mu^{\Lambda\Lambda_0;p_\nu} - \mathcal{I}_\mu^{\Lambda\Lambda_0;q_\nu}$. At zero external momenta and $\Lambda = M$ we have $\Delta_{\mu\nu}^{M\Lambda_0}(0) = 0$. Then using the bounds of theorem 3 we get

$$|\Delta_{\mu\nu}^{M\Lambda_0}(\vec{q})| \leq \int_0^1 dt M |\partial \Gamma_{\beta}^{M\Lambda_0;cA}| \stackrel{\beta}{\sim} 0. \quad (312)$$

The term $\Delta_{\mu\nu}^{\Lambda\Lambda_0}$ obeys the FE, see (129). It remains to integrate the FE from 0 to Λ and use inequality (410) to obtain

$$\Delta_{\mu\nu}^{0\Lambda_0}(\vec{q}) - \Delta_{\mu\nu}^{M\Lambda_0}(\vec{q}) \stackrel{\beta}{\sim} 0, \quad \text{and thus} \quad \Delta_{\mu\nu}^{0\Lambda_0}(\vec{q}) \stackrel{\beta}{\sim} 0. \quad (313)$$

Hence in the monomial basis $\{\delta^s Q^k\}_2$ with $Q = (\bar{p}, \bar{q})$ the δ -component of $\mathcal{I}_\mu^{A\bar{c}c;q_\nu}$ is *small* at the renormalization point if the following condition holds

$$R_1^{A\bar{c}c} - gR_2 - \zeta_{\beta^c A}^{\Lambda_0} \stackrel{\beta}{\sim} 0. \quad (314)$$

This gives the renormalization condition for $R_1^{A\bar{c}c}$.

5.3 Γ_β^{cAA} and $\Gamma_\beta^{cc\bar{c}}$

The renormalization point is $\vec{q} = (\bar{k}, \bar{l}, \bar{p}, \bar{q}) \in \mathbb{M}_4^s$, see (132).

$$\Gamma_{\beta^b}^{c^a A_\mu^t A_\nu^s}(l, p, q) \stackrel{\beta}{\sim} \Gamma^{c^a \bar{c}^b A_\mu^t A_\nu^s}(k, p, q) - ik_\rho \Gamma_{\gamma_\rho^b}^{c^a A_\mu^t A_\nu^s}(l, p, q) \quad (315)$$

$$\Gamma_{\beta^b}^{c^a c^t \bar{c}^s}(l, p, q) \stackrel{\beta}{\sim} \Gamma^{c^a \bar{c}^b c^t \bar{c}^s}(k, p, q) - ik_\rho \Gamma_{\gamma_\rho^b}^{c^a c^t \bar{c}^s}(l, p, q) \quad (316)$$

At $\Lambda = M$ it follows from property (148) that these terms vanish at zero momenta. Denoting the renormalization point by \vec{q} , using the bounds of theorem 3 and integrating the FE from M to 0 we obtain

$$|\Gamma_\beta^{M\Lambda_0; \vec{\phi}}(\vec{q})| \leq \int_0^1 dt |\vec{q}| |\partial \Gamma_\beta^{M\Lambda_0; \vec{\phi}}(t\vec{q})| \stackrel{\beta}{\sim} 0, \quad (317)$$

$$|\Gamma_\beta^{0\Lambda_0; \vec{\phi}}(\vec{q}) - \Gamma_\beta^{M\Lambda_0; \vec{\phi}}(\vec{q})| \leq \frac{M}{\Lambda_0} \mathcal{P}_{2(l-1)}^{(1)}(\log_+ \frac{\Lambda_0}{M}) \stackrel{\beta}{\sim} 0. \quad (318)$$

5.4 $\Gamma_1^{c\bar{c}c\bar{c}c}$ and Γ_1^{cAAAA}

These functions do not have nonvanishing marginal terms:

$$\epsilon^{dab} \langle \check{c}^d \check{A}_\mu^a \check{A}_\nu^b \check{A}_\mu^s \check{A}_\nu^s \rangle = 0, \quad \epsilon^{dab} \langle \check{c}^d \check{c}^s \check{c}^a \check{c}^s \check{c}^b \rangle = 0. \quad (319)$$

5.5 $\Gamma_1^{\bar{c}ccAA}$

From equation (148) it follows that for $\Lambda = M$ the function vanishes if the antighost momentum is zero. Using the bounds of theorem 4 first we obtain at the renormalization point $|\Gamma_1^{M\Lambda_0; \bar{c}ccAA}(\vec{q})| \sim 0$ where $\vec{q} \in \mathbb{M}_5^s$, and then integrating the FE from M to 0 we show that the term is *small* at $\Lambda = 0$.

5.6 Γ_1^{cA}

The renormalization point is $\vec{q} = (-\bar{q}, \bar{q}) \in \mathbb{M}_2^s$, see (132).

$$\Gamma_1^{c^a A^b}(p) \sim i\delta^{ab} F_{T; \mu\nu}^{AA}(p) R_1 p_\nu = i\delta^{ab} p_\mu f(p^2), \quad (320)$$

$$f(x) := \frac{1}{\xi} x R_1(x) \Sigma_L^{AA}(x). \quad (321)$$

The marginal term satisfies

$$\Gamma_1^{c^a A^b; ppp}(p) \sim i\delta^{ab} \left(2f'(p^2)(\Sigma_{t \in \{\delta^2\}_4} t) + 4f''(p^2)(\Sigma_{t \in \{\delta pp\}_4} t) + 8ppppf'''(p^2) \right).$$

For the coefficient of δ -tensors we have

$$\xi f'(p^2) = R_1(p^2) \left(\Sigma_L^{AA}(p^2) + p^2 \frac{\partial \Sigma_L^{AA}(p^2)}{\partial p^2} \right) + p^2 \Sigma_L^{AA}(p^2) \frac{\partial R_1(p^2)}{\partial p^2}. \quad (322)$$

Recalling the definition of $r_{1,2}^{AA}$ in appendix E,

$$r_2^{AA}(p^2) + r_1^{AA}(p^2) + p^2 \frac{\partial r_2^{AA}(p^2)}{\partial p^2} = \frac{1}{\xi} \left(\Sigma_L^{AA}(p^2) + p^2 \frac{\partial \Sigma_L^{AA}(p^2)}{\partial p^2} \right). \quad (323)$$

We then obtain the following sufficient condition

$$u^{cA} \sim 0 \iff R_1(r_2^{AA} + r_1^{AA}) + \zeta_{cA}^{\Lambda_0} \sim 0, \quad (324)$$

where

$$\zeta_{cA}^{\Lambda_0}(p^2) := p^2 \left(R_1(p^2) \frac{\partial r_2^{AA}(p^2)}{\partial p^2} + \frac{1}{\xi} \Sigma_L^{AA}(p^2) \frac{\partial R_1(p^2)}{\partial p^2} \right). \quad (325)$$

See section 5.0.2 for the definition of u^{cA} . Relation (324) gives us the renormalization condition for r_2^{AA} .

5.7 Γ_1^{cAA}

The renormalization point is $\vec{q} = (\bar{k}, \bar{p}, \bar{q}) \in \mathbb{M}_3^s$, see (132). With $k = -p - q$,

$$\Gamma_1^{c^s A_\mu^a A_\nu^b}(p, q) \sim \Gamma^{A_\rho^t A_\mu^a A_\nu^b}(p, q) \Gamma_{\gamma_\rho^t}^{c^s}(k) + \sum_{Z_2} F_{T; \mu\rho}^{AA}(p) \Gamma_{\gamma_\rho^a}^{c^s A_\nu^b}(k, q), \quad (326)$$

where the sum \sum_{Z_2} runs over all cyclic permutations of $\{(\mu, p, a), (\nu, q, b)\}$. The marginal terms are: $\Gamma_1^{cAA; pp}$, $\Gamma_1^{cAA; pq}$, $\Gamma_1^{cAA; qq}$. Using equation (300) we see that $u_2^{cAA} \sim 0 \implies u_1^{cAA} \sim 0$. Acting with $\partial_p \partial_p$ on both sides of (326) we obtain

$$u_2^{cAA} \sim 0 \iff gR_2(1 + r_1^{AA}) - 2R_1 R^{AAA} + \zeta_{cAA}^{\Lambda_0} \sim 0. \quad (327)$$

This gives the renormalization condition for R_2 .

5.8 $\Gamma_{1\gamma}^{\bar{c}ccA}$ and $\Gamma_{1\gamma}^{ccA}$

The renormalization point is $\vec{q} = (\bar{l}, \bar{k}, \bar{q}, \bar{p}) \in \mathbb{M}_4^s$, see (132). With $l = -k - q - p$,

$$\begin{aligned} \Gamma_1^{\bar{c}^a c^b c^d A_\mu^s}(k, q, p) &\sim \Gamma^{c^d \bar{c}^a A_\mu^s A_\rho^t}(l, p, k) \Gamma_{\gamma_\rho^t}^{c^b}(k) - \Gamma^{c^b \bar{c}^a A_\mu^s A_\rho^t}(l, p, q) \Gamma_{\gamma_\rho^t}^{c^d}(q) \\ &\quad + \Gamma^{A_\rho^t c^d \bar{c}^a}(q, l) \Gamma_{\gamma_\rho^t}^{A_\mu^s c^b}(p, k) - \Gamma^{A_\rho^t c^b \bar{c}^a}(k, l) \Gamma_{\gamma_\rho^t}^{A_\mu^s c^d}(p, q) \\ &\quad + F_{T;\mu\rho}^{AA}(p) \Gamma_{\gamma_\rho^s}^{c^b c^d \bar{c}^a}(k, q, l) + \Gamma^{c^t \bar{c}^a A_\mu^s}(l, p) \Gamma_{\omega^t}^{c^b c^d}(k, q) \\ &\quad + \Gamma^{c^t \bar{c}^a}(l) \sigma_{0\Lambda_0}(l^2) \Gamma_{\omega^t}^{c^b c^d A_\mu^s}(k, q, p). \end{aligned} \quad (328)$$

From equation (293) it follows that $u_4^{\bar{c}ccA} \sim 0$ and

$$u_1^{\bar{c}ccA} \sim 0 \implies u^{\gamma Acc} \sim 0, u_2^{\bar{c}ccA} \sim 0, u_3^{\bar{c}ccA} \sim 0. \quad (329)$$

Consequently, we need only one condition

$$u_1^{\bar{c}ccA} \sim 0 \iff g(R_2 - R_3)R_1^{A\bar{c}c} + \zeta_{\bar{c}ccA}^{\Lambda_0} \sim 0. \quad (330)$$

This gives the renormalization condition for R_3 .

5.9 $\Gamma_{1\gamma}^{\bar{c}cc}$ and $\Gamma_{1\gamma}^{cc}$

From equations (293), (300) we have

$$u_1^{\bar{c}ccA} \sim 0 \implies u^{\gamma Acc} \sim 0 \implies u^{\gamma cc} \sim 0 \implies u_{i \in \{1,2\}}^{\bar{c}cc} \sim 0. \quad (331)$$

Consequently, the marginal contribution to the functions is *small*.

5.10 Γ_1^{cAAA}

The renormalization point is $\vec{q} = (\bar{l}, \bar{k}, \bar{q}, \bar{p}) \in \mathbb{M}_4^{cp}$, see (133). With $l = -k - q - p$,

$$\begin{aligned} \Gamma_1^{c^s A_\mu^t A_\nu^b A_\rho^d}(p, q, k) &\sim \Gamma^{A_\sigma^a A_\mu^t A_\nu^b A_\rho^d}(p, q, k) \Gamma_{\gamma_\sigma^s}^{c^s}(l) \\ &\quad + \sum_{Z_3} \Gamma^{A_\sigma^a A_\mu^t A_\nu^b}(p, q) \Gamma_{\gamma_\sigma^a}^{A_\rho^d c^s}(k, l) \\ &\quad + \sum_{Z_3} F_{T;\rho\alpha}^{AA}(k) \Gamma_{\gamma_\alpha^d}^{c^s A_\mu^t A_\nu^b}(l, p, q), \end{aligned} \quad (332)$$

where F_T^{AA} is defined in (491), and the sum \sum_{Z_3} runs over all cyclic permutations of $\{(d, \rho, k), (b, \nu, q), (t, \mu, p)\}$. From (301) it follows that we need two equations

$$u_1^{cAAA} \sim 0 \quad \Longleftrightarrow \quad 8R_1R_2^{AAAA} - 4gR_2R^{AAA} + \zeta_{cAAA,1}^{\Lambda_0} \sim 0, \quad (333)$$

$$u_3^{cAAA} \sim 0 \quad \Longleftrightarrow \quad 4R_1R_1^{AAAA} + 2gR_2R^{AAA} + \zeta_{cAAA,3}^{\Lambda_0} \sim 0. \quad (334)$$

These equations give the renormalization conditions for $R_{1,2}^{AAAA}$.

6 Acknowledgements

We want to thank the reviewer for careful reading of the paper, useful remarks and criticism to which we paid serious attention.

A.N. Efremov is indebted to the Institute for Theoretical Physics of the University of Leipzig whose facilities and financial support were used during the work on the final version of the manuscript.

We are especially grateful to S. Hollands for the time he devoted to study this work, numerous discussions and warm hospitality.

We also thank H. Gies from the University of Jena for his interest in our paper, valuable remarks and questions.

