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Excited State Pathways Leading to Formation of Adenine Dimers

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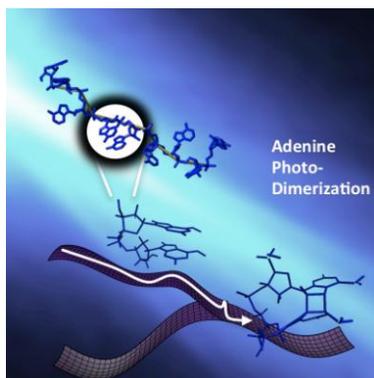
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ABSTRACT

The reaction intermediate in the path leading to UV-induced formation of adenine dimers A=A and AA* is identified for the first time quantum mechanically, using PCM/TD-DFT calculations on (dA)₂. In parallel, its fingerprint is detected in the absorption spectra recorded on the millisecond time-scale for the single strand (dA)₂₀.



The characterization of photochemical reactions of nucleic acids is important in connection with the UV-induced damage to the genetic code. Although a large number of publications deal with pyrimidine dimerization, very few focus on adenine dimer formation.¹⁻⁸ Yet, the photodimerization quantum yield of adenine single strands is the order of 10^{-3} ,⁶ comparable to that of (6-4)pyrimidine-pyrimidone adducts (64PPs) in thymine single strands,⁹ which have received considerable attention.¹⁰ Intriguingly, the photodimerization yields dramatically drop in DNA duplexes and in RNA strands.^{3, 6, 11} Experimental studies demonstrated that irradiation around 260 nm leads to the formation of two different types of adenine dimers AA* and A=A (Fig. 1) and suggested that they have a common azetidine intermediate.⁵ Here, we report the first theoretical study exploring the excited state pathways that lead to the azetidine precursor. The fingerprint of the reaction intermediate, which, according to our calculations, presents weak electronic transitions at longer wavelengths compared to AA* and A=A, is detected in the absorption spectra recorded at 20 ms following 266 nm excitation by 5 ns laser pulses of the single strand (dA)₂₀.

Using extensive molecular dynamics (MD) analysis of (dA)₅ (see SI, Fig. SI-1), we found that a significant percentage of adenines undergoes *syn/anti* conformational equilibrium, *syn* glycosidic conformers being particularly stable for the 5'-terminal base. Then, the four (AA) dinucleotides corresponding to the different combinations of conformers, namely (AA)_{*anti-anti*}, (AA)_{*syn-anti*}, (AA)_{*anti-syn*}, and (AA)_{*syn-syn*}, have been optimized at the M052X/6-31G(d) level,¹² including bulk solvent effects by means of the Polarizable Continuum Model (PCM).¹³ We have focused on C2'-endo/C2'-endo conformers, since MD simulations predict that most of the 'inner bases' (those more representative of the

behavior in long strands) adopts a C2'-endo puckering. At the PCM/M052X level of theory, the most stable conformer is (AA)_{syn-anti} (Fig. SI-3 and Table SI-1), which is stabilized by a strong hydrogen bond between the terminal CH₂OH group and N3 atom of the 5'-terminal adenine (⁵N3) (Fig. 1). When the molecules of the first solvation shell (see SI, Fig. SI-2) are included, (AA)_{anti-anti} is found to be the most stable conformer, but the difference in energy with respect to (AA)_{syn-anti} is only 4.0 kcal/mol, suggesting that a very small, but non-zero, percentage of the AA steps adopts this conformation.