Appendices

A Properties of Gaussian measures

In the following $d\nu = d\nu_C(A)d\nu_S(c, \bar{c})$ is the measure given in (31), (32)

$$d\nu_{C_1+C_2}(A)f(A) = d\nu_{C_1}(A_1)d\nu_{C_2}(A_2)f(A_1 + A_2), \quad (335)$$

$$d\nu_{tC}(A)f(A) = d\nu_C(A)f(t^{\frac{1}{2}}A), \quad (336)$$

$$d\nu_C(A - \delta A) = d\nu_C(A)e^{-\frac{1}{2\hbar}\langle \delta A, C_{\Lambda\Lambda_0}^{-1} \delta A \rangle} e^{\frac{1}{\hbar}\langle A, C_{\Lambda\Lambda_0}^{-1} \delta A \rangle}, \quad (337)$$

$$\frac{d}{d\Lambda}d\nu_C f(A) = \frac{1}{2}d\nu_C \left\langle \frac{\delta}{\delta A}, \hbar \dot{C}^{\Lambda\Lambda_0} \frac{\delta}{\delta A} \right\rangle f(A), \quad (338)$$

$$d\nu_C G(A) \left(AC_{\Lambda\Lambda_0}^{-1} - \hbar \frac{\delta}{\delta A} \right) F(A) = d\nu_C \left(\hbar \frac{\delta}{\delta A} G(A) \right) F(A). \quad (339)$$

When integrating over Grassmann variables one obtains

$$d\nu_{S_1+S_2}(c, \bar{c})f(\bar{c}, c) = d\nu_{S_1}(c_1, \bar{c}_1)d\nu_{S_2}(c_2, \bar{c}_2)f(\bar{c}_1 + \bar{c}_2, c_1 + c_2), \quad (340)$$

$$d\nu_{tS}(c, \bar{c})f(\bar{c}, c) = d\nu_S(c, \bar{c})f(t^{\frac{1}{2}}\bar{c}, t^{\frac{1}{2}}c), \quad (341)$$

$$d\nu_S(c - \delta c, \bar{c} - \delta\bar{c}) = d\nu(c, \bar{c})e^{\frac{1}{\hbar}\langle\delta\bar{c}, S_{\Lambda\Lambda_0}^{-1}\delta c\rangle}e^{-\frac{1}{\hbar}(\langle\bar{c}, S_{\Lambda\Lambda_0}^{-1}\delta c\rangle + \langle\delta\bar{c}, S_{\Lambda\Lambda_0}^{-1}c\rangle)}, \quad (342)$$

$$\frac{d}{d\Lambda}d\nu_S f(\bar{c}, c) = d\nu_S \left\langle \frac{\delta}{\delta c}, \hbar \dot{S}^{\Lambda\Lambda_0} \frac{\delta}{\delta \bar{c}} \right\rangle f(\bar{c}, c), \quad (343)$$

$$d\nu_S \tilde{G}(\bar{c}, c) \left(-\bar{c} S_{\Lambda\Lambda_0}^{-1} + \hbar \frac{\delta}{\delta c} \right) \tilde{F}(\bar{c}, c) = d\nu_S \left(\hbar \frac{\delta_R}{\delta c} \tilde{G}(\bar{c}, c) \right) \tilde{F}(\bar{c}, c), \quad (344)$$

$$d\nu_S \tilde{G}(\bar{c}, c) \left(S_{\Lambda\Lambda_0}^{-1} c + \hbar \frac{\delta}{\delta \bar{c}} \right) \tilde{F}(\bar{c}, c) = d\nu_S \left(\hbar \frac{\delta_R}{\delta \bar{c}} \tilde{G}(\bar{c}, c) \right) \tilde{F}(\bar{c}, c), \quad (345)$$

where right functional derivatives are distinguished from left ones by the label R . Properties (339), (344), (345) are proved for

$$G = e^{\frac{i}{\hbar}\langle j, A \rangle}, \quad F = e^{\frac{i}{\hbar}\langle j', A \rangle}, \quad \tilde{G} = e^{\frac{i}{\hbar}(\langle \bar{c}, \eta \rangle + \langle \bar{\eta}, c \rangle)}, \quad \tilde{F} = e^{\frac{i}{\hbar}(\langle \bar{c}, \eta' \rangle + \langle \bar{\eta}', c \rangle)}, \quad (346)$$

and extended to polynomials in the fields by functional differentiation.

B Chains of vertex functions

For the purpose of example we give the complete list of reduced chains which appear in the loop integrals for $\dot{\Gamma}^{AAc\bar{c};w}$, together with the corresponding "dotted" propagators. The external fields are underlined>. Moreover, $\sum_{i=0}^{k-1} w_i = w$, k being the number of vertex functions in each chain.

$$\begin{aligned}
& \dot{C}\Gamma^{AAA;w_0} C\Gamma^{AAA;w_1} C\Gamma^{Ac\bar{c};w_2} S\Gamma^{c\bar{c}A;w_3}, & \dot{C}\Gamma^{AAAAc\bar{c};w}, \\
& \dot{S}\Gamma^{c\bar{c};w_0} S\Gamma^{c\bar{c};w_1} S\Gamma^{Ac\bar{c};w_2} C\Gamma^{c\bar{c}A;w_3}, & \dot{S}\Gamma^{c\bar{c}AAc\bar{c};w}, \\
& \dot{C}\Gamma^{AAA;w_0} C\Gamma^{Ac\bar{c};w_1} S\Gamma^{c\bar{c}A;w_2}, & \dot{C}\Gamma^{AAAc\bar{c};w_0} C\Gamma^{AAA;w_1}, \\
& \dot{S}\Gamma^{c\bar{c}A;w_0} S\Gamma^{Ac\bar{c};w_1} C\Gamma^{c\bar{c}A;w_2}, & \dot{S}\Gamma^{Ac\bar{c}c\bar{c};w_0} S\Gamma^{Ac\bar{c};w_1}, \\
& \dot{C}\Gamma^{AAc\bar{c};w_0} S\Gamma^{Ac\bar{c};w_1} S\Gamma^{Ac\bar{c};w_2}, & \dot{C}\Gamma^{AAAc\bar{c};w_0} S\Gamma^{Ac\bar{c};w_1}, \\
& \dot{S}\Gamma^{c\bar{c}A;w_0} C\Gamma^{AAA;w_1} C\Gamma^{Ac\bar{c};w_2}, & \dot{S}\Gamma^{AAAc\bar{c};w_0} C\Gamma^{Ac\bar{c};w_1}, \\
& \dot{C}\Gamma^{AAc\bar{c};w_0} S\Gamma^{Ac\bar{c};w_1} C\Gamma^{AAA;w_2}, & \dot{C}\Gamma^{AAAc\bar{c};w_0} S\Gamma^{Ac\bar{c};w_1}, \\
& \dot{S}\Gamma^{c\bar{c}A;w_0} C\Gamma^{Ac\bar{c};w_1} S\Gamma^{Ac\bar{c};w_2}, & \dot{S}\Gamma^{AAAc\bar{c};w_0} C\Gamma^{Ac\bar{c};w_1}, \\
& \dot{C}\Gamma^{Ac\bar{c}A;w_0} C\Gamma^{AAA;w_1} C\Gamma^{AAA;w_2}, & \dot{C}\Gamma^{AAAc\bar{c};w_0} S\Gamma^{AAc\bar{c};w_1}, \\
& \dot{S}\Gamma^{c\bar{c}c\bar{c};w_0} S\Gamma^{Ac\bar{c};w_1} S\Gamma^{Ac\bar{c};w_2}, & \dot{S}\Gamma^{AAc\bar{c};w_0} C\Gamma^{AAc\bar{c};w_1}, \\
& \dot{C}\Gamma^{AAA;w_0} C\Gamma^{Ac\bar{c};w_1}, & \dot{S}\Gamma^{AAc\bar{c};w_0} S\Gamma^{c\bar{c}c\bar{c};w_1}.
\end{aligned} \quad (347)$$

C Tensors

For the definition of the tensor monomial sets $\{\delta^s q^n\}$, $\{\delta^s q^n\}_r$ see beginning of page 23.

Lemma 28 *Let $q = (q_1, \dots, q_m)$ where $q_i \in \mathbb{R}^D$ are $m \in \mathbb{N}$ linearly independent vectors. Then the tensor monomials $\{\delta^s q^n\}_r$ of positive rank $r = 2s + n \leq 2(D - m) + 1$ are linearly independent,*

$$\sum_{t \in \{\delta^s q^n\}_r} c_t t = 0 \implies c_t = 0, \forall t. \quad (348)$$

Proof Observe that, for $m, r, D \in \mathbb{N}$, the inequality $r \leq 2(D - m) + 1$ is equivalent to $m + s \leq D$ for all $s, n \in \mathbb{N}$ such that $r = n + 2s$. Let $I := \{1, \dots, r\}$. Let \mathcal{P}_s be the set of all divisions of the set I in $m + s$ pairwise-disjoint, possibly-empty sets,

$$I = \left(\cup_{j=1}^m V_j \right) \cup \left(\cup_{k=1}^s S_k \right), \quad (349)$$

such that $S_k = \{s_k^1, s_k^2\}$, $s_k^1 < s_k^2$, and $\min S_1 < \dots < \min S_s$. There is a bijection that maps a division $(V_j, S_k) \in \mathcal{P}_s$ to a tensor monomial $t \in \{\delta^s q^n\}$, constructed by the relation

$$t_{\mu_1 \dots \mu_r} = \prod_{j=1}^m \prod_{v_j \in V_j} q_{j; \mu_{v_j}} \prod_{k=1}^s \delta_{\mu_{s_k^1} \mu_{s_k^2}}. \quad (350)$$

Let us first prove the statement of the lemma for orthonormal q_j . In an appropriate basis of \mathbb{R}^D , their components are

$$q_{j; \mu} = \delta_{j\mu}, \quad j \in \{1, \dots, m\}, \quad \mu \in \{1, \dots, D\}. \quad (351)$$

Let us assume that $\sum_t c_t t = 0$, with $t \in \{\delta^s q^n\}_r$. We will proceed by proving inductively that $c_t = 0$ for all $t \in \{\delta^s q^n\}$, from $s = D - m$ down to $s = 0$. Fix $\bar{s} \leq D - m$ and assume that $c_t = 0$ for all t involving more than \bar{s} Kronecker's tensors (which is vacuously true for $\bar{s} = D - m$, due to the rank constraint). Let us prove that $c_{\bar{t}} = 0$ for an arbitrary $\bar{t} \in \{\delta^{\bar{s}} q^{\bar{n}}\}$, which is associated to a division $(\bar{V}_j, \bar{S}_k) \in \mathcal{P}_{\bar{s}}$. Fix the values of the indices $\bar{\mu}_i$, with $i \in I$, by

$$\bar{\mu}_i = \begin{cases} j, & \text{if } \exists j \text{ such that } i \in \bar{V}_j, \\ m + k, & \text{if } \exists k \text{ such that } i \in \bar{S}_k, \end{cases} \quad (352)$$

Note that this choice is possible because $m + \bar{s} \leq D$. It is enough to show that whenever $t_{\bar{\mu}_1 \dots \bar{\mu}_r} \neq 0$ for $t \in \{\delta^s q^n\}$ and $s \leq \bar{s}$ (i.e. $n \geq \bar{n}$) then $s = \bar{s}$ and $t = \bar{t}$: in fact this property, the inductive hypothesis, and the vanishing of the sum $\sum_t c_t t$ imply that $c_{\bar{t}} = 0$. To prove the aforementioned property, introduce the division $(V_j, S_k) \in \mathcal{P}_s$ defining the tensor t and, using (351), correspondingly write

$$0 \neq t_{\bar{\mu}_1 \dots \bar{\mu}_r} = \prod_{j=1}^m \prod_{v_j \in V_j} \delta_{j \bar{\mu}_{v_j}} \prod_{k=1}^s \delta_{\bar{\mu}_{s_1} \bar{\mu}_{s_2} \dots \bar{\mu}_{s_k}}. \quad (353)$$

Relations (352) and (353) imply that $V_j \subseteq \bar{V}_j$ for all j , which, together with the inductive condition $n \geq \bar{n}$, leads to $n = \bar{n}$ and, because the rank r is fixed, to $s = \bar{s}$. Relations (352), (353), and $s = \bar{s}$ imply that there is an injective map $f : I \rightarrow I$ such that $S_j = \bar{S}_{f(j)}$. By definition of the S_j , it then follows that $\min \bar{S}_{f(1)} < \dots < \min \bar{S}_{f(1)}$: this is only possible if f is the identity, which concludes the first part of the proof.

Let us now prove the statement for m linearly independent vectors p_1, \dots, p_m . The sum $\sum_t c_t t = 0$, with $t \in \{\delta^s p^n\}_r$, may be rewritten as

$$\sum_{2s+n=r} \sum_{1 \leq k_1, \dots, k_n \leq m} \sum_{\pi} c_{k_1, \dots, k_n; \pi} \prod_{j=1}^n (p_{k_j})_{\mu_j^\pi} \prod_{j'=1}^s \delta_{\mu_{m+2j'-1}^\pi \mu_{m+2j'}^\pi} = 0, \quad (354)$$

where $\mu_j^\pi := \mu_{\pi(j)}$ and the sum over π runs over the right coset of permutation groups $S_r \setminus (S_n \times S_s \times S_2)$. Expressing the p_k in terms of m orthonormal vectors $q_{k'}$, $p_k = A_{kk'} q_{k'}$, gives a tensor transformation leading to the coefficients in the $\{\delta^s q^n\}_r$ basis:

$$c'_{k'_1, \dots, k'_n; \pi} = c_{k_1, \dots, k_n; \pi} \prod_{j=1}^n A_{k_j k'_j}. \quad (355)$$

The validity of equation (348) for the q_k implies that $c'_{k'_1, \dots, k'_n; \pi} = 0$, which, by invertibility of the matrix $A \in GL(m, \mathbb{R})$, gives $c_{k_1, \dots, k_n; \pi} = 0$. \blacksquare

Note that for $m = n_\phi - 1$ and $D = 4$, the condition for linear independence of the monomials in $\{\delta^s q^k\}_r$ reads

$$r \leq 2(4 - (n_\phi - 1)) + 1 = 11 - 2n_\phi. \quad (356)$$

Lemma 29 *Let $q = (q_1, \dots, q_m)$ with $q_i \in \mathbb{R}^D$ and m nonnegative integer. The tensor monomials $\{\delta^s q^n\}_r$ of positive rank $r \geq 2(D - m + 1)$ are linearly dependent.*

Proof The signature of a permutation π is denoted by $(-)^{\pi}$.

- The Gram matrix $g := g(t_1, \dots, t_k)$ of k tensors t_i of equal rank is defined by its components: $g_{ij} := (t_i, t_j)$ for all $i, j \in [k]$. The scalar product (2) is $O(D)$ -invariant: $(\underline{R}t, \underline{R}u) = (t, u)$ for all tensors t, u of rank r and all $\underline{R} := R \otimes \dots \otimes R$ (r times) with $R \in O(D)$. As a consequence, $g(\underline{R}t_1, \dots, \underline{R}t_k) = g(t_1, \dots, t_k)$ for k tensors of rank r . It is a well-known fact that tensors of equal rank t_1, \dots, t_k are linearly independent iff their Gram matrix $g(t_1, \dots, t_k)$ is invertible. It follows that t_1, \dots, t_k are linearly independent iff $\underline{R}t_1, \dots, \underline{R}t_k$ are linearly independent for some \underline{R} as above stated.
- We assume that $m > 0$ because whenever $m = 0$ for each $s \geq D + 1$ one has

$$\sum_{\pi \in S_s} (-)^{\pi} \delta_{\mu_1 \nu_{\pi(1)}} \cdots \delta_{\mu_s \nu_{\pi(s)}} = 0. \quad (357)$$

- We assume that q_1, \dots, q_m are linearly independent because otherwise there exist c_i not all vanishing such that $\sum_{i \in [m]} c_i q_i = 0$, which, for every $r > 0$, yields

$$\sum_{i_1, \dots, i_r \in [m]} c_{i_1} \cdots c_{i_r} q_{i_1} \otimes \cdots \otimes q_{i_r} = 0. \quad (358)$$

- If the tensors $\{\delta^s q^n\}_r$ are linearly dependent then the tensors $\{\delta^s q^n\}_{r'}$ are linearly dependent, for every $r' > r$. Proof: linear dependence of the tensors $\{\delta^s q^n\}_r$ yields

$$\sum_{t \in \{\delta^s q^n\}_r} c_t t = 0,$$

where not all c_t vanish. Linear independence of q_1, \dots, q_m implies that $q_1 \neq 0$. The proof is concluded by applying $r' - r$ times the tensorial product $\otimes q_1$, which gives

$$\sum_{t \in \{\delta^s q^n\}_r} c_t t \otimes (q_1)^{\otimes r' - r} = 0.$$

- For q_1, \dots, q_D linearly independent the statement of the lemma follows from the previous fact and a well-known relation involving the Gram matrix $g := g(q_1, \dots, q_D)$:

$$\delta_{\mu\nu} = \sum_{i,j \in [D]} q_{i,\mu} (g^{-1})_{ij} q_{j,\nu}. \quad (359)$$

- It is then enough to prove the statement of the lemma for $r = 2(D - m + 1)$, $0 < m < D$, and q_1, \dots, q_m linearly independent and such that $q_i^\mu = 0$ for all $\mu \in [m + 1 : D]$ and $i \in [m]$. There exist q_{m+1}, \dots, q_D orthonormal vectors such that $q_i^\mu = \delta_i^\mu$ for all $\mu \in [D]$ and $i \in [m + 1 : D]$. Relation (359) then gives

$$\delta_{\mu\nu} = \sum_{i,j \in [m]} q_{i,\mu} (g_{\parallel}^{-1})_{ij} q_{j,\nu} + \delta_{\mu\nu}^\perp, \quad (360)$$

where $g_{\parallel} := g(q_1, \dots, q_m)$ and $\delta_{\mu\nu}^\perp := \sum_{i \in [m+1:D]} q_{i,\mu} q_{i,\nu}$. For $s = D - m + 1$, one has

$$\sum_{\pi \in S_s} (-)^{\pi} \delta_{\mu_{m+1} \nu_{m+\pi(1)}}^\perp \cdots \delta_{\mu_{m+s} \nu_{m+\pi(s)}}^\perp = 0. \quad (361)$$

Combining (360) and (361) concludes the proof of the lemma. ■

The following lemma states a necessary condition for a regular, $O(4)$ -invariant tensor field.

Lemma 30 *Let $f(\underline{y})$ be a regular, $O(4)$ -invariant tensor field of rank r where $\underline{y} := (y_1, \dots, y_m)$ with $y_j \in \mathbb{R}^4$. Assume that the tensor monomials $\{\delta^s \underline{y}^k\}_r$ as well as $\{\delta^s \underline{y}^k\}_{r+1}$ are linearly independent pointwise for all $\underline{y} \in O$ where O is some open set. Then on O we have*

$$f = \sum_{t \in \{\delta^s \underline{y}^k\}_r} f_t t, \quad \partial_j f = \sum_{t' \in \{\delta^s \underline{y}^k\}_{r+1}} f_{j,t'} t'. \quad (362)$$

Furthermore for every $t \in \{\delta^s \underline{y}^{k>0}\}_r$ there exist j and $t' \in \{\delta^s \underline{y}^k\}_{r+1}$ such that $f_t = f_{j,t'}$ on O .

Proof For shortness we consider only the case $m = 2$. We have

$$f = \sum_{t \in \{\delta^s\}_r} u_t t + \sum_{t \in \{\delta^s \underline{y}^k > 0\}_r} \zeta_t t, \quad (363)$$

where u_t, ζ_t are regular functions of the scalar parameters $\mathbb{X} = \{\frac{1}{2}y_1^2, \frac{1}{2}y_2^2, y_1y_2\}$. Apply the operator ∂_j to both sides of (363). The Leibniz rule gives

$$\partial_j(u_t \delta^s) = \sum_{x \in \mathbb{X}} (\partial_x u_t) \delta^s \partial_j x, \quad (364)$$

$$\partial_j(\zeta_t \delta^s \underline{y}^k) = \sum_{x \in \mathbb{X}} (\partial_x \zeta_t) \delta^s \underline{y}^k \partial_j x + \zeta_t \delta^s \partial_j \underline{y}^k, \quad (365)$$

where $\partial_j x \in \{0, y_1, y_2\}$, $\underline{y}^k := y_{i_1} \dots y_{i_k}$ with $i_l \in \{1, 2\}$, and

$$\partial_j \underline{y}^k = \partial_j \prod_{l=1}^k y_{i_l} = \sum_{q=1}^k \delta_{j, i_q} \prod_{\substack{l=1 \\ l \neq q}}^k y_{i_l}. \quad (366)$$

The nonvanishing tensor monomials arising at a given $j \in \{1, 2\}$ from the rhs of (364) and (365) have rank $r + 1$, are linearly independent by assumption, and are pairwise different. Each coefficient ζ_t of the tensor $\delta^s y_{i_1} \dots y_j \dots y_{i_k}$ in the decomposition (363) appears also as the coefficient of $\delta^s y_{i_1} \dots \delta \dots y_{i_k}$ in the decomposition (365) for $\partial_j f$. \blacksquare

Lemma 31 below relies on lemma 28, which shows that, for m linearly independent vectors $\underline{e} = (e_1, \dots, e_m)$, the relation $r + 1 \leq 9 - 2m$ is a sufficient condition for the linear independence of the tensor monomials $\{\delta^s \underline{e}^k\}_r$ and $\{\delta^s \underline{e}^k\}_{r+1}$, see lemma 29 for a necessary condition. The renormalization points in appendix H are chosen to comply with the aforementioned relation. The proof of lemma 31 is in the same spirit as the one of the preceding lemma 30.

Lemma 31 *Let F be a regular, $O(4)$ -invariant tensor field of rank $r \in \{2, 4\}$ on \mathbb{P}_n . Let be given $\vec{q} \in \mathbb{M}_n$ and $m \geq 2$ linearly independent vectors $\underline{e} = (e_1, \dots, e_m)$, such that $\text{span}(\vec{q}) = \text{span}(\underline{e})$. Assume that $r + 1 \leq 9 - 2m$. By lemma 28 there exist unique coefficients F_t such that $F(\vec{q}) = \sum_{t \in \{\delta^s \underline{e}^k\}_r} F_t t$. Furthermore,*

$$|F(\vec{q})| \leq c \max \left((|F_t|)_{t \in \{\delta^s\}_r}, (M |\partial_k F(\vec{q})|)_{k \in [n-1]} \right), \quad (367)$$

$$\left| \sum_{t \in \{\delta^s \underline{e}^k > 0\}_r} F_t t \right| \leq c \max \left((M |\partial_k F(\vec{q})|)_{k \in [n-1]} \right). \quad (368)$$

The bounds hold with the same constant c for all F of equal rank.

Proof The coefficients $(|F_t|)_{t \in \{\delta^s\}_r}$ in the basis $\{\delta^s \underline{e}^k\}_r$ do not depend on the choice of the vectors \underline{e} . Hence it is enough to prove (367) in the case when $e_i e_j = M^2 \delta_{ij}$. For simplicity we assume that $m = 2$, the extension to other m being clear.

By hypothesis, there exists a $(n-1) \times 2$ matrix L such that $q_k = L_{ki} e_i$ and

$$|L| := \sqrt{L_{ki} L_{ki}} = \frac{1}{M} \sqrt{L_{ki} L_{kj} e_i e_j} = \frac{1}{M} \left(\sum_{k=1}^{n-1} q_k^2 \right)^{\frac{1}{2}} \leq (n-1)^{\frac{1}{2}}. \quad (369)$$

Denote by $\mathbb{E} \subset \mathbb{R}^4$ the linear span of the vectors q_0, \dots, q_{n-1} . The matrix L induces a linear map $L : \mathbb{E}^2 \rightarrow \mathbb{R}^{4(n-1)}$, $\underline{y} \mapsto L\underline{y}$, where $\underline{y} = (y_1, y_2)$ and $(L\underline{y})_k = L_{ki} y_i$. We also define an auxiliary function on \mathbb{E}^2 : $f(\underline{y}) := F(L\underline{y})$. Setting $\partial_{y_i} := \partial/\partial y_i$ and $\partial_k := \partial/\partial p_k$, the Cauchy–Schwarz inequality and (369) imply that,

$$\sum_{i=1}^2 |\partial_{y_i} f(\underline{e})|^2 \leq |L|^2 \sum_{k=1}^{n-1} |\partial_k F(\bar{q})|^2 \leq (n-1)^2 \max_k (|\partial_k F(\bar{q})|^2). \quad (370)$$

For all $y_1, y_2 \in \mathbb{E}$, denote as usual by $\{\delta^s \underline{y}^k\}$ the set of all monomials being a tensor product of k vectors in $\underline{y} = (y_1, y_2)$ and of s Kronecker tensors, and by $\{\delta^s \underline{y}^k\}_r$ the union of all $\{\delta^s \underline{y}^k\}$ such that $2s + k = r$. By (348), whenever $r \leq 5$ and y_1, y_2 are linearly independent, the elements of $\{\delta^s \underline{y}^k\}_r$ are linearly independent. In this case, we label the tensor monomials by fixing a family of disjoint sets $A_{k,s}$ and a family of bijections $\alpha \mapsto t_\alpha$ from each $A_{k,s}$ to $\{\delta^s \underline{y}^k\}$. Furthermore, we define the auxiliary sets

$$A^r := \bigcup_{2s+k=r} A_{k,s}, \quad A_+^r := \bigcup_{2s+k=r, k \geq 1} A_{k,s}, \quad A_0^r := A^r \setminus A_+^r. \quad (371)$$

For $r \in \{2, 4\}$ and $\underline{y} = (y_1, y_2)$ in an open neighborhood of \underline{e} (in \mathbb{E}^2) for which y_1, y_2 are linearly independent, we write the following tensor decomposition:

$$f(\underline{y}) = \sum_{\alpha \in A_0^r} u_\alpha t_\alpha + \sum_{\alpha \in A_+^r} \zeta_\alpha t_\alpha, \quad (372)$$

where u_α, ζ_α are regular functions of the scalar parameters $\mathbb{X} = \{\frac{1}{2}y_1^2, \frac{1}{2}y_2^2, y_1 y_2\}$. Evaluating (372) at $\underline{y} = \underline{e}$ and using the general fact that $|t| = 2^s M^k$ for

$t \in \{\delta^s \underline{e}^k\}$ we obtain

$$|f(\underline{e})| \leq \sum_{\alpha \in A_0^r} 2^s |u_\alpha| + \sum_{\alpha \in A_+^r} |\zeta_\alpha| 2^s M^k. \quad (373)$$

We now want to prove the existence of a constant $c_1 > 0$ such that

$$\sum_{\alpha \in A_+^r} |\zeta_\alpha| 2^s M^k \leq c_1 M \left(\sum_{i \in \{1,2\}} |\partial_{y_i} f(\underline{e})|^2 \right)^{1/2}. \quad (374)$$

Apply the operator ∂_{y_i} to both sides of (372). The Leibniz rule gives

$$\partial_{y_i}(u_\alpha \delta^s) = \sum_{x \in \mathbb{X}} (\partial_x u_\alpha) \delta^s \partial_{y_i} x, \quad (375)$$

$$\partial_{y_i}(\zeta_\alpha \delta^s y^k) = \sum_{x \in \mathbb{X}} (\partial_x \zeta_\alpha) \delta^s y^k \partial_{y_i} x + \zeta_\alpha \delta^s \partial_{y_i} y^k, \quad (376)$$

where $\partial_{y_i} x \in \{0, y_1, y_2\}$, $y^k = y_{i_1} \dots y_{i_k}$ with $i_l \in \{1, 2\}$, and

$$\partial_{y_i} y^k = \partial_{y_i} \prod_{l=1}^k y_{i_l} = \sum_{j=1}^k \delta_{i,i_j} \prod_{\substack{l=1 \\ l \neq j}}^k y_{i_l}. \quad (377)$$

Fix $\underline{y} = \underline{e}$. The nonvanishing tensor monomials arising at a given $i \in \{1, 2\}$ from the rhs of (375) and (376) have rank $r+1 \in \{3, 5\}$, are pairwise different, and are a subset of the tensor monomials in A^{r+1} , themselves linearly independent by (348). Denote by $B_i^{r+1} \subset A_+^{r+1}$ the subset labeling the monomials of type $\delta^s \delta_{i,i_j} \prod_{l \neq j} y_{i_l}$ arising from (376) and (377) at the given i . By construction, we can define the maps $\pi_i : B_i^{r+1} \rightarrow A_+^r$ with $i \in \{1, 2\}$ such that:

- i) if $t_{\pi_i(\beta)} \in \{\delta^s \underline{e}^k\}$ then $t_\beta \in \{\delta^{s+1} \underline{e}^{k-1}\}$ (in this case, $|t_{\pi_i(\beta)}| = 2^s M^k$ and $|t_\beta| = 2^{s+1} M^{k-1}$);
- ii) for each $\beta \in B_i^{r+1}$ the coefficient of t_β in (376) and that of $t_{\pi_i(\beta)}$ in (372) are the same, namely $\zeta_{\pi_i(\beta)}$;
- iii) $\pi_1(B_1^{r+1}) \cup \pi_2(B_2^{r+1}) = A_+^r$.