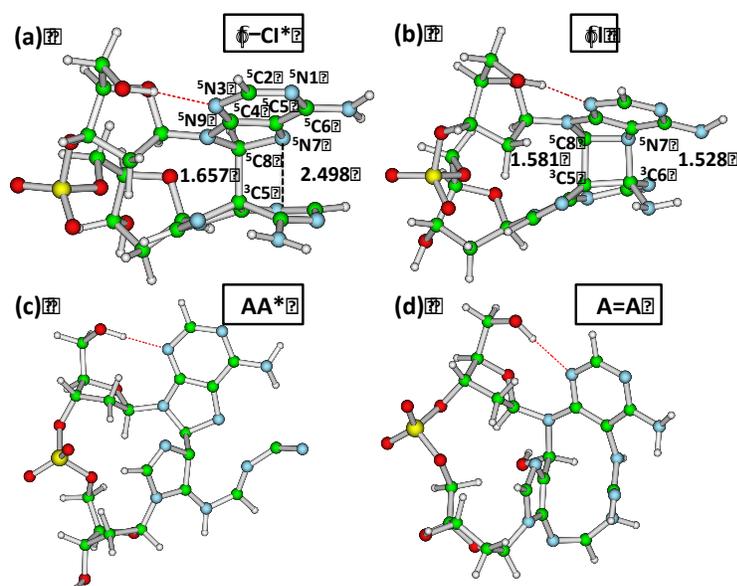


Figure 1. Schematic drawing, atom labeling and selected bond distances (in Å) of relevant structures in the photodimerization path of (AA)_{syn-anti} conformers of (dA)_n oligonucleotides: (a) ϕ -CI*, corresponding to the conical intersection, (b) reaction intermediate ϕ I, (c) final photo-product AA* and (d) final photo-product A=A.

Excited state PCM/TD-M052X calculations, which have been successfully applied to the study of DNA photochemistry,^{9, 14-15} indicate that, for all conformers, the two lowest energy excited states in the Franck-Condon region correspond to symmetric/antisymmetric combinations (Frenkel excitons) of the lowest energy $\pi\pi^*$

excited state of adenine (Table SI-2), denoted as L_a . Geometry optimization of the S_1 excited state leads to minima (Fig. 2 and Fig. SI-3) characterized by a close approach of the two bases. The S_1 minimum found for $(AA)_{syn-anti}$ exhibits very short ${}^5C8-{}^3C5$ and ${}^5N7-{}^3C6$ distances (the superscripts 5 and 3 correspond to the bases at the 5' and 3'-end, respectively). The HOMO \rightarrow LUMO electronic transition involves transfer of an electron to an orbital with bonding character in respect to these atom pairs (Fig. 2). An extensive analysis of the potential energy surfaces along the ${}^5C8-{}^3C5$ coordinate shows that an extremely low energy barrier (<0.08 eV) separates the S_1 minimum from a crossing region with S_0 in the path leading to photodimerization.

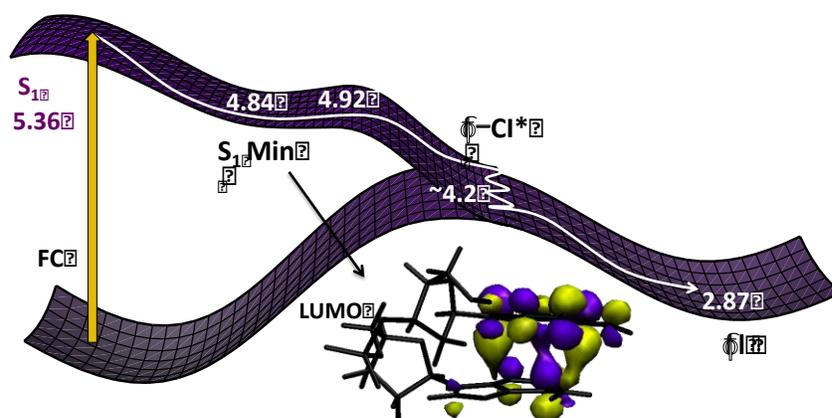


Figure 2. Schematic description of the potential energy surface in the path leading from the Franck-Condon state to the minimum of the S_1 excited state and, then, to the formation of the azetidine precursor ϕ . PCM/TD-M052X/6-31G(d) calculations for the in $(AA)_{syn-anti}$ conformer. The energy is given in eV.