The following bound holds at $\underline{y} = \underline{e}$ for every tensor $\Psi = \sum_{\beta} \Psi_\beta t_\beta / |t_\beta|$ with $\beta \in A^{r'}$ and $r' \leq 5$, and for every nonempty $B \subset A^{r'}$:

$$|\Psi|^2 = (\Psi, \Psi) = \Psi_\beta^* G_{\beta\beta'} \Psi_{\beta'} \geq \lambda_1 \sum_{\beta \in A^{r'}} |\Psi_\beta|^2 \geq \lambda_1 \sum_{\beta \in B} |\Psi_\beta|^2, \quad (378)$$

where $\lambda_1 > 0$ is the smallest eigenvalue of the Gram matrix of components $G_{\beta\beta'} := (t_\beta, t_{\beta'}) / (|t_\beta||t_{\beta'}|)$, which is positive definite by (348). Application of (378) to $\Psi_i := \partial_{y_i} f(\underline{e})$ for each $i \in \{1, 2\}$, with $r' = r + 1$ and $B = B_i^{r+1}$, gives

$$\sum_{i \in \{1, 2\}} |\partial_{y_i} f(\underline{e})|^2 \geq \lambda_1 \sum_{\substack{i \in \{1, 2\} \\ \beta \in B_i^{r+1}}} |\zeta_{\pi_i(\beta)}|^2 |t_\beta|^2 \geq \frac{4\lambda_1}{M^2} \sum_{\alpha \in A_+^r} |\zeta_\alpha|^2 (2^s M^k)^2, \quad (379)$$

from which follows the bound (374).

Inequalities (373), (379) and (370) lead to the bounds (367), (368), with a constant

$$c = \max \left(2^{\frac{r}{2}} |A_0|, (n-1) \sqrt{\frac{|A_+^r|}{4\lambda_1}} \right). \quad (380)$$

■

D Basic estimates

Lemma 32 *Let $0 \leq q \leq p \leq P$, $k \in \mathbb{N}$. Then $\exists C_k > 0$ such that*

$$\int_{\Lambda}^{+\infty} \frac{d\lambda \log_+^k \frac{P}{\lambda}}{(\lambda+p)(\lambda+q)} \leq C_k \frac{1 + \log_+^{k+1} \frac{P}{\Lambda+q}}{\Lambda+p+q}. \quad (381)$$

Proof Let I^k be the left hand side of the inequality and

$$I_{[a,b]}^k := \int_a^b \frac{d\lambda}{(\lambda+p)(\lambda+q)} \log_+^k \frac{P}{\lambda}. \quad (382)$$

We begin with the case $k \geq 1$.

- $\Lambda > P$, $I^k = 0$.
- $q \leq \Lambda \leq P$, $I^k = I_{[\Lambda, P]}^k$.

$$\begin{aligned} I^k &\leq \frac{1}{\Lambda+p} \int_{\Lambda}^P \frac{d\lambda}{\lambda+q} \log^k \frac{2P}{\lambda+q} \leq \frac{2}{k+1} \frac{1}{\Lambda+p+q} \log^{k+1} \frac{2P}{\Lambda+q} \\ &\leq \frac{A_k}{\Lambda+p+q} \left(\log_+^{k+1} \frac{P}{\Lambda+q} + 1 \right), \end{aligned} \quad (383)$$

where

$$A_k := \frac{2(\log 2 + 1)^{k+1}}{k + 1}, \quad (384)$$

and we have used the inequality

$$\frac{\log_+ x + \log 2}{(1 + \log_+^n x)^{\frac{1}{n}}} < 1 + \log 2. \quad (385)$$

- $\Lambda < q \leq P$, $I^k = I_{[\Lambda, q]}^k + I_{[q, P]}^k$.

$$\begin{aligned} I_{[q, P]}^k &\leq \frac{1}{\Lambda + p} \int_q^P \frac{d\lambda}{\lambda + q} \log^k \frac{2P}{\lambda + q} \leq \frac{1}{\Lambda + p} \int_\Lambda^P \frac{d\lambda}{\lambda + q} \log^k \frac{2P}{\lambda + q} \\ &\leq \frac{A_k}{\Lambda + p + q} (\log_+^{k+1} \frac{P}{\Lambda + q} + 1), \end{aligned} \quad (386)$$

$$\begin{aligned} I_{[\Lambda, q]}^k &\leq \frac{1}{(\Lambda + p)(\Lambda + q)} \int_0^q d\lambda \log^k \frac{P}{\lambda} < \frac{qk!}{(\Lambda + p)(\Lambda + q)} \sum_{j=0}^k \frac{1}{j!} \log^j \frac{P}{q} \\ &< \frac{ek!}{\Lambda + p} (\log^k \frac{2P}{\Lambda + q} + 1) < \frac{2ek!}{\Lambda + p + q} (\log_+^k \frac{P}{\Lambda + q} + 2). \end{aligned} \quad (387)$$

It remains to consider the case $k = 0$.

- $\Lambda > p$. This implies $2(\lambda + q) > \lambda + p + 2q > \lambda + p$, $\forall \lambda > \Lambda$.

$$I^0 = 2 \int_\Lambda^{+\infty} \frac{d\lambda}{(\lambda + p)^2} < \frac{4}{\Lambda + p + q}. \quad (388)$$

- $\Lambda \leq p$, $I^0 = I_{[\Lambda, p]}^0 + I_{[p, +\infty]}^0$.

$$I_{[p, +\infty]}^0 < 2 \int_p^{+\infty} \frac{d\lambda}{(\lambda + p)^2} < \frac{1}{p} < \frac{3}{\Lambda + p + q}, \quad (389)$$

$$\begin{aligned} I_{[\Lambda, p]}^0 &< \frac{1}{\Lambda + p} \int_\Lambda^p \frac{d\lambda}{\lambda + q} < \frac{2}{\Lambda + p + q} \log \frac{2P}{\Lambda + q} \\ &< \frac{2}{\Lambda + p + q} (\log 2 + \log_+ \frac{P}{\Lambda + q}). \end{aligned} \quad (390)$$

■

Lemma 33 *Let $0 \leq q \leq p$, $\eta := \min(M, q)$, $k \in \mathbb{N}$. Then $\exists C_k > 0$ such that*

$$\int_{\Lambda}^{+\infty} \frac{d\lambda \log_+^k \frac{\lambda}{M}}{(\lambda + p)(\lambda + q)} \leq C_k \frac{1 + \log_+^{k+1} \frac{p}{\Lambda + \eta} + \log_+^{k+1} \frac{\Lambda}{M}}{\Lambda + p + q}. \quad (391)$$

Proof Denote by I^k the lhs of equation (391). If $k = 0$ then the inequality follows from (381) with $P = p$. Let $k \geq 1$ and $\mu = \max(\Lambda, M)$.

- $p \leq \Lambda$.

$$I^k \leq \int_{\mu}^{+\infty} d\lambda \frac{2 \log_+^k \frac{\lambda}{M}}{(\lambda + p)^2} \leq \int_{\mu}^{+\infty} d\lambda f'(\lambda) \leq 2^{k+2} k! \sqrt{e} \frac{1 + \log_+^k \frac{\Lambda}{M}}{\Lambda + p + q}, \quad (392)$$

$$f(\lambda) := -\frac{k! 2^{k+1}}{\lambda + p} \sum_{j=0}^k \frac{\log_+^k \frac{\lambda}{M}}{j! 2^j}. \quad (393)$$

- $\Lambda < p$.

$$J^k := \int_{\Lambda}^p d\lambda \frac{\log_+^k \frac{\lambda}{M}}{\lambda + q}, \quad I^k \leq \frac{2J^k}{\Lambda + p + q} + \int_{\mu}^{+\infty} d\lambda \frac{2 \log_+^k \frac{\lambda}{M}}{(\lambda + p)^2}. \quad (394)$$

The integral on the rhs of I^k is exactly the same as in the case $p \leq \Lambda$. For the integral J^k using the inequality

$$\log_+ \frac{\lambda}{M} \leq \log 2 + \log_+ \frac{\lambda}{\Lambda + M} + \log_+ \frac{\Lambda}{M}, \quad (395)$$

we have

$$J^k < 3^{k+1} \left(1 + \log_+^{k+1} \frac{p}{\Lambda + \eta} + \log_+^{k+1} \frac{\Lambda}{M} \right). \quad (396)$$

■

Lemma 34 *For $0 < \Lambda < \eta \leq M$, $k \in \mathbb{N}$*

$$\int_{\Lambda}^{\eta} d\lambda \frac{1}{\lambda + \eta} \log_+^k \frac{\lambda}{M} < 1. \quad (397)$$

Proof Denote by I^k the lhs of equation (397). If $k > 0$ then $I^k = 0$. It remains to consider the integral

$$\int_{\Lambda}^{\eta} \frac{d\lambda}{\lambda + \eta} = \log \frac{2\eta}{\Lambda + \eta} < \log 2. \quad (398)$$

■

Lemma 35 Let $0 < \Lambda < \eta \leq P$, $k \in \mathbb{N}$. Then $\exists C_k > 0$ such that

$$\int_{\Lambda}^{\eta} \frac{d\lambda}{\lambda + \eta} \log_+^k \frac{P}{\lambda} < C_k \left(1 + \log_+^k \frac{P}{\Lambda + \eta} \right). \quad (399)$$

Proof Denote by $I_{[\Lambda, \eta]}^k$ the lhs of equation (399).

$$I_{[\Lambda, \eta]}^k < I_{[0, \eta]}^k < \frac{1}{\eta} \int_0^{\eta} d\lambda \log_+^k \frac{P}{\lambda} < 2k!e^2 \left(1 + \log_+^k \frac{P}{\eta + \Lambda} \right). \quad (400)$$

■

Lemma 36 Let $0 < \Lambda < M \leq P$, $0 \leq \eta \leq M$, $k \in \mathbb{N}$. Then $\exists C_k > 0$ such that

$$\int_{\Lambda}^M \frac{d\lambda}{\lambda + \eta} \log_+^k \frac{P}{\lambda} < C_k \left(1 + \log_+^{k+1} \frac{P}{\Lambda + \eta} \right). \quad (401)$$

Proof Denote by $I_{[\Lambda, M]}^k$ the lhs of equation (401). If $\Lambda \geq \eta$ then

$$I_{[\Lambda, M]}^k \leq \int_{\Lambda}^M \frac{d\lambda}{\lambda + \eta} \log_+^k \frac{2P}{\lambda + \eta} \leq \frac{2^{k+1}}{k+1} \left(1 + \log_+^{k+1} \frac{P}{\Lambda + \eta} \right). \quad (402)$$

If $\eta \geq \Lambda$ then $I^k < I_{[0, \eta]}^k + I_{[\eta, M]}^k$ where

$$I_{[0, \eta]}^k = \int_0^{\eta} \frac{d\lambda}{\lambda + \eta} \log_+^k \frac{P}{\lambda} < 2k!e^2 \left(1 + \log_+^k \frac{P}{\eta + \Lambda} \right), \quad (403)$$

$$I_{[\eta, M]}^k = \int_{\eta}^M \frac{d\lambda}{\lambda + \eta} \log_+^k \frac{2P}{\lambda + \eta} < \frac{2^{k+1}}{k+1} \left(1 + \log_+^{k+1} \frac{P}{\Lambda + \eta} \right). \quad (404)$$

Lemma 37 For $q \geq 0$, $k \in \mathbb{N}$, $P > 0$ there exists a constant $C_k > 0$ such that

$$\int_{\Lambda}^{\Lambda_0} \frac{d\lambda}{\lambda + q} \log_+^k \frac{P}{\lambda} < C_k \left(1 + \log_+^{k+1} \frac{P}{\Lambda + q} + \log_+ \frac{\Lambda_0}{\Lambda + q} \right). \quad (405)$$

Proof Denote by $I_{[\Lambda, \Lambda_0]}^k$ the lhs of equation (405). If $k = 0$ then

$$I_{[\Lambda, \Lambda_0]}^0 < 1 + \log_+ \frac{\Lambda_0}{\Lambda + q}. \quad (406)$$

For $k > 0$ and $\Lambda \geq q$

$$I_{[\Lambda, \Lambda_0]}^k < \frac{2^{k+1}}{k+1} \left(1 + \log_+^{k+1} \frac{P}{\Lambda + q} \right). \quad (407)$$

If $k > 0$ and $\Lambda < q$ then $I_{[\Lambda, \Lambda_0]}^k < I_{[0, q]}^k + I_{[q, \Lambda_0]}^k$ where

$$I_{[0, q]}^k < \frac{1}{q} \int_0^q d\lambda \log_+^k \frac{P}{\lambda} < k! e^2 \left(1 + \log_+^{k+1} \frac{P}{\Lambda + q} \right), \quad (408)$$

$$I_{[q, \Lambda_0]}^k < \frac{2^{k+1}}{k+1} \left(1 + \log_+^{k+1} \frac{P}{\Lambda + q} \right). \quad (409)$$

Lemma 38 Let $a, d > 0$, $b \geq 0$ and $m, k \in \mathbb{N}$. Then $\exists C_{k,m} > 0$ such that

$$\int_0^a dx x^m \log_+^k \frac{d}{b+x} \leq C_{k,m} a^{m+1} \left(1 + \log_+^k \frac{d}{a+b} \right). \quad (410)$$

Proof By direct calculations it is easy to show that

$$x^m \log_+^k \frac{d}{b+x} \leq f', \quad f := \frac{k! x^{m+1}}{(m+1)^{k+1}} \sum_{j=0}^k \frac{(m+1)^j}{j!} \log_+^j \frac{d}{b+x}. \quad (411)$$

Consequently, the lhs of equation (410) is bounded above by $f(a)$,

$$\begin{aligned} f(a) &= \frac{k!}{(m+1)^{k+1}} a^{m+1} \sum_{j=0}^k \frac{(m+1)^j}{j!} \log_+^j \frac{d}{a+b} \\ &\leq \frac{e^{m+1} k!}{(m+1)^{k+1}} a^{m+1} \left(1 + \log_+^k \frac{d}{a+b}\right). \end{aligned} \quad (412)$$

■

Lemma 39 *Let $p, q \in \mathbb{R}^4$, $\Lambda' > 0$, $\Lambda' \geq \eta \geq 0$. Then*

$$\int_0^1 dt \frac{|p|}{\Lambda' + |tp + q|} \leq 2 \left(\log 4 + \log_+ \frac{|p|}{\Lambda' + \eta} \right). \quad (413)$$

Proof Let $I_{[0,1]}$ denote the lhs of equation (413). There exists $t_1 \in [0, 1]$ such that $|tp + q| \geq |p||t - t_1|$ for all $t \in [0, 1]$. (To prove this fact: write $q = -t_q p + q^\perp$, with $p \cdot q^\perp = 0$, in such a way that $|tp - q| \geq |t - t_q| |p|$; then set $t_1 := 0$ for $t_q < 0$, $t_1 := 1$ for $t_q > 1$, and $t_1 = t_q$ otherwise.)

$$\begin{aligned} I_{[0,1]} &= I_{[0,t_1]} + I_{[t_1,1]} \leq 2 \log \frac{\Lambda' + |p|}{\Lambda'} \leq 2 \log 4 \max\left(1, \frac{|p|}{\Lambda' + \eta}\right) \\ &\leq 2 \left(\log 4 + \log_+ \frac{|p|}{\Lambda' + \eta} \right). \end{aligned} \quad (414)$$

■

Lemma 40 *Let $r > 0$, $w \in \mathbb{N}$ and $x, y \in \mathbb{R}^4$*

$$\frac{e^{-rx^2}}{(1 + |x - y|)^w} \leq \frac{w! \max\left(2, 1 + \frac{1}{2r}\right)^w}{(1 + |y|)^w}. \quad (415)$$

Proof Choosing the Cartesian coordinate system such that one of the basis vectors e_L is along y we have

$$\frac{e^{-rx^2}}{(1 + |x - y|)^w} \leq f(t), \quad f(t) := \frac{e^{-rt^2}}{(1 + |t - t_0|)^w}. \quad (416)$$

where t, t_0 are the longitudinal coordinates, $x = x_T + te_L$ and $y = t_0e_L$, $t_0 = |y|$. If $t \geq t_0$ then f is strictly decreasing. If $t \leq t_0$ then $f(t) = g(t)$ where

$$g(t) := \frac{e^{-rt^2}}{(t_1 - t)^w}, \quad t_1 := 1 + t_0. \quad (417)$$

If $t < t_1$ then g is either increasing or has a local maximum at t_- .

$$g'(t) \leq 0 \text{ if } (\Delta \geq 0) \wedge (t_- \leq t \leq t_+), \quad t_{\mp} := \frac{t_1 \mp \sqrt{\Delta}}{2}, \quad (418)$$

$$g'(t) \geq 0 \text{ otherwise,} \quad \Delta := t_1^2 - 2\frac{w}{r}. \quad (419)$$

Consequently, $f(t) \leq \max(g(t_-), g(t_0))$ where

$$g(t_-) = \frac{e^{-rt_-^2}}{t_+^w} \leq \frac{2^w}{t_1^w} = \frac{2^w}{(1+t_0)^w}, \quad g(t_0) = e^{-rt_0^2} \leq \frac{w!(1+\frac{1}{2r})^w}{(1+t_0)^w}. \quad (420)$$

■

Lemma 41 *Let $r > 0$. There is a constant C such that*

$$e^{-r\frac{s^2}{\Lambda^2}} \log_+ \frac{\max(M, \sqrt{p^2 + s^2})}{\Lambda} < C + \frac{1}{2} \log_+ \frac{1}{r} + \log_+ \frac{\max(M, p)}{\Lambda}. \quad (421)$$

Proof Using the following inequality

$$\frac{\max(M, \sqrt{p^2 + s^2})}{\Lambda} \leq 2 \max(1, \frac{s}{\Lambda}) \max(1, \frac{\max(M, p)}{\Lambda}), \quad (422)$$

we bound the lhs of the statement by

$$\log_+ \frac{\max(M, p)}{\Lambda} + \log 2 + \frac{1}{2} \log_+ \frac{1}{r} + \frac{1}{2} e^{-z} \log_+ z, \quad z := r \frac{s^2}{\Lambda^2}. \quad (423)$$

The inequality $e^{-z} \log_+ z < e^{-1}$ finishes the proof. ■

Lemma 42 *Let $x, y, m \geq 0$, $\mathcal{P}^{(0)}(x) = \sum_{k=0}^n c_k x^k$ a polynomial of the degree n , $\log_m x := \log_+ \max(x, m)$. Then there exist polynomials $\mathcal{P}^{(1)}, \mathcal{P}^{(2)}$ of the degree n such that*

$$\mathcal{P}^{(0)}(\log_m \sqrt{y^2 + x^2}) \leq \mathcal{P}^{(1)}(\log_m y) + \mathcal{P}^{(2)}(\log_m x). \quad (424)$$

Proof Substitution of the inequalities

$$\max(\sqrt{y^2 + x^2}, m) \leq \max(y + x, m) \leq \max(y, m) + \max(x, m), \quad (425)$$

$$\max(a + b, 1) \leq \max(a, 1) + \max(b, 1) \leq 2 \max(a, 1) \max(b, 1), \quad (426)$$

into the definition $\log_+ a := \log \max(a, 1)$ yields

$$\log_m \sqrt{y^2 + x^2} \leq \log_m y + \log_m x + \log 2. \quad (427)$$

This gives

$$\mathcal{P}^{(0)}(\log_m \sqrt{x^2 + y^2}) \leq \sum_{k=0}^n c_k 3^k (\log_m^k y + \log_m^k x + 1). \quad (428)$$

■

Lemma 43 *For a fixed $s \in \mathbb{N}$ there exists a constant c such that $\forall u \leq w_{max}$ and $\forall x \in \mathbb{R}^4$*

$$|\partial^u (x^{\otimes s} (1 - e^{-x^4}))| \leq c \begin{cases} |x|^{s+1-u}, & u \leq s, \\ 1, & \text{otherwise.} \end{cases} \quad (429)$$

Proof First we consider the case $u \leq s$

$$|\partial^u (x^{\otimes s} (1 - e^{-x^4}))| \leq |\partial^u x^{\otimes s}| (1 - e^{-x^4}) + \sum_{u_2 > 0} \frac{u! 4^{u_2}}{u_1! u_2!} |\partial^{u_1} x^{\otimes s+3u_2}| e^{-x^4}. \quad (430)$$

For the derivatives on the right we have

$$|\partial^{u_1} x^{\otimes s+3u_2}| \leq (s + 3u_2)! \frac{|x|^{s-u+1}}{(s-u+1)!} \frac{|x|^{4u_2-1}}{(4u_2-1)!}. \quad (431)$$

Furthermore, $\forall k \in \mathbb{N}$

$$\frac{|x|^k}{k!} e^{-x^4} < \sum_{n=0}^{\infty} \frac{|x|^n}{n!} e^{-x^4} = e^{-x^4+|x|} < e. \quad (432)$$

Noting that $1 - e^{-x^4} \leq |x|$, we obtain the bound in the case $u \leq s$

$$\frac{s! |x|^{s+1-u}}{(s-u)!} + \frac{e |x|^{s+1-u}}{(s-u+1)!} \sum_{u_2 > 0} \frac{u! 4^{u_2} (s+3u_2)!}{u_1! u_2!} \leq c |x|^{s+1-u}. \quad (433)$$

For $u > s$ using (432) we have

$$\sum \frac{u!4^{u_2}(s+3u_2)!}{u_1!u_2!} \frac{|x|^{s+3u_2-u_1}}{(s+3u_2-u_1)!} e^{-x^4} \leq (s+3u)!5^u e \leq c. \quad (434)$$

■

Lemma 44 *For a fixed $s \in \mathbb{N} \setminus \{0\}$ there exists a constant c such that for all $u, v \leq w_{max}$ and all $x = (x_1, x_2)$ with $x_i \in \mathbb{R}^4$*

$$|\partial_2^v \partial_1^u (x_2 \otimes x_1^{\otimes s-1} (1 - e^{-x_1^4}))| \leq c \begin{cases} |x|^{s+1-u-v}, & u+v \leq s, \\ |x|+1, & \text{otherwise.} \end{cases} \quad (435)$$

Proof First let $v \in \{0, 1\}$. Using lemma 43 we obtain for $0 \leq u \leq s$

$$|\partial_1^u (x_2 \otimes x_1^{\otimes s-1} (1 - e^{-x_1^4}))| \leq c|x|^{s+1-u}, \quad (436)$$

$$|\partial_2 \partial_1^u (x_2 \otimes x_1^{\otimes s-1} (1 - e^{-x_1^4}))| \leq c|x|^{s-u}, \quad (437)$$

and for $u > s$ we also have

$$|\partial_1^u (x_2 \otimes x_1^{\otimes s-1} (1 - e^{-x_1^4}))| \leq c|x|, \quad (438)$$

$$|\partial_2 \partial_1^u (x_2 \otimes x_1^{\otimes s-1} (1 - e^{-x_1^4}))| \leq c. \quad (439)$$

Finally, for $v > 1$ $\partial_2^v \partial_1^u (x_2 \otimes x_1^{\otimes s-1} (1 - e^{-x_1^4})) = 0$. ■

Lemma 45 *For $s \in \{0, 1\}$ there exists a constant c such that for all $u, v \leq w_{max}$ and all $x = (x_1, x_2)$ with $x_i \in \mathbb{R}^4$*

$$|\partial_2^v \partial_1^u (x_1^{\otimes s} (e^{-x_1^4} - e^{-(x_1+x_2)^4}))| \leq c \begin{cases} |x|^{s+1-u-v}, & u+v \leq s, \\ |x|+1, & \text{otherwise.} \end{cases} \quad (440)$$

Proof First let $s = 0$. For $u = v = 0$ put $y = x_1 + x_2$ and assume $|x_1| \leq |y| \leq 1$

$$e^{-x_1^4} - e^{-y^4} = e^{-x_1^4} (1 - e^{-(y^4-x_1^4)}) \leq y^4 - x_1^4 \leq |y| \leq 2|x|. \quad (441)$$

Inequality (432) implies that

$$\partial^w e^{-x^4} < e(3w)!5^4. \quad (442)$$

Consequently for $u + v > 0$

$$|\partial_2^v \partial_1^u (e^{-x_1^4} - e^{-(x_1+x_2)^4})| \leq |\partial_1^u e^{-x_1^4}| + |\partial_2^v \partial_1^u e^{-(x_1+x_2)^4}| \leq c_2. \quad (443)$$

Finally, we consider the case $s = 1$. Using (441) and (442) we get

$$|x_1(e^{-x_1^4} - e^{-(x_1+x_2)^4})| \leq 2|x|^2, \quad (444)$$

$$|\partial_2(x_1(e^{-x_1^4} - e^{-(x_1+x_2)^4}))| \leq |x| |\partial_2(e^{-x_1^4} - e^{-(x_1+x_2)^4})| \leq c_3|x|, \quad (445)$$

$$|\partial_1(x_1(e^{-x_1^4} - e^{-(x_1+x_2)^4}))| \leq c_4|x|, \quad (446)$$

and for $u + v > 1$

$$|\partial_1^u \partial_2^v (x_1(e^{-x_1^4} - e^{-(x_1+x_2)^4}))| \leq c_5 + |x|c_6. \quad (447)$$