In general, TD-DFT cannot provide an accurate description of conical intersections. However, in the case of several photodimerization processes in DNA, the description obtained for the S_1/S_0 crossing region by TD-M052X is very similar to that obtained by wavefunction based methods.^{9, 14-15} A representative structure corresponding to this

region (ϕ -CI*), is shown in Fig. 1a. Subsequent S_1 geometry optimization leads to the photo-adduct ϕ I presented in Fig. 1b.

According to our calculations, the three other conformers do not exhibit any significant tendency to photodimerization (see SI for details, Figs. SI-4, 5 and 6). The general picture obtained for $(AA)_{anti-syn}$ is qualitatively similar to that shown in Fig. 2 but the energy barrier separating the S_1 minimum from the crossing region is much higher (≥ 0.33 eV). In the S_1 minimum of $(AA)_{syn-syn}$ the 5C8 and 3C8 atoms of adenines are very close (~ 1.6 Å); further approach leads to a crossing region with S_0 and decay back to the ground state of $(AA)_{syn-syn}$. For $(AA)_{anti-anti}$, the calculation along the 5C8 - 3C5 coordinate involves a large increase in the energy (> 0.5 eV), making the access to ϕ I impossible. Finally, in the case of $(AA)_{anti-anti}$, we located another possible photo-adduct (see SI) where 5C5 - 3C5 and 5C6 - 3C6 are bonded. However, the latter photo-adduct is less stable than ϕ I by 1.2 eV. Therefore, ϕ I appears the only relevant intermediate in the path to A=A and AA* and corresponds to the species postulated 25 years ago by Sharma and Davies.⁵

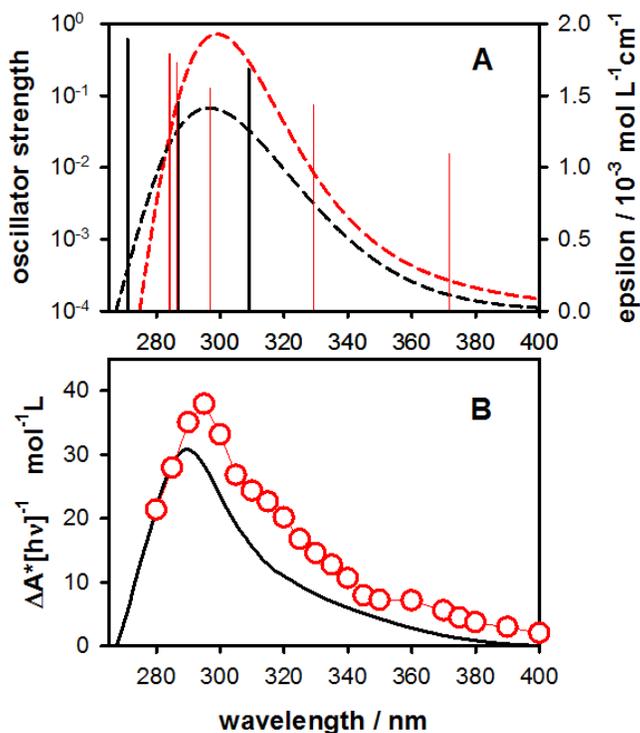


Figure 3. (A) Low energy electronic transitions of A=A and AA* dimers (black) and the precursor ϕI (red); PCM/TD-CAM-B3LYP/6-31+G(d,p)//PCM/M052X/6-31G(d) calculations. The energies have been homogenously decreased by 0.45 eV, which is the difference between the experimental absorption maximum of adenine and the vertical absorption energy computed at the adopted level of theory. In dashes: calculated differential absorption spectra corresponding to the disappearance of reacting adenines and the formation of ϕI (red) and A=A and AA* (black, assuming that their yields are roughly equal, in line with literature⁶); the oscillator strength is proportional to the spectral area; the spectral width is 0.4 eV. (B) Spectrum obtained for (dA)₂₀ at 20 ms (red) compared to the difference of steady-state absorption spectra recorded after and before irradiation (black). The spectra are normalized by the concentration of absorbed photons [hv].