■

Lemma 46 *Let $0 \leq \beta \leq 1$, $x \geq 0$ and*

$$f(x) := \frac{1}{x}(e^{-\beta x} - e^{-x}) \quad (448)$$

Then $\forall w \in \mathbb{N}$

$$|\partial^w f(x)| < e \frac{w!}{(1+x)^{w+1}}. \quad (449)$$

Proof We have an identity

$$\partial^w f(x) = (-1)^w g_w(x), \quad g_w(x) := \int_{\beta}^1 d\gamma \gamma^w e^{-\gamma x}, \quad (450)$$

It follows

$$\begin{aligned} 0 \leq g_w(x) &\leq \int_0^1 d\gamma \gamma^w e^{-\gamma(x+1)} e^{\gamma} < e \int_0^1 d\gamma \gamma^w e^{-\gamma(x+1)} \\ &< e \frac{1}{(x+1)^{w+1}} \int_0^{\infty} dz z^w e^{-z} = e \frac{\Gamma(w+1)}{(x+1)^{w+1}}. \end{aligned} \quad (451)$$

■

Lemma 47 Let $0 \leq \beta \leq 1$, $x \geq 0$ and

$$h(x) := \frac{1}{x}(e^{-\beta x^2} - e^{-x^2}) \quad (452)$$

Then $\forall w \in \mathbb{N}$, $\forall C > 1$

$$|\partial^w h(x)| < w!e(C+1) \frac{(2\sqrt{e}C)^w}{(1+x)^{w+1}}. \quad (453)$$

Proof

$$|\partial^w h(x)| \leq |x\partial^w f(x^2)| + |w\partial^{w-1} f(x^2)|, \quad f(x^2) := \frac{1}{x}h(x). \quad (454)$$

Using an auxiliary variable y

$$\begin{aligned} \frac{\partial^w}{(\partial x)^w} f(x^2) &= \left(\frac{\partial}{\partial x} + 2x \frac{\partial}{\partial y} \right)^w f(y)|_{y=x^2} \\ &= \sum_{k=0}^{2k \leq w} \frac{w!(2k-1)!!}{(w-2k)!(2k)!} \left(2 \frac{\partial}{\partial y} \right)^k \left(2x \frac{\partial}{\partial y} \right)^{w-2k} f(y)|_{y=x^2} \\ &= \sum_{k=0}^{2k \leq w} \frac{w!}{(w-2k)!k!} (2x)^{w-2k} \left(\frac{\partial}{\partial y} \right)^{w-k} f(y)|_{y=x^2}. \end{aligned} \quad (455)$$

Equation (449) gives

$$|x\partial^w f(x^2)| < w!2^w e \sum_{k=0}^{2k \leq w} \frac{1}{k!} \frac{(w-k)!}{(w-2k)!} 2^{-2k} \frac{x^{w-2k+1}}{(1+x^2)^{w-k+1}}. \quad (456)$$

With an arbitrary constant $C > 1$ we have

$$\frac{x}{1+x^2} < \frac{C}{1+x}, \quad (1+x^2) \geq \frac{1}{2}(1+x)^2, \quad \frac{(w-k)!}{(w-2k)!} \leq w^k. \quad (457)$$

Consequently,

$$|x\partial^w f(x^2)| < \frac{w!Ce(2C)^w}{(1+x)^{w+1}} \sum_{k=0}^{2k \leq w} \frac{1}{k!} \left(\frac{w}{2C^2} \right)^k < \frac{w!Ce(2\sqrt{e}C)^w}{(1+x)^{w+1}}. \quad (458)$$

Similarly, we obtain

$$|w\partial^{w-1} f(x^2)| < \frac{2w!e(2\sqrt{e}C)^{w-1}}{(1+x)^{w+1}}. \quad (459)$$

■

Lemma 48 *Let $f(p^2)$ be a scalar function. Then*

$$\left| \prod_{i=1}^w \frac{\partial}{\partial p_{\mu_i}} f(p^2) \right| \leq 2^w \sum_{k=0}^{2k \leq w} \frac{w!}{(w-2k)!k!} |p|^{w-2k} \left(\frac{\partial}{\partial p^2} \right)^{w-k} f(p^2). \quad (460)$$

Proof With the aid of an auxiliary variable y

$$\prod_{i=1}^w \frac{\partial}{\partial p_{\mu_i}} f(p^2) = \prod_{i=1}^w \left(\frac{\partial}{\partial p_{\mu_i}} + 2p_{\mu_i} \frac{\partial}{\partial y} \right) f(y)|_{y=p^2}. \quad (461)$$

A partial derivative wrt p_μ contributes only if it can be paired with $2p_\mu$ term. Consequently, we can compute the right hand side by considering the possible pairs,

$$\left(\frac{\partial}{\partial p_{\mu_i}} + 2p_{\mu_i} \frac{\partial}{\partial y} \right) \left(\frac{\partial}{\partial p_{\mu_j}} + 2p_{\mu_j} \frac{\partial}{\partial y} \right) = 2\delta_{\mu_i \mu_j} \frac{\partial}{\partial y} + 2p_{\mu_i} 2p_{\mu_j} \left(\frac{\partial}{\partial y} \right)^2. \quad (462)$$

It gives

$$\sum_{k=0, \pi}^{2k \leq w} \frac{1}{(w-2k)!(2k)!} \sum_{\sigma, \pi' := \sigma\pi} \frac{1}{2^k k!} \prod_{i=1}^k \left(2\delta_{\pi'_i \pi_{i+1}} \frac{\partial}{\partial y} \right) \prod_{i=2k+1}^w \left(2p_{\pi_i} \frac{\partial}{\partial y} \right), \quad (463)$$

where the outer and inner sums run over $w!$, $(2k)!$ permutations, respectively. Using the inequality $|A_{\vec{\mu}} B_{\vec{\nu}}| \leq |A_{\vec{\mu}}| |B_{\vec{\nu}}|$ we obtain the upper bound. \blacksquare

Lemma 49 *Let $C > 1$*

$$\left| \prod_{i=1}^w \frac{\partial}{\partial p_{\mu_i}} S^{\Lambda \Lambda_0}(p) \right| < 2w! e(C+1) \frac{(2^2 e^{\frac{3}{2}} C^2)^w}{(\Lambda + |p|)^{w+2}}. \quad (464)$$

Proof We change the variable $x_\mu = p_\mu/\Lambda$

$$\prod_{i=1}^w \frac{\partial}{\partial p_{\mu_i}} S^{\Lambda \Lambda_0}(p) = \frac{1}{\Lambda^{2w+2}} \prod_{i=1}^w \frac{\partial}{\partial x_{\mu_i}} h(x^2). \quad (465)$$

Equations (460), (453) yield

$$|\partial^w h(x^2)| < 2^w e(C+1) \sum_{k=0}^{2k \leq w} \frac{w!}{k!} \frac{(w-k)!}{(w-2k)!} \frac{x^{w-2k} (2\sqrt{e}C)^{w-k}}{(1+x^2)^{w-k+1}}. \quad (466)$$

\blacksquare

Lemma 50 *Let $C > 1$*

$$\left| \prod_{i=1}^w \frac{\partial}{\partial p_{\mu_i}} C_{\mu\nu}^{\Lambda\Lambda_0}(p) \right| < 2^3 w! e(C + 1 + (\xi + 1)(C^2 + 3)) \frac{(2^2 e^{\frac{3}{2}} C^2)^w}{(\Lambda + |p|)^{w+2}}. \quad (467)$$

Proof

$$|\partial^w C_{\mu\nu}^{\Lambda\Lambda_0}(p)| \leq 4 |\partial^w S^{\Lambda\Lambda_0}(p)| + |\xi - 1| \left| \partial^w \frac{p_\mu p_\nu}{p^2} S^{\Lambda\Lambda_0}(p) \right|. \quad (468)$$

The first term is bounded in (464). Using $x_\mu = p_\mu/\Lambda$ we have an upper bound for the last term

$$x^2 |\partial^w f(x^4)| + 2^3 w |x| |\partial^{w-1} f(x^4)| + 4^2 w(w-1) |\partial^{w-2} f(x^4)|, \quad (469)$$

where $f(x^4)$ is the same as in (449). Equation 449 gives

$$|\partial^w f(x^2)| < \frac{2w! e(2\sqrt{e}C)^w}{(1+x)^{w+2}}, \quad |\partial^w f(x^4)| < \frac{2^3 w! e(2^2 e^{\frac{3}{4}} C^2)^w}{(1+x)^{w+4}}. \quad (470)$$

Consequently,

$$|\partial^w x_\mu x_\nu f(x^4)| < 2^3 (C^2 + 3) e w! \frac{(2^2 C^2 e^{\frac{3}{4}})^w}{(1+x)^{w+2}}. \quad (471)$$

■

Lemma 51 *For all $p \in \mathbb{R}^4$ there exists a constant C such that*

$$|\dot{\mathbf{C}}^{\Lambda\Lambda_0}(p)| \leq C \frac{1}{\Lambda^3} e^{-\frac{p^2}{\Lambda^2}}. \quad (472)$$

Proof Using the inequality $x e^{-x^2} \leq x e^{1-2x} = x e^{1-x} e^{-x} \leq e^{-x}$ we obtain

$$|\dot{S}^{\Lambda\Lambda_0}(p)| = 4 \frac{1}{\Lambda^3} x e^{-x^2} \leq \frac{4}{\Lambda^3} e^{-x}, \quad x := \frac{p^2}{\Lambda^2}, \quad (473)$$

$$|\dot{C}_{\mu\nu}^{\Lambda\Lambda_0}(p)| \leq 4 \frac{1 + |\xi - 1|}{\Lambda^3} x e^{-x^2}. \quad (474)$$

■

Lemma 52 *For all $p \in \mathbb{R}^4$ there exists a constant C such that*

$$|\partial_p^w \partial_{\Lambda_0} \mathbf{C}^{\Lambda\Lambda_0}(p)| \leq C \frac{1}{\Lambda_0 (\Lambda_0 + |p|)^{2+\|w\|}}. \quad (475)$$

Proof Let $f = e^{-x^4}xx$ then using $|\partial_x^w|g| \leq |\partial_x^w g|$ we have

$$|\partial_p^w \partial_{\Lambda_0} C^{\Lambda\Lambda_0}(p)| \leq \frac{4(2 + |\xi - 1|)|\partial_x^w f|}{\Lambda_0^{\|w\|+3}}, \quad |\partial_p^w \partial_{\Lambda_0} S^{\Lambda\Lambda_0}(p)| \leq \frac{4|\partial_x^w f|}{\Lambda_0^{\|w\|+3}}, \quad (476)$$

where introducing $C_1 = 2^{\frac{3}{2}w_{max}}(4w_{max} + 2)!$

$$|\partial_x^w f| \leq C_1 e^{-x^4}(|x|^{3\|w\|+2} + 1) \leq eC_1 e^{-2x^2}(|x|^{3\|w\|+2} + 1). \quad (477)$$

Then it is easy to see that for $0 \leq m \leq 3w_{max} + 2$

$$\max(e^{-x^2}(x^m + 1)) \leq \max(e^{-x^2}x^m) + 1 \leq e^{\frac{m}{2}(\log \frac{m}{2} - 1)} + 1 \leq C_2. \quad (478)$$

To go further we need the following inequality for all $y \geq 0$ and $k \in \mathbb{N}$

$$g_k(y) \leq k!, \quad g_k(y) := (1 + y)^k e^{-y}, \quad (479)$$

which is obtained looking for the maximum $\bar{y}_k = k - 1$ and using $g_0(\bar{y}_0) = 1$.

$$g_k(\bar{y}_k) = kg_{k-1}(\bar{y}_k) \leq kg_{k-1}(\bar{y}_{k-1}) \implies g_k(y) \leq g_k(\bar{y}_k) \leq k! \quad (480)$$

Defining $C_3 := e^2 C_1 C_2$ and using inequality (479) we have

$$|\partial_x^w f| \leq C_3 e^{-2x} \leq C_3 \frac{2!}{(1 + |x|)^2} \frac{\|w\|!}{(1 + x)^{\|w\|}} \leq \frac{2C_3 w_{max}!}{(1 + |x|)^{2+\|w\|}}. \quad (481)$$

■

Lemma 53 Let $\mathcal{C}^{\Lambda\Lambda_0}$ be one of the propagators $S^{\Lambda\Lambda_0}$ or $C_{\mu\nu}^{\Lambda\Lambda_0}$, as defined in (35), (36). There are positive constants c_0, c_1, d such that for all $w \in \mathbb{N}$, $p \in \mathbb{R}^4$, $0 < \Lambda \leq \Lambda_0$, and with $c_\xi := c_0 + \xi c_1$,

$$\left| \left(\prod_{i=1}^w \frac{\partial}{\partial p_{\mu_i}} \right) \mathcal{C}^{\Lambda\Lambda_0}(p) \right| < \frac{w! d^w c_\xi}{(|p| + \Lambda)^{w+2}}, \quad |\dot{\mathcal{C}}^{\Lambda\Lambda_0}(p)| < \frac{c_\xi}{\Lambda^3} e^{-\frac{p^2}{\Lambda^2}}. \quad (482)$$

Proof The statement follows from (464), (467), (472). ■

E The functional $\Gamma_{n \leq 4}$

We expand the generating functional $\tilde{\Gamma}^{0\Lambda_0}(\underline{A}, \underline{c}, \underline{\bar{c}})$ of (82) as formal series in $\underline{A}, \underline{c}, \underline{\bar{c}}$. As usual, we adopt the shorthand notation $\Gamma^{0\Lambda_0}(A, \bar{c}, c)$ for $\tilde{\Gamma}^{0\Lambda_0}(\underline{A}, \underline{c}, \underline{\bar{c}})$.

$$\Gamma^{0\Lambda_0} = \Gamma_{n \leq 4}^{0\Lambda_0} + \Gamma_{n \geq 5}^{0\Lambda_0}, \quad \Gamma_{n \leq 4}^{0\Lambda_0} := \sum_{n=1}^4 \Gamma_n^{0\Lambda_0}, \quad (483)$$

where n counts the number of fields. The functionals $\Gamma_n^{0\Lambda_0}$ with $n \leq 4$ contain both relevant and irrelevant terms. We assume hypothesis RC1. In general the tensors $\zeta_{\mu_1 \dots \mu_r}^{\vec{\phi}}$ appearing in the form factors $F_{\mu_1 \dots \mu_r}^{\vec{\phi}}(\vec{p})$ are elements of $\text{span}(\{\delta^s \underline{p}^{k > 0}\}_{r > 0})$ where $\underline{p} = (p_1, \dots, p_{n-1})$.

1. One-point function

There are no local terms that preserve Euclidean invariance and global SU(2) symmetry. It follows that $\Gamma_1 = 0$.

2. Two-point functions

$$\Gamma_2^{0\Lambda_0} = \frac{1}{2} \langle F_{\mu\nu}^{AA} \check{A}_\mu^a \check{A}_\nu^a \rangle + \langle F^{\bar{c}c} \check{c}^a \check{c}^a \rangle, \quad (484)$$

$$\begin{aligned} F_{\mu\nu}^{AA}(p) &:= (\delta_{\mu\nu} p^2 - p_\mu p_\nu) (\sigma_{0\Lambda_0}^{-1}(p^2) + \Sigma_T^{AA}(p^2)) \\ &\quad + \frac{1}{\xi} p_\mu p_\nu (\sigma_{0\Lambda_0}^{-1}(p^2) + \Sigma_L^{AA}(p^2)), \end{aligned} \quad (485)$$

$$F^{\bar{c}c}(p) := -p^2 (\sigma_{0\Lambda_0}^{-1}(p^2) + \Sigma^{\bar{c}c}(p^2)). \quad (486)$$

We assume that the form factors Σ^{AA} and $\Sigma^{\bar{c}c}$ include all loop corrections. Note that for the functional $\underline{\Gamma}_2^{0\Lambda_0}$ we have

$$\underline{\Gamma}_2^{0\Lambda_0} = \frac{1}{2} \langle \underline{F}_{\mu\nu}^{AA} \check{A}_\mu^a \check{A}_\nu^a \rangle + \langle F^{\bar{c}c} \check{c}^a \check{c}^a \rangle, \quad (487)$$

$$\underline{F}_{\mu\nu}^{AA}(p) := F_{\mu\nu}^{AA}(p) - \frac{1}{2\xi} p_\mu p_\nu. \quad (488)$$

With $p^2 = M^2$ substitution of the above definitions into the expressions $\sigma_{0\Lambda_0} F^{AA}$, $\sigma_{0\Lambda_0} F^{\bar{c}c}$ appearing in AGE(123) and STI(124) gives

$$\sigma_{0\Lambda_0}(p^2) F^{\bar{c}c}(p) \stackrel{\beta}{\sim} -p^2 (1 + \Sigma^{\bar{c}c}(p^2)), \quad (489)$$

$$\sigma_{0\Lambda_0}(p^2) \underline{F}_{\mu\nu}^{AA}(p) \stackrel{\beta}{\sim} F_{T;\mu\nu}^{AA}(p), \quad (490)$$

$$F_{T;\mu\nu}^{AA}(p) := (\delta_{\mu\nu}p^2 - p_\mu p_\nu)(1 + \Sigma_T^{AA}(p^2)) + \frac{1}{\xi}p_\mu p_\nu \Sigma_L^{AA}(p^2). \quad (491)$$

Using (85) for the functional $\Gamma_2^{0\Lambda_0}$ we have

$$\Gamma_2^{0\Lambda_0} = \frac{1}{2}\langle F_{\mu\nu}^{AA} \check{A}_\mu^a \check{A}_\nu^a \rangle + \langle F^{\bar{c}c} \check{c}^a \check{c}^a \rangle, \quad (492)$$

$$F_{\mu\nu}^{AA}(p) := (\delta_{\mu\nu}p^2 - p_\mu p_\nu)\Sigma_T^{AA}(p^2) + \frac{1}{\xi}p_\mu p_\nu \Sigma_L^{AA}(p^2), \quad (493)$$

$$F^{\bar{c}c}(p) := -p^2 \Sigma^{\bar{c}c}(p^2). \quad (494)$$

For marginal terms we obtain

$$F_{\mu\nu}^{AA;p\rho p\sigma}(p) = 2\delta_{\mu\nu}\delta_{\rho\sigma}r_1^{AA} + 2(\delta_{\mu\rho}\delta_{\nu\sigma} + \delta_{\nu\rho}\delta_{\mu\sigma})r_2^{AA} + \zeta_{\mu\nu\rho\sigma}^{AA}, \quad (495)$$

$$F^{\bar{c}c;p\rho p\sigma}(p) = 2\delta_{\rho\sigma}r^{\bar{c}c}(p^2) + \zeta_{\rho\sigma}^{\bar{c}c}, \quad (496)$$

$$r_1^{AA}(p^2) := \Sigma_T^{AA}(p^2) + p^2 \frac{\partial \Sigma_T^{AA}(p^2)}{\partial p^2}, \quad (497)$$

$$r_2^{AA}(p^2) := \frac{1}{\xi}\Sigma_L^{AA}(p^2) - \Sigma_T^{AA}(p^2), \quad (498)$$

$$r^{\bar{c}c}(p^2) := -\Sigma^{\bar{c}c}(p^2) - p^2 \frac{\partial \Sigma^{\bar{c}c}(p^2)}{\partial p^2}. \quad (499)$$

3. Three-point functions

$$\Gamma_3^{0\Lambda_0} = \langle \epsilon_{abd} F_{\rho\mu\nu}^{AAA} \check{A}_\rho^a \check{A}_\mu^b \check{A}_\nu^d \rangle + \langle \epsilon_{adb} F_\mu^{A\bar{c}c} \check{A}_\mu^a \check{c}^b \check{c}^d \rangle, \quad (500)$$

$$F_\mu^{A\bar{c}c}(k, p, q) := ip_\mu R_1^{A\bar{c}c}(p, q) + iq_\mu r_2^{A\bar{c}c}(p, q), \quad (501)$$

$$R_1^{A\bar{c}c}(p, q) := g + r_1^{A\bar{c}c}(p, q), \quad (502)$$

$$F_{\rho\mu\nu}^{AAA}(k, p, q) := i\delta_{\mu\nu}(p_\rho - q_\rho)R^{AAA}(p, q) + i\delta_{\mu\nu}k_\rho \zeta_-^{AAA}(p, q) + i\zeta_{\rho\mu\nu}^{AAA}(p, q), \quad (503)$$

$$R^{AAA}(p, q) := \frac{1}{2}g + r^{AAA}(p, q). \quad (504)$$

Here $R^{AAA}(p, q)$ is a symmetric function whereas $\zeta_-^{AAA}(p, q)$ is antisymmetric.

4. Four-point functions

$$\Gamma_4^{0\Lambda_0} = \langle F_{\sigma\rho\mu\nu}^{AAAA} \check{A}_\sigma^b \check{A}_\rho^b \check{A}_\mu^a \check{A}_\nu^a + F_{1,\mu\nu}^{\bar{c}cAA} \check{c}^b \check{c}^b \check{A}_\mu^a \check{A}_\nu^a + F_{2,\mu\nu}^{\bar{c}cAA} \check{c}^a \check{c}^b \check{A}_\mu^a \check{A}_\nu^b + r^{\bar{c}c\bar{c}c} \check{c}^b \check{c}^b \check{c}^a \check{c}^a \rangle, \quad (505)$$

$$F_{n,\mu\nu}^{\bar{c}cAA} := \delta_{\mu\nu} r_n^{\bar{c}cAA} + \zeta_{n,\mu\nu}^{\bar{c}cAA}, \quad (506)$$

$$F_{\sigma\rho\mu\nu}^{AAAA} := R_{\sigma\rho\mu\nu}^{AAAA} + \zeta_{\sigma\rho\mu\nu}^{AAAA}, \quad (507)$$

$$R_{\sigma\rho\mu\nu}^{AAAA} := \frac{1}{2}(\delta_{\mu\rho}\delta_{\nu\sigma} + \delta_{\mu\sigma}\delta_{\rho\nu})R_1^{AAAA} + \delta_{\mu\nu}\delta_{\sigma\rho}R_2^{AAAA} + \frac{1}{2}(\delta_{\mu\rho}\delta_{\nu\sigma} - \delta_{\mu\sigma}\delta_{\rho\nu})\zeta_-^{AAAA}, \quad (508)$$

$$R_1^{AAAA} := -\frac{g^2}{4} + r_1^{AAAA}, \quad (509)$$

$$R_2^{AAAA} := \frac{g^2}{4} + r_2^{AAAA} \quad (510)$$

Here the terms

$$r_1^{AA}, \quad r_2^{AA}, \quad r^{\bar{c}c}, \quad R_1^{A\bar{c}c}, \quad r_2^{A\bar{c}c}, \quad R^{AAAA}, \quad R_1^{AAAA}, \quad R_2^{AAAA}, \quad r_1^{\bar{c}cAA}, \quad r_2^{\bar{c}cAA}, \quad r^{\bar{c}c\bar{c}c}, \quad (511)$$

are scalar functions of momenta, Λ_0 , and M . All 11 renormalization constants are fine-tuned by imposing appropriate renormalization conditions.

F The functionals $\Gamma_{\gamma;n \leq 2}$ and $\Gamma_{\omega;n \leq 2}$

With $\Gamma_{\varkappa}^{0\Lambda_0}(p) := \tilde{\delta}_{\varkappa(p)}\Gamma^{0\Lambda_0}|_{\varkappa=0}$,

$$\Gamma_{\gamma_\mu^a;n \leq 2}^{0\Lambda_0}(p) = R_1 i p_\mu c^a(p) + g \epsilon^{abs} \langle F_{\mu\nu}^{\gamma Ac} | \check{A}_\nu^b \check{c}^s; p \rangle, \quad (512)$$

$$F_{\mu\nu}^{\gamma Ac}(k, q) := \delta_{\mu\nu} R_2 + \zeta_{\mu\nu}^{\gamma Ac}(k, q), \quad (513)$$

$$\zeta_{\mu\nu}^{\gamma Ac}(k, q) := k_\mu q_\nu \zeta_1^{\gamma Ac} + k_\nu q_\mu \zeta_2^{\gamma Ac} + k_\mu k_\nu \zeta_3^{\gamma Ac} + q_\mu q_\nu \zeta_4^{\gamma Ac}, \quad (514)$$

$$\Gamma_{\omega^a;n \leq 2}^{0\Lambda_0}(p) = \frac{1}{2} g \epsilon^{abs} \langle R_3 | \check{c}^b \check{c}^s; p \rangle. \quad (515)$$

Here R_1, R_2, R_3 are scalar functions of momenta, Λ_0 , and M .

G Violated STI for $\Lambda > 0$

In this section we present, omitting the details of the calculation, the extension of some results of section 2.4 to arbitrary $0 < \Lambda < \Lambda_0$.

$$\tilde{\Gamma}_\beta^{\Lambda\Lambda_0} = \sigma_{0\Lambda_0} * \left(\frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \underline{\mathcal{C}}} - \partial \frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \gamma} \right), \quad (516)$$

$$\begin{aligned} \tilde{\Gamma}_1^{\Lambda\Lambda_0} = & \left\langle \frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \underline{A}}, \sigma_{0\Lambda_0} * \frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \gamma} \right\rangle - \left\langle \frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \underline{\mathcal{C}}}, \sigma_{0\Lambda_0} * \frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \omega} \right\rangle \\ & - \frac{1}{\xi} \left\langle \partial \underline{A}, \sigma_{0\Lambda_0} * \frac{\delta \tilde{\mathcal{F}}^{\Lambda\Lambda_0}}{\delta \underline{\mathcal{C}}} \right\rangle + \hbar \tilde{\Delta}^{\Lambda\Lambda_0}, \end{aligned} \quad (517)$$

$$\tilde{\mathcal{F}}^{\Lambda\Lambda_0} := \tilde{\Gamma}^{\Lambda\Lambda_0} + \frac{1}{2} \langle \tilde{\Phi}, \tilde{\mathbf{C}}_{0\Lambda_0}^{-1} \tilde{\Phi} \rangle, \quad (518)$$

$$\begin{aligned} \tilde{\Delta}^{\Lambda\Lambda_0} := & \langle (\sigma_\Lambda, 0, 0) \left(1 + \frac{\delta^2 \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \tilde{\Phi} \delta \tilde{\Phi}} \hat{\mathbf{1}} \tilde{\mathbf{C}}^{\Lambda\Lambda_0} \right)^{-1} \frac{\delta^2 \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \tilde{\Phi} \delta \gamma} \rangle \\ & + \langle (0, \sigma_\Lambda, 0) \left(1 + \frac{\delta^2 \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \tilde{\Phi} \delta \tilde{\Phi}} \hat{\mathbf{1}} \tilde{\mathbf{C}}^{\Lambda\Lambda_0} \right)^{-1} \frac{\delta^2 \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \tilde{\Phi} \delta \omega} \rangle \\ & - \langle (\sigma_\Lambda \partial, 0, 0) \left(1 + \frac{\delta^2 \tilde{\Gamma}^{\Lambda\Lambda_0}}{\delta \tilde{\Phi} \delta \tilde{\Phi}} \hat{\mathbf{1}} \tilde{\mathbf{C}}^{\Lambda\Lambda_0} \right)^{-1} \frac{\delta \underline{\mathcal{C}}}{\delta \tilde{\Phi}} \rangle. \end{aligned} \quad (519)$$

Note that in relations (516) and (517) there still appears a convolution with $\sigma_{0\Lambda_0}$ since we have chosen to define the regularized BRST transformation to include a convolution with this function, see (90).