The computed absorption spectrum of ϕI (Fig. SI-7) exhibits two transitions located at energies lower than those of both the reacting adenines (AA) and the dimers A=A and AA* by at least 0.25 and 0.6 eV (Fig. 3A). This finding suggests that, at wavelengths longer than 300 nm, the ϕI spectrum is more intense compared to those of the reactants and the final photoproducts. On this ground, we searched the fingerprints of ϕI in the time-resolved spectra recorded by laser flash photolysis. This technique has been applied

successfully to study the formation of 64PPs in thymine single strands from the oxetane intermediate.¹⁶

The studied adenine eicosamers (dA)₂₀ (HPLC purified from Eurogentec Europe) were dissolved in phosphate buffer (0.15 mol·L⁻¹ NaH₂PO₄, 0.15 mol·L⁻¹ Na₂HPO₄). The solutions were excited at 266 nm with 5 ns pulses. We stress that such experiments are very delicate and specific experimental protocols, preventing excitation of UV-damaged oligomers are required (see SI). The average power used in our experiments varied from 0.5 to 2x10⁶ W·cm⁻². Under such experimental conditions, in addition to UV-induced dimerization, ionization processes give rise to adenine radicals.¹⁷ The typical absorption bands of the adenine radical are easily distinguishable in the (dA)₂₀ spectrum recorded at 1.7 μs. This is shown in Fig. 4A together with the spectrum of the deprotonated radical cation reported by Candeias and Steenken.¹⁸ The intensity of the band at 600 nm varies linearly with the concentration of ejected hydrated electrons [e⁻] (Fig. SI-8), determined from their absorption band peaking at 720 nm.^{17, 19} This is not the case for the UV band, whose shape changes with the intensity of the excitation pulse. When the ratio of ejected electrons versus absorbed photons [e⁻]/[hv] drops from 1.5x10⁻³ to 0.3x10⁻³, we can clearly distinguish a peak at 300 nm while the radical peak at 335 nm appears only as a shoulder (inset in Fig. 4A).

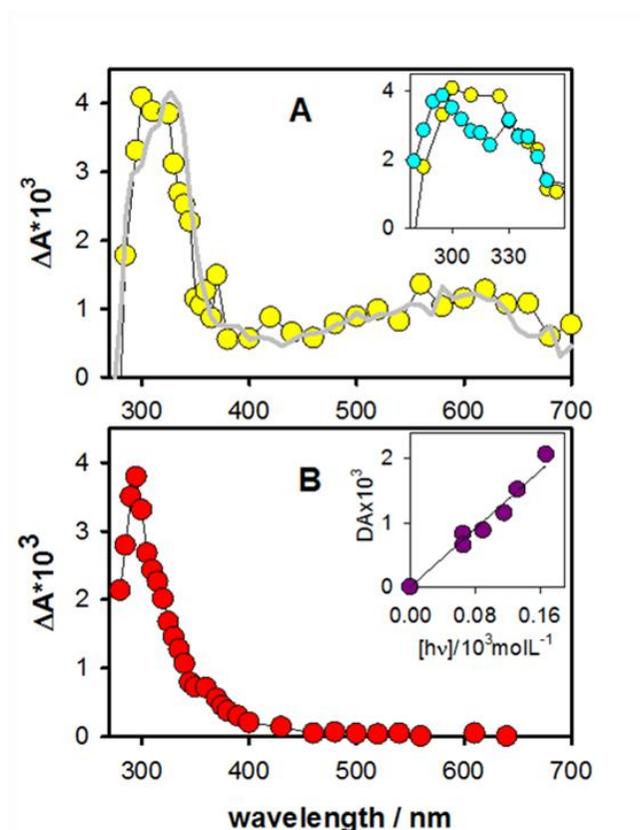


Figure 4. Time-resolved absorption spectra (circles) obtained for (dA)₂₀ at 1.7 μs (**A**) and 20 ms (**B**) following laser excitation at 266 nm. The grey line in **A** corresponds to the spectrum of the adenine radical (from reference ¹⁸; in arbitrary units). Inset in **A**: normalized spectra obtained for different ratios of ejected electron versus absorbed photons; 1.5 × 10⁻³ (yellow) and 0.3 × 10⁻³ (cyan). Inset in **B**: dependence of the transient absorbance at 340 nm and 20ms on the concentration of absorbed photons.