Using the bounds of theorem 1 one can show that $\lim_{\Lambda \rightarrow 0} \Delta^{\Lambda\Lambda_0} = 0$ at nonexceptional momenta. It follows that $\lim_{\Lambda \rightarrow 0} \tilde{\Gamma}_1^{\Lambda\Lambda_0} = \tilde{\Gamma}_1^{0\Lambda_0}$ at nonexceptional momenta. More information can be found in [Efr17].

H List of the renormalization points

X	$[X]$	n_X	r_m	r_X	ren.p.
$\partial\partial\Gamma^{\bar{c}c}$	0	2	7	0	$\Lambda = 0, \vec{q} \in \mathbb{M}_2^s$
$\partial\Gamma_\gamma^c$	0	2	7	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_2^s$
$\partial\partial\Gamma^{AA}$	0	2	7	4	$\Lambda = 0, \vec{q} \in \mathbb{M}_2^s$
Γ_ω^{cc}	0	3	5	0	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
$\partial_A\Gamma^{c\bar{c}A}$	0	3	5	2	$\Lambda = M, \vec{q} = 0$
$\partial_c\Gamma^{c\bar{c}A}$	0	3	5	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
Γ_γ^{cA}	0	3	5	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
Γ_γ^{AA}	0	3	5	3	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
$\partial\Gamma^{AAA}$	0	3	5	4	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
$\Gamma^{\bar{c}\bar{c}c}$	0	4	3	0	$\Lambda = M, \vec{q} = 0$
$\Gamma^{\bar{c}cAA}$	0	4	3	2	$\Lambda = M, \vec{q} = 0$
Γ^{AAAA}	0	4	3	4*	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^{cp}$
$\partial\partial\Gamma_\beta^c$	0	2	7	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_2^s$
$\partial\partial\partial\Gamma_1^{cA}$	0	2	7	4	$\Lambda = 0, \vec{q} \in \mathbb{M}_2^s$
$\partial\partial\Gamma_1^{c\bar{c}c}$	0	3	5	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
$\partial\Gamma_{1,\gamma}^{cc}$	0	3	5	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
$\partial\partial\Gamma_1^{cAA}$	0	3	5	4	$\Lambda = 0, \vec{q} \in \mathbb{M}_3^s$
$\Gamma_\beta^{c\bar{c}c}$	0	4	3	0	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^s$
$\Gamma_{1,\omega}^{ccc}$	0	4	3	0	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^s$
Γ_β^{cAA}	0	4	3	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^s$
$\partial\Gamma_1^{c\bar{c}cA}$	0	4	3	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^s$
$\Gamma_{1,\gamma}^{ccA}$	0	4	3	2	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^s$
$\partial\Gamma_1^{cAAA}$	0	4	3	4*	$\Lambda = 0, \vec{q} \in \mathbb{M}_4^{cp}$
$\Gamma_1^{c\bar{c}\bar{c}c}$	0	5	1	0	$\Lambda = 0, \vec{q} \in \mathbb{M}_5^s$
$\Gamma_1^{c\bar{c}cAA}$	0	5	1	2*	$\Lambda = 0, \vec{q} \in \mathbb{M}_5^s$
Γ_1^{cAAAA}	0	5	1	4*	$\Lambda = 0, \vec{q} \in \mathbb{M}_5^{cp}$

X	$[X]$	n_X	r_X	ren.p.
$\partial\Gamma^{\bar{c}c}$	1	2	1	$\Lambda = 0, \vec{q} = 0$
$\partial\Gamma^{AA}$	1	2	3	$\Lambda = 0, \vec{q} = 0$
$\Gamma^{\bar{c}cA}$	1	3	1	$\Lambda = 0, \vec{q} = 0$
Γ^{AAA}	1	3	3	$\Lambda = 0, \vec{q} = 0$
$\Gamma^{\bar{c}c}$	2	2	0	$\Lambda = 0, \vec{q} = 0$
Γ^{AA}	2	2	2	$\Lambda = 0, \vec{q} = 0$
$\partial\Gamma_\beta^c$	1	2	1	$\Lambda = 0, \vec{q} = 0$
$\partial\partial\Gamma_1^{cA}$	1	2	3	$\Lambda = 0, \vec{q} = 0$
Γ_β^{cA}	1	3	1	$\Lambda = 0, \vec{q} = 0$
$\partial\Gamma_1^{c\bar{c}c}$	1	3	1	$\Lambda = 0, \vec{q} = 0$
$\Gamma_{1,\gamma}^{cc}$	1	3	1	$\Lambda = 0, \vec{q} = 0$
$\partial\Gamma_1^{cAA}$	1	3	3	$\Lambda = 0, \vec{q} = 0$
$\Gamma_1^{c\bar{c}cA}$	1	4	1	$\Lambda = 0, \vec{q} = 0$
Γ_1^{cAAA}	1	4	3	$\Lambda = 0, \vec{q} = 0$
Γ_β^c	2	2	0	$\Lambda = 0, \vec{q} = 0$
$\partial\Gamma_1^{cA}$	2	2	2	$\Lambda = 0, \vec{q} = 0$
$\Gamma_1^{c\bar{c}c}$	2	3	0	$\Lambda = 0, \vec{q} = 0$
Γ_1^{cAA}	2	3	2	$\Lambda = 0, \vec{q} = 0$
Γ_1^{cA}	3	2	1	$\Lambda = 0, \vec{q} = 0$

List of all terms preserving the global symmetries, with an arbitrary number of γ, ω insertions, and with at most one β or 1 insertion (not both). Notation: $[X]$ is the mass dimension of X (reduced Fourier transform); r_X is the tensor rank; n_X is the total number of fields and sources (not including 1); ∂ stands for a momentum derivative; “ren.p.” stands for “renormalization point”. A * in the rank entry means that the condition $r_m := 11 - 2n_X \geq r + 1$ is violated for a term X : as stated in Lemma 28 the tensor monomials $\{\delta^s \underline{q}^k\}_{r+1}$ are not linearly independent for $\vec{q} = (q_0, \underline{q}) \in \mathbb{M}_{n_X}^s$, hence they are not suitable as a basis for the form-factor decomposition of ∂X . See lemma 31 and sections 5.4, 5.5, 5.10.

I List of insertions

X	$[X]$	gh(X)	def
ψ_μ^a	2	1	(88)
γ_μ^a	2	-1	
Ω^a	2	2	(88)
ω^a	2	-2	
Q_ρ	5	1	(100)
$Q_{\rho\gamma}$	3	2	(95)
$Q_{\rho\omega}$	3	3	(96)
ρ	-1	-1	
Q_β	3	1	(101)
β	1	-1	

List of operators and sources, and their quantum numbers. Notation: $[X]$ stands for the mass dimension of X in position space; the ghost charge of the ghost field is $\text{gh}(c) := 1$.

References

- [BDM93] M. Bonini, M. D’Attanasio, and G. Marchesini. Perturbative renormalization and infrared finiteness in the Wilson renormalization group: the massless scalar case. *Nucl. Phys. B*, 409:441 – 464, 1993.
- [BDM95] M. Bonini, M. D’Attanasio, and G. Marchesini. BRS symmetry for Yang-Mills theory with exact renormalization group. *Nucl. Phys. B*, 437:163–186, 1995.
- [Bec96] C. Becchi. On the construction of renormalized gauge theories using renormalization group techniques. *arXiv*, 9607188v1, 1996.
- [Ber66] F.A. Berezin. *The method of second quantization. Translated by N. Mugibayashi and A. Jeffrey.* Academic Press New York, 1966.
- [BRS75] C. Becchi, A. Rouet, and R. Stora. Renormalization of the abelian Higgs-Kibble model. *Comm. Math. Phys.*, 42(2):127–162, 1975.
- [DF91] Yu.L. Dalecky and S.V. Fomin. *Measures and differential equations in infinite-dimensional space.* Springer, 1991.
- [Efr17] A.N. Efremov. *Renormalization of $SU(2)$ Yang-Mills theory with flow equations.* PhD thesis, Université Paris-Saclay, 2017.
- [FHH16] M.B. Fröb, J. Holland, and S. Hollands. All-order bounds for correlation functions of gauge-invariant operators in Yang-Mills theory. *Journal Math. Phys.*, 57(12), 2016.
- [FP67] L.D. Faddeev and V.N. Popov. Feynman diagrams for the Yang-Mills field. *Physics Letters B*, 25:29–30, 1967.
- [GG00] P. Gambino and P. A. Grassi. Nielsen identities of the SM and the definition of mass. *Phys. Rev. D*, 62(7):076002, 2000.
- [GJ87] J. Glimm and A. Jaffe. *Quantum physics.* Springer, 1987.
- [GK] R. Guida and Ch. Kopper. All-order uniform momentum bounds for the massless ϕ_4^4 field theory. *In preparation.*

- [GK11] R. Guida and Ch. Kopper. All-order uniform momentum bounds for the massless ϕ^4 theory in four dimensional euclidean space. *arXiv*, 1103.5692, 2011.
- [KK92] G. Keller and Ch. Kopper. Perturbative renormalization of composite operators via flow equations. *Comm. Math. Phys.*, 148:445–467, 1992.
- [KKS92] G. Keller, Ch. Kopper, and M. Salmhofer. Perturbative renormalization and effective lagrangians in ϕ^4 in four-dimensions. *Helv. Phys. Acta*, 65:32–52, 1992.
- [KKS97] G. Keller, Ch. Kopper, and C. Schophaus. Perturbative renormalization with flow equations in minkowski space. *Helv.Phys.Acta*, (70):247–274, 1997.
- [KM02] Ch. Kopper and F. Meunier. Large momentum bounds from flow equations. *Annales Henri Poincaré*, 3:435–449, 2002.
- [KM09] Ch. Kopper and V.F. Müller. Renormalization of spontaneously broken SU(2) Yang-Mills theory with flow equations. *Rev. in Math. Phys.*, 21(6):781–820, 2009.
- [KSZ75] H. Kluberg-Stern and J.B. Zuber. Ward identities and some clues to the renormalization of gauge-invariant operators. *Phys. Rev. D*, 12:467–481, Jul 1975.
- [Lai81] C.H. Lai, editor. *Gauge theory of weak and electromagnetic interactions*. World Scientific, 1981.
- [LZJ72] B.W. Lee and J. Zinn-Justin. Spontaneously broken gauge symmetries. I,II,III. *Phys. Rev. D*, 5:3121–3137, 3137–3155,3155–3160., 1972.
- [MD96] T.R. Morris and M. D’Attanasio. Gauge Invariance, the Quantum Action Principle, and the Renormalization Group. *Physics Letters B*, 378:213–221, 1996.
- [Mor94] T.R. Morris. The exact renormalization group and approximate solutions. *Int. Journ. of Mod. Phys. A*, 09(14):2411–2449, 1994.

- [Mül03] V.F. Müller. Perturbative renormalization by flow equations. *Rev. in Math. Phys.*, 15(05):491–558, 2003.
- [Nie75] N.K. Nielsen. On the gauge dependence of spontaneous symmetry breaking in gauge theories. *Nucl. Phys. B*, 101:173–188, 1975.
- [PF67] V.N. Popov and L.D. Faddeev. Perturbation theory for gauge-invariant fields. *preprint ITF-67-036 Kiev*, 1967. In russian. English translation: NAL-THY-57.
- [Pol84] J. Polchinski. Renormalization and effective Lagrangians. *Nucl. Phys. B*, 231:269–295, 1984.
- [PS85] O. Piguet and K. Sibold. Gauge independence in ordinary Yang-Mills theories. *Nucl. Phys. B*, 253:517–540, 1985.
- [Ree73] M. Reed. Functional analysis and probability theory. In G. Velo and A. Wightman, editors, *Constructive Quantum Field Theory*, volume 25 of *Lecture Notes in Physics*. Springer, 1973.
- [RW94] M. Reuter and C. Wetterich. Effective average action for gauge theories and exact evolution equations. *Nucl. Phys. B*, 417:181–214, 1994.
- [SF91] A.A. Slavnov and L.D. Faddeev. *Gauge fields: an introduction to quantum theory*. Addison-Wesley Publishing Company, 1991.
- [Sla72] A.A. Slavnov. Ward identities in gauge theories. *Theo. and Math. Phys.*, 10(2):99–104, 1972.
- [Tay71] J.C. Taylor. Ward identities and charge renormalization of the Yang-Mills field. *Nucl. Phys. B*, 33:436–444, 1971.
- [tH71] G. 't Hooft. Renormalization of massless Yang-Mills fields. *Nucl. Phys. B*, 33:173–199, 1971.
- [tHV72] G. 't Hooft and M. Veltman. Regularization and renormalization of gauge fields. *Nucl. Phys. B*, 44:189–213, 1972.
- [Tyu75] I.V. Tyutin. Gauge invariance in field theory and statistical physics in operator formalism. *arXiv*, 0812.0580, 1975.

- [Wet93] C. Wetterich. Exanct evolution equation for the effective potential. *Physics Letters B*, 301:90–94, 1993.
- [WH73] F. Wegner and A. Houghton. Renormalization group equations for critical phenomena. *Phys. Rev. A*, 8(1):401–412, 1973.
- [Wil71] K.G. Wilson. Renormalization group and critical phenomena. *Phys. Rev. B*, 4(9):3174–3183, 1971.
- [YM54] C.N. Yang and R.L. Mills. Conservation of isotopic spin and isotopic gauge invariance. *Phys. Rev.*, 96:191–195, Oct 1954.
- [ZJ75] J. Zinn-Justin. *Trends in elementary particle theory*, volume 37. Springer Berlin Heidelberg, 1975. International summer institute on theoretical physics in Bonn 1974.