At much longer times, when the adenine radicals have disappeared, we observe a large absorption band peaking at 300 nm with a weak tail extending to 420 nm (Fig. 4B). In contrast to the species issued from electron ejection, the intensity of the band at 20 ms scales with the concentration of absorbed photons (inset in Fig. 4B). The transient spectra obtained for aerated and Argon bubbled solutions are practically the same. These observations indicate that the chemical species responsible for the 300 nm band, which appears at times shorter than the time resolution of the experimental setup (50 ns; Fig. SI-

9) originates from a singlet excited state. We have not observed any significant decay within 100 ms. Beyond this time, the instability of the detector hinders the measurements. The spectrum at 20 ms is compared in Fig. 3B with the difference of the steady-state spectra obtained before and after 266 nm irradiation. The same differential spectrum is obtained using either a continuous light source or laser pulses of varying energy. In all cases, the spectrum intensity scales with the number of absorbed photons (Fig. SI-10) and it is not altered for several hours after irradiation; from these steady-state spectra and using the average molar absorption coefficient of the adenine dimers reported by Pörschke for (dA)₂,¹ we estimated that the dimerization quantum yield is $\sim 7 \times 10^{-3}$, which is of the same order of magnitude as that reported by Clingen and Davies for poly(dA) (3×10^{-3}).⁶

The comparison of the two spectra in Fig. 3B shows that spectral changes occur between 100 ms and 5 min, which is the time between irradiation and recording of the steady-state spectrum. The peak shifts from 300 to 295 nm and the intensity between 295 and 400 nm decreases significantly. These modifications compare nicely with the differences in the corresponding theoretical spectra (Fig. 3A), computed assuming that the two final photoproducts are formed with the same quantum yield.

In view of the above discussion, we assign the transient spectrum to the precursor ϕI . Its conversion to A=A and AA* is obviously slower compared with the evolution of oxetane to 64PPs in thymine strands, which takes place within 4 ms, suggesting that the stability of the studied azetidine intermediate is larger than that of the oxetane.¹⁶ The stability of four-ring dimerization precursors has been correlated with the electronegativity of the involved heteroatoms. In the presence of sulphur atom, which is less electronegative than nitrogen and

oxygen atoms, the reaction intermediate thietane is so stable that it has been isolated and characterized.²⁰ Thus, it is not surprising that the Nitrogen containing precursor ϕI , is less stable than thietane but longer lived than oxetane.

In conclusion, we studied adenine dimerization combining theoretical tools (molecular dynamics simulations and quantum mechanical methods) with steady-state and time-resolved absorption spectroscopy. We identified a reaction intermediate which is formed from a singlet excited state and characterized by a lifetime comprised between 0.1 s and a few minutes. We showed that the photodimerization reaction is governed by conformational equilibria. Different interbase stacking geometries lead to different patterns of bonding/non-bonding interactions between atoms of the two adenines, and, therefore, determine which photodimers can be potentially formed without a significant energy barrier. Moreover, some of the possible photoproducts do not correspond to stable minima, as shown by our results on $(\text{AA})_{\text{syn-syn}}$, or they are less stable than the reactants. Consequently, adenine photodimerization is possible only for *syn-anti* conformers. This finding explains why base pairing, dramatically destabilizing *syn* conformers, diminishes drastically the quantum yield of the UV-induced adenine dimers.⁶ It would be interesting to perform the same type of investigation for other photodimerization reactions, not only in DNA strands but also in other multichromophoric systems whose geometry is characterized by many degrees of freedom.

Supporting Information: Computational Details, Additional Computational Results, Experimental Details, Additional Experimental Results, Cartesian Coordinates.

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The authors declare no competing financial interests.

